

LINE VORTICES IN THE U(1) - HIGGS MODEL

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ABSTRACT. For a given U(1)-bundle E over $M = \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$, where the x_i are n distinct points of \mathbb{R}^3 , we minimise the U(1)-Higgs action and we make an asymptotic analysis of the minimizers when the coupling constant tends to infinity. We prove that the curvature (= magnetic field) converges to a limiting curvature that we give explicitly and which is singular along line vortices which connect the x_i . This work is the three dimensional equivalent of previous works in dimension two (see [3] and [4]). The results presented here were announced in [12].

1. INTRODUCTION

Consider n distinct points $\{x_1, \dots, x_n\}$ in \mathbb{R}^3 and E a U(1)-bundle over $M_{\bar{x}} = \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$.

Recall that such a bundle is uniquely determined by its first Chern class $c_1(E) \in H^2(M_{\bar{x}}, \mathbb{Z}) = \mathbb{Z} \oplus \dots \oplus \mathbb{Z}$ (n summands), that is n integers d_1, \dots, d_n which correspond to the integrals of $c_1(E)$ on small spheres surrounding the points x_1, \dots, x_n . In other words E is uniquely determined by a curvature $h_0 \in H^2(M_{\bar{x}})$ on E and for such a curvature one has

$$\frac{1}{2\pi} \int_S h_0 = \sum_{i: x_i \in \text{Int}(S)} d_i,$$

where S is any closed, oriented surface of \mathbb{R}^3 enclosing some of the points x_i in its interior, $\text{Int}(S)$.

We suppose that $\forall i \ d_i \neq 0$ and we make the following neutrality hypothesis

$$\sum_{i=1 \dots n} d_i = 0.$$

Under the previous hypothesis, we can write the sequence of points $\{x_i, \dots, x_n\}$ and the corresponding degrees $\{d_1, \dots, d_n\}$ in the form $\{P_1, \dots, P_k, N_1, \dots, N_k\}$ where the P_i (resp. N_i) correspond to the x_i having positive (resp. negative) integer d_i and are repeated as many time as to the multiplicity $|d_i|$.

This paper deals with a variational problem suggested in [8] by J. Fröhlich and M. Struwe. For any section φ of the \mathbb{C} -line bundle associated to E and

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any connection \tilde{A} on E , consider the $U(1)$ - Higgs action

$$S_\lambda(\varphi, \tilde{A}) = \int_{M_{\tilde{x}}} |\nabla_{\tilde{A}}\varphi|^2 + |d\tilde{A}|^2 + \frac{\lambda}{2}(1 - |\varphi|^2)^2,$$

where $\nabla_{\tilde{A}}\varphi$ is the covariant derivative of φ relative to \tilde{A} , $d\tilde{A}$ is the curvature of \tilde{A} and $V(|\varphi|) = \frac{\lambda}{2}(1 - |\varphi|^2)^2$ is the Higgs potential depending on $|\varphi|$, which, from a mathematical point of view, could be replaced, in the remain of the paper, by any “reasonable” potential on $|\varphi|$ (see [7] part 3) which forces $|\varphi|$ to be close to 1. The aim is to minimize $S_\lambda(\varphi, \tilde{A})$ and to describe the minimizers.

The words used above (section, connection, covariant derivative...) come from geometry but there are also parallel words coming from abelian gauge theory in physics, that we will mix with the previous ones, to designate the same objects. We will also say that φ is a scalar (Higgs) field, \tilde{A} is the gauge field (or vector potential), $d\tilde{A}$ is the magnetic field. We will also called $|\varphi|$ the density of the scalar field and $(i\varphi, \nabla_A\varphi)$ the current where $(,)$ is the usual scalar product on each fiber \mathbb{C} and $i\varphi$ the $\frac{\pi}{2}$ rotation of φ on the oriented fibre.

As it is pointed out in [8] for $|d_i| \neq 0$ $S(\varphi, \tilde{A}) = +\infty$, in fact for any $i \in [1, n]$ and r sufficiently small

$$\int_{\partial B_r(x_i)} d\tilde{A} = 2\pi d_i \quad ,$$

thus by Cauchy-Schwartz inequality

$$\frac{\pi|d_i|^2}{2r^2} \leq \int_{\partial B_r(x_i)} |d\tilde{A}|^2 \quad , \tag{1. 1}$$

integrating (1. 1) in r we get $\int |d\tilde{A}|^2 = +\infty$.

Thus we will adopt the renormalisation proposed in [8].

Consider the following closed 2-form on $M_{\tilde{x}}$

$$h_0 = -\frac{1}{2} \sum_{i=1, \dots, n} d_i \quad * d \left(\frac{1}{|x - x_i|} \right). \tag{1. 2}$$

Where $*$ denote the Hodge operator. h_0 is a curvature on E and it verifies the following equations

$$\begin{cases} dh_0 &= 2\pi \sum_{i=1}^n d_i \delta_{x_i} & \text{in } \mathcal{D}'(\mathbb{R}^3) \\ d^*h_0 &= 0 & \text{in } \mathcal{D}'(\mathbb{R}^3) \end{cases} \tag{1. 3}$$

where $d^* = *d*$. h_0 will be called the reference curvature (or magnetic field). Consider a connection \tilde{A} on E and $\tilde{h} = d\tilde{A}$ it's curvature, we denote by h the difference $h = \tilde{h} - h_0$. Formally since $dh = 0$ in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ and $\int_{\partial\Omega} h.\nu = 0$ for any regular bounded domain Ω of \mathbb{R}^3 , thus there exists a

global 1-form A on \mathbb{R}^3 such that $h = dA$ and

$$\begin{aligned} \int_{\mathbb{R}^3} (\tilde{h} - h_0).h_0 &= \int_{\mathbb{R}^3} h.h_0 = \int_{\mathbb{R}^3} dA \wedge *h_0 \\ &= - \int_{\mathbb{R}^3} A \wedge d * h_0 \\ &= 0 \end{aligned} \tag{1. 4}$$

(The previous affirmations and equalities could be justified rigourously by assuming some regularity on \tilde{h} in the neighborhood of the x_i and infinity. For instance suppose $\tilde{h} \in L^p(\mathbb{R}^3)$ for some $1 < p < \frac{3}{2}$ and $d^*\tilde{h} \in L^q(\mathbb{R}^3)$ for some $q > \frac{3}{2}$. This would imply $d\tilde{h} = 2\pi \sum_{i=1}^n d_i \delta_{x_i}$ in $\mathcal{D}'(\mathbb{R}^3)$ and the existence of $A \in W_{loc}^{1,s}(\mathbb{R}^3)$ for some $s > 3$ such that $h = dA$ in $\mathcal{D}'(\mathbb{R}^3)$).

In view of (1. 4) we deduce that, by considering $\int_{M_{\tilde{x}}} |\tilde{h} - h_0|^2$ instead of $\int_{M_{\tilde{x}}} |\tilde{h}|^2$, on one hand we are subtracting the “same infinite quantity” $\int_{M_{\tilde{x}}} |h_0|^2$ independent of \tilde{A} and, on the other hand, $\int_{M_{\tilde{x}}} |\tilde{h} - h_0|^2$ has no reason, any more, to be infinite (take $\tilde{h} = h_0$!).

Thus we are tempted to replace the classical Higg’s action by the following one

$$\tilde{S}_\lambda(\varphi, \tilde{A}) = \int_{M_{\tilde{x}}} |\nabla_{\tilde{A}} \varphi|^2 + |d\tilde{A} - h_0|^2 + \frac{\lambda}{2}(1 - |\varphi|^2)^2.$$

In view of formulating a well posed variational problem we introduce a finite covering U_α of $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ in each open set of which we can trivialize the bundle. The transition functions are choosen to be regular and we introduce the following set of minimisation

$$V = \left\{ \begin{array}{l} (\varphi_\alpha, \tilde{A}_\alpha) \in H_{loc}^1(U_\alpha, \mathbb{C}) \times H_{loc}^1(U_\alpha; \mathbb{R}^3) \\ s.t. \text{ on } U_\alpha \cap U_\beta \quad \varphi_\alpha = g_{\alpha\beta} \varphi_\beta \text{ and } i\tilde{A}_\alpha = g_{\alpha\beta}^{-1} dg_{\alpha\beta} + i\tilde{A}_\beta \\ \tilde{S}_\lambda(\varphi, \tilde{A}) < +\infty \end{array} \right\}$$

We have the following theorem which was already established in [8] and whose proof is essentially based on the choice of a good gauge which makes the functional coercive.

THEOREM 1.1.

$$\min_{(\varphi, \tilde{A}) \in V} \tilde{S}_\lambda(\varphi, \tilde{A}) \text{ is achieved.}$$

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After having given a proof of theorem 1.1, in the remains of the paper, we will follow the approach of [2] and [4], that is, we will make the parameter $\lambda = \frac{1}{\varepsilon}$ tend to infinity in view of forcing the density $|\varphi|$ to be close to 1.

The problem is that we cannot expect, for topological reasons, to find a section φ of the previous non trivial bundle E taking it’s values in the unit sphere of each fiber. This occurs only out of lines which connect the points x_i according to their multiplicity. The problem is similar to the one in dimension 2, here in dimension 3 φ necessarily has to have a degree around those lines which is independent of the local trivialisation and this will makes the energy tend to infinity. That is the main difficulty that one encounters when one tries to understand the behaviour of the minimizers as λ tends to infinity. The different quantities seem to diverge and one has to work

harder to get a norm for which something will remain bounded and converge in some space.

The notation $\tilde{S}_\lambda(\varphi, \tilde{A})$ is used in [8] where this variational problem was first introduced but, since we make the parameter λ tend to infinity, to be consistent with [4], the energy denoted by $\tilde{S}_\lambda(\varphi, \tilde{A})$ becomes $G_\varepsilon(\varphi, \tilde{A})$ where $\varepsilon = \frac{1}{\sqrt{\lambda}}$.

In fact, as we will see later, the scale ε will be much more relevant than λ .

Before we state our main result, we need to introduce the notion of minimal connection already used in [6]. Let \mathcal{S}_k be the set of all the permutations in $\{1, \dots, k\}$. The following quantity L is called *minimal connection* of the sequence $(x_i, d_i)_{i=1, \dots, n}$

$$L = \min_{\sigma \in \mathcal{S}_k} \sum_{i=1}^k |P_i - N_{\sigma(i)}|.$$

Consider any $\sigma \in \mathcal{S}_k$ which realizes the previous minimum (σ is perhaps not unique), we will call also *minimal connection* the 1 dimensional current \mathbb{L} with support in $\bigcup_{i=1}^k [P_i, N_{\sigma(i)}]$ oriented by $P_i - N_{\sigma(i)}$ and having the integer multiplicity

$$\theta(x) = \# \{i \text{ s.t. } x \in [P_i, N_{\sigma(i)}]\}.$$

L is of course the mass of \mathbb{L} .

Our main result is the following:

THEOREM 1.2. For any sequence ε_n tending to zero and for any sequence (φ_n, \tilde{A}_n) of minimizers of G_{ε_n} on V , one can extract a subsequence, still denoted (φ_n, \tilde{A}_n) , such that, for any $1 < p < 3/2$

$$h_n = d\tilde{A}_n - h_0 \text{ strongly converges in } H_{loc}^1(\mathbb{R}^3 \setminus \text{supp } \mathbb{L}) \cap W_{loc}^{1,p}(\mathbb{R}^3)$$

where \mathbb{L} is a minimal connection associated to the sequence $(x_i, d_i)_{i=1, \dots, n}$. Moreover the limit h_\star is the solution of

$$*(-\Delta h_\star + h_\star + h_0) = -2\pi \mathbb{L} \quad . \quad (1.5)$$

■

REMARK 1.1. If, for instance, the minimal connection is unique, thus the complete sequence converges.

REMARK 1.2. Contrary to the 2 dimensional case, we are not able to give a precise asymptotic expansion of the energy which should be of the form, when $d_i = \pm 1$,

$$G_{\varepsilon_n}(\varphi_n, \tilde{A}_n) = 2\pi L \log \frac{1}{\varepsilon_n} + W(x_i, d_i) + o(1) \quad .$$

Here $W(x_i, d_i)$ is a renormalised energy which represents the self interaction of the minimal connection on itself. But this has not the same importance as in the 2-dimensional problem considered in [3] and [4] since, here, we know where the singularities are going to be located and the renormalized energy $W(x_i, d_i)$ could be directly computed once we have fixed the x_i and d_i .

The asymptotic expansion of the minimal energy when all the minimal connections contain a segment with integer multiplicity different from 1 is an interesting question see remarks 2.1 and 2.2.

REMARK 1.3. The analysis developed in this paper could be adapted for solving non-gauge invariant problems in dimension 3, that is, variational problems for functionals of the form

$$E_\varepsilon(u) = \int_{\Omega \subset \mathbb{R}^3} |\nabla u|^2 + \frac{1}{\varepsilon^2}(1 - |u|^2)^2 \quad .$$

For instance, take Ω to be a convex domain of \mathbb{R}^3 and g_ε a boundary condition which simulates the vorticity at the boundary. A particular choice of g_ε could be a “ H^1 -approximation” of a map $g_\star : \partial\Omega \rightarrow S^1$ having some degree around different points of $\partial\Omega$ i.e.

- $|g_\varepsilon| \leq 1$ on $\partial\Omega$
- $\exists n$ distinct points (a_1, \dots, a_n) of $\partial\Omega$ and n integers (d_1, \dots, d_n) such that

$$\sum_{i=1 \dots n} d_i = 0$$

$$|g_\varepsilon| = 1 \text{ on } \partial\Omega \setminus (\bigcup_{i=1}^n B_\varepsilon(a_i) \cap \partial\Omega)$$

$$\text{deg}(g_\varepsilon, \partial B_\varepsilon(a_i) \cap \partial\Omega) = d_i$$
- an upper bound for the Dirichlet energy of g_ε

$$\int_{\partial\Omega} |\nabla_\tau g_\varepsilon|^2 \leq \left(2\pi \sum_{i=1}^n d_i^2 \right) \log \frac{1}{\varepsilon} + K.$$

One can prove that, given ε_n tending to zero, from a sequence u_{ε_n} of minimizers of E_{ε_n} among $H_{g_{\varepsilon_n}}^1(\Omega, \mathbb{C})$ one can extract a subsequence which converges in $H_{loc}^1(\bar{\Omega} \setminus \text{Supp} \mathbb{L})$ to an harmonic map u into S^1 , where \mathbb{L} is a minimal connection between the (a_i, d_i) and the degree of u around the different segments of \mathbb{L} corresponds to the multiplicity of the minimal connection along its segments. Moreover, if the minimal connection is unique, the complete sequence converges. The proof of this theorem can be established by following the main ideas of the proof of theorem 2. For instance one can transpose word by word the η -compactness lemma (part 3) to the non-gauge invariant situation just omitting to write h, A, A_0, h_0 ! The case where Ω is not convex is more complicated since it may happen that the lines of singularities (i.e. the bad set $|u_\varepsilon| \leq 1/2$) tend to touch $\partial\Omega$ as $\varepsilon \rightarrow 0$.

REMARK 1.4. With some more work, one could establish a stronger convergence than H^1 away from the line singularities as in the two dimensional case (see [2]).

REMARK 1.5. As in the two dimensional case, the description of the minimizers for $\varepsilon > 0$ is much more interesting and also more difficult than the one of the limiting situation (see in dimension 2 for $\varepsilon > 0$ [11], [1],...). In dimension 3, also, we can ask similar question to the ones in dimension 2: Are there several minimizers ? In the particular case of a dipole with multiplicity 1, is the minimizer axially symmetric around the dipole ? We can ask also questions specific to the 3 dimension: For instance when the multiplicity d of our dipole is no more 1 how does the zero set of $|\varphi|$ look like ? We can expect to get d different lines connecting the two poles in a way that has to be described. The configuration of those lines seems to be given by a

1-dimensional variational problem, similar to the zero dimensional problem in 2 dimensions given by the renormalised energy but depending strongly on ε . What could be the resulting interaction energy between those lines ?

The paper is devoted to the proof of theorem 1.2. The organisation of the proof is similar to the one developed in [4], but the dimension 3 induces a lot of new technical difficulties. It is for instance much more complicated to detect and to “catch” lines (which can be a-priori very bad) than a finite number of points.

One of the main differences comes also from the stress-energy tensor (see part 2-5) which contains, in dimension 3, a new term preventing a direct local estimate of the density and its closeness to 1 (see [4] part 3). In dimension 2 this closeness ensured quite directly the convergence except on a finite number of small balls. Here, following ideas and classical methods of the regularity theory for non-linear elliptic problems, we deduce, from the conservation law of the stress-energy tensor, some kind of monotonicity formula which is the main tool for proving what we call an “ η -compactness” lemma, (with reference to the classical ε -regularity lemma of the above mentioned theory). This η -compactness lemma says that, if the energy on a ball is sufficiently small compared to the global bound (in $\log \frac{1}{\varepsilon}$) established in 2-3, one can ensure that, in the ball of half radius, the density $|\varphi|$ is close to 1, this is necessary for being far away from the non-compactness locus (the bad part of $\mathbb{R}^3 \setminus \{x_i, \dots, x_n\}$). From the η -compactness Lemma 3.1 we deduce, in part 4, that, for any $\alpha < 1$ the bad set is contained in $\frac{c(\alpha)}{\varepsilon^\alpha}$ balls of radius ε^α and that, in some sense, we can bound the vorticity around this union of balls (this is Lemma 4.1). Having located the bad set in something which looks like a line when $\varepsilon \rightarrow 0$ ($\frac{c(\alpha)}{\varepsilon^\alpha}$ ball or radius ε^α), we finally establish a first bound, independent of ε , which is a bound of the $W^{1,p}$ norm of the magnetic field for any $1 < p < 3/2$ (see Lemma 5.1). This gives us a weak convergence, in $W^{1,p}$, of the magnetic field and Theorem 5.1 describes the limit. In part 6 we work out of the limiting bad set for establishing really the H^1 compactness, in fact we were not sure that, away from the limiting bad set, we could ensure the closeness to 1 for $|\varphi|$ which, combined with the weak $W^{1,p}$ convergence of the magnetic field, leads directly to the compactness. Here again the main ingredient is the η -compactness lemma. Finally, in part 7, we show that the limiting bad set which, for topological reasons, has to connect the different points x_i according to the multiplicities d_i , generates an energy, at the level ε , which is asymptotically greater than $2\pi \log \frac{1}{\varepsilon}$ times the connection it realizes between the (x_i, d_i) minus a positive quantity as small as we want. This fact combined with the initial lower bound established in 2.3 forces the limiting bad set to be a minimal connection between the (x_i, d_i) . The difficulty in this part is essentially technical, since, the limiting bad set, is a-priori only a 1-dimensional rectifiable current whose support is \mathcal{H}^1 -bounded.

In the appendix we prove some technical and basic Lemmas we use in the paper.

2. PRELIMINARIES AND NOTATIONS

2.1. CHOICE OF A COVERING AND A COULOMB GAUGE

First of all we fix a particular covering of $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ on which we will be able to trivialize E and that will be used in the remains of the paper in particular in 2.3 for establishing an upper bound for the energy. Suppose first that the indexing of the (N_1, \dots, N_k) is chosen such that the trivial permutation $(i \mapsto i)$ realizes a minimal connection i.e.

$$\sum_{i=1}^k |P_i - N_i| = \min_{\sigma \in S_k} \sum_{i=1}^k |P_i - N_{\sigma(i)}|.$$

Suppose now there exists i and i' such that $[P_i, N_i] \cap [P_{i'}, N_{i'}]$ contains more than 1 point.

1st case: $[P_i, N_i] = [P_{i'}, N_{i'}]$. We do not change the sequence $(P_i, \dots, P_k, N_i, \dots, N_k)$.

2nd case: $P_i = P_{i'}$ but $N_i \neq N_{i'}$ (or the opposite). Suppose $N_i \in [P_{i'}, N_{i'}]$ thus we repeat one more time N_i in the negative terms and we add N_i in the positive term thus $[P_i, N_i] \cup [P_{i'}, N_{i'}]$ is replaced by $[P_i, N_i] \cup [P_i, N_i] \cup [N_i, N_{i'}]$

3rd case: $[P_i, N_i] \subset [P_{i'}, N_{i'}]$ with $P_i \neq P_{i'}$ and $N_i \neq N_{i'}$.

In the sequence $(P_i, \dots, P_k, N_i, \dots, N_k)$ we repeat P_i and N_i such that

$$[P_i, N_i] \cup [P_{i'}, N_{i'}] = [P_{i'}, P_i] \cup [P_i, N_i] \cup [P_i, N_i] \cup [N_i, N_{i'}]$$

4th case: $[P_{i'}, N_{i'}] \subset [P_i, N_i]$ with $P_{i'} \neq P_i$ and $N_{i'} \neq N_i$

In the sequence $(P_1, \dots, P_k, N_1, \dots, N_k)$ we repeat $P_{i'}$ and N_i such that

$$[P_i, N_i] \cup [P_{i'}, N_{i'}] = [P_i, P_{i'}] \cup [P_{i'}, N_{i'}] \cup [P_{i'}, N_i] \cup [N_i, N_{i'}]$$

Thus in all the cases we can change the sequence $(P_i, N_i)_{i=1, \dots, k}$ and the number k without changing neither the bundle nor the minimal connection in such a way that for $i \neq i'$ $[P_i, N_i] \cap [P_{i'}, N_{i'}] = \emptyset$ or a point or $[P_i, N_i] = [P_{i'}, N_{i'}]$.

We proceed to a last reindexing by associating to a couple (P_i, N_i) the positive integer n_i which is the number of times the couple is repeated in $((P_i, N_i))_{i=1, \dots, k}$. The previous sequence is now represented by $((P_i, N_i, n_i))_{i=1, \dots, k'}$ with $k' \leq k$. k' will be denoted by k .

Let O_i be the middle point of the segment $[P_i, N_i]$, and $C_i(r)$ be the union of the two half cones of vertices P_i and N_i and of base the disk perpendicular to $[P_i, N_i]$, of centre O_i and radius r .

We can choose r sufficiently small such that $\forall i \neq i' \quad C_i(r) \cap C_{i'}(r) = \emptyset$ or $C_i(r) = C_{i'}(r)$.

Let $U_i = C_i(r)$ for $i \leq k$ and $U_{k+1} = \mathbb{R}^3 \setminus \bigcup_{i=1}^k \overline{C_i(\frac{r}{2})}$. $(U_i)_{i=1, \dots, k+1}$ realizes a covering of $\mathbb{R}^3 \setminus \{x_i, \dots, x_n\}$ on each open set U_i of which we can make a

trivialisations of E . By $d\theta_i$ we denote the angular 1-form around the axis (P_i, N_i) oriented by $P_i - N_i$.

We will only consider trivialisations of E over the U_i whose transition functions $g_{i,k+1}(x) \in S^1$ on $U_i \cap U_{k+1} = C_i(r) \setminus C_i(r/2)$ verify

$$\forall x \in C_i(r) \setminus C_i(r/2) \quad g_{i,k+1}^{-1} dg_{i,k+1}(x) = i n_i d\theta_i \tag{2. 1}$$

Suppose we have a trivialisations of E over the U_i whose transition functions verify (2. 1) we get all the others (verifying (2.1)) simply by multiplying by $e^{i\psi}$ where ψ is a global function on all of $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$. That is the change of trivialisations we will consider most of the time.

Consider the following 1-forms

$$\begin{aligned} A_{k+1}^0(x) &= \frac{1}{2} \sum_{i=1 \dots k} n_i \left(\frac{x-N_i}{|x-N_i|} - \frac{x-P_i}{|x-P_i|} \right) \cdot \frac{P_i-N_i}{|P_i-N_i|} d\theta_i \text{ in } U_{k+1} \\ A_j^0(x) &= \frac{1}{2} \sum_{i=1 \dots k} n_i \left(\frac{x-N_i}{|x-N_i|} - \frac{x-P_i}{|x-P_i|} \right) \cdot \frac{P_i-N_i}{|P_i-N_i|} d\theta_i \\ &\quad - n_j d\theta_j \text{ in } U_j, \quad j \leq k \end{aligned} \tag{2. 2}$$

Clearly $A_l^0(l = 1, \dots, k + 1)$ is regular in U_l and we have

$$\forall j \leq k \quad A_{k+1}^0(x) = A_j^0(x) + n_j d\theta_j \quad \text{in } U_{k+1} \cap U_j. \tag{2. 3}$$

Moreover one verifies that

$$\forall l \leq k + 1 \quad dA_l^0(x) = h_0(x) \quad \text{in } U_l \tag{2. 4}$$

Thus, suppose we have chosen a trivialisations of E over the U_i which verifies (2. 1), (2. 2) is the expression in this trivialisations of a connection whose curvature is h_0 .

Consider now a H^1 -connection \tilde{A} on E that is $\tilde{A}_j \in H_{loc}^1(U_j, \mathbb{R}^3)$ for $j \leq k + 1$ such that

$$\tilde{A}_{k+1}(x) = \tilde{A}_j(x) + n_j d\theta_j \quad \text{in } U_{k+1} \cap U_j. \tag{2. 5}$$

Let $\tilde{h} = d\tilde{A}$ in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ and suppose $h = \tilde{h} - h_0 \in L^2(\mathbb{R}^3)$. Denote by A the following global 1-form on \mathbb{R}^3

$$A(x) = -\frac{1}{4\pi} d^* \left(h * \frac{1}{|x|} \right) = -\frac{1}{4\pi} \sum_{k=1}^3 \int_{\mathbb{R}^3} \sum_{l=1}^3 h_{kl}(y) \frac{x^l - y^l}{|x - y|^3} dy dx_k \tag{2. 6}$$

Since $h \in L^2(\mathbb{R}^3)$, $dh \in W^{-1,2}(\mathbb{R}^3)$ and since $\text{supp}(dh) \subset \{x_i, \dots, x_n\}$ we have $dh \equiv 0$ in $\mathcal{D}'(\mathbb{R}^3)$. Thus we have

$$\begin{aligned} dA(x) &= dd^* \left(h * -\frac{1}{4\pi|x|} \right) + d^* d \left(h * -\frac{1}{4\pi|x|} \right) \\ &= \Delta \left(h * -\frac{1}{4\pi|x|} \right) = h(x) \quad \text{in } \mathcal{D}'(\mathbb{R}^3) \end{aligned} \tag{2. 7}$$

This implies

$$d \left(\tilde{A} - A^0 - A \right) = 0 \quad \text{in } \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$$

where $\tilde{A} - A^0$ is a global 1-form defined by

$$\tilde{A} - A^0(x) = \tilde{A}_l(x) - A_l^0(x) \text{ in } U_l \quad .$$

Since $H_{dR}^1(\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}) = 0$ there exists $\psi \in W_{loc}^{2,2}(\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}; \mathbb{R})$ such that

$$\tilde{A} - A^0 - A(x) = d\psi(x) \text{ in } \mathbb{R}^3 \setminus \{x_1, \dots, x_n\} \quad .$$

Thus, if we make the global change of trivialisation by multiplying the fiber over x by $e^{-i\psi(x)}$ the new expression of the connection \tilde{A} is given by

$$\tilde{A}'_l(x) = \tilde{A}_l - d\psi(x) = A_l^0(x) + A(x) \quad \text{in } U_l. \quad (2. 8)$$

Consider now the set among which we are going to minimize G_ε ,

$$V = \left\{ \begin{array}{l} (\varphi_l, \tilde{A}_l) \in H_{loc}^1(U_l, \mathbb{C}) \times H_{loc}^1(U_l, \mathbb{R}^3) \text{ s.t.} \\ \forall j \leq k \quad \varphi_{k+1} = \varphi_j e^{in_j \theta_j}, \quad \tilde{A}_{k+1} = \tilde{A}_j + n_j d\theta_j \text{ in } U_j \cap U_{k+1} \\ G_\varepsilon(\varphi, \tilde{A}) < +\infty \end{array} \right\} \quad (2. 9)$$

where $e^{i\theta_j}$ is the angular coordinate around the axis $[P_j, N_j]$ in $C_j(r) \setminus C_j(\frac{r}{2})$ where we have fixed some horizontal direction equal to 1.

Among V we say that (φ, \tilde{A}) is equivalent to (φ', \tilde{A}') if there exists $\psi \in W_{loc}^{2,2}(\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}, \mathbb{R})$ such that

$$\forall 1 \leq l \leq k + 1 \quad \tilde{A}_l = \tilde{A}'_l + d\psi \quad \text{and} \quad \varphi_l = e^{i\psi} \varphi'_l \quad \text{in } U_l \quad (2. 10)$$

PROPOSITION 2.1. Let $(\varphi, \tilde{A}) \in V$, there exists $(\varphi', \tilde{A}') \in V$ equivalent to (φ, \tilde{A}) such that

$$\forall 1 \leq l \leq k + 1 \quad \tilde{A}'_l(x) = A_l^0(x) + A(x) \quad \text{in } U_l$$

where $A(x) = -\frac{1}{4\pi} d^* \left((\tilde{h} - h_0) * \frac{1}{|x|} \right)$: This is what we call the Coulomb gauge of \tilde{h} . ■

The proposition is proved just above, the name Coulomb gauge comes from the fact that

$$d^* A = 0 \quad \text{in } \mathcal{D}'(\mathbb{R}^3). \quad (2. 11)$$

Moreover one easily verifies that

$$\forall l \in \{1, \dots, k\} \quad d^* A_l^0 = 0 \quad \text{in } \mathcal{D}'(U_l) \quad . \quad (2. 12)$$

Thus

$$\forall l \in \{1, \dots, k\} \quad d^* \tilde{A}'_l = 0 \quad \text{in } \mathcal{D}'(U_l) \quad . \quad (2. 13)$$

2.2. EXISTENCE OF MINIMIZERS. PROOF OF THEOREM 1.1

In this part we prove the theorem 1 stated in the introduction. More precisely we prove the following,

THEOREM 2.3. For any $\varepsilon > 0$ the minimum of G_ε among V (defined by (2. 9)) is achieved by, at least, one class of couples $(\varphi_\varepsilon, \tilde{A}_\varepsilon)$ (for the relation 2. 10). The $(\varphi_\varepsilon, \tilde{A}_\varepsilon)$ verify

$$-\frac{1}{2} \Delta |\varphi_\varepsilon|^2 = \frac{1}{\varepsilon^2} |\varphi_\varepsilon|^2 (1 - |\varphi_\varepsilon|^2) - |\nabla_{\tilde{A}_\varepsilon} \varphi_\varepsilon|^2 \text{ in } \mathcal{D}'(\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}), \quad (2. 14)$$

and

$$-d^*h_\varepsilon = J_\varepsilon \quad \text{in } \mathcal{D}'(\mathbb{R}^3) \quad , \tag{2. 15}$$

where J_ε , the current, is the gauge invariant quantity equal to $(i\varphi_\varepsilon, d_{\tilde{A}_\varepsilon}\varphi_\varepsilon)$ in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$. Moreover $|\varphi_\varepsilon|^2, J_\varepsilon, h_\varepsilon$ are regular in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$. ■

Proof of theorem 2.3

Let ε be a fixed positive real for the remain of the proof. Let (φ^p, \tilde{A}^p) be a minimizing sequence of G_ε . We have

$$\int_{\mathbb{R}^3} |h^p|^2 = \int_{\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}} |\tilde{h}^p - h_0|^2 < C \quad \text{indep. of } p \quad . \tag{2. 16}$$

Thus by classical estimates for Calderon - Zygmund operators we have

$$\left(\int_{\mathbb{R}^3} |A^p|^6 \right)^{1/6} + \left(\int_{\mathbb{R}^3} |\nabla A^p|^2 \right)^{1/2} \leq C \left(\int_{\mathbb{R}^3} |h^p|^2 \right)^{1/2} \leq C \quad , \tag{2. 17}$$

where A^p is defined by (2. 6).

In the class of (φ, \tilde{A}) consider the Coulomb gauge given by Proposition II.1. For any $l \leq k + 1$ and K compact set in U_l , we have

$$\int_{U_l \cap K} |\nabla_{\tilde{A}^p} \varphi^p|^2 \leq C \quad ,$$

thus

$$\begin{aligned} \int_{U_l \cap K} |\nabla \varphi_l^p|^2 &\leq C + 2 \int_{U_l \cap K} |\tilde{A}_l^p|^2 |\varphi^p|^2 \\ &\leq C + 2 \int_{U_l \cap K} |\tilde{A}_l^p|^2 + \\ &\quad + 2 \left(\int_{U_l \cap K} |\tilde{A}_l^p|^4 \right)^{1/2} \left(\int_{U_l \cap K} (1 - |\varphi^p|^2)^2 \right)^{1/2} \end{aligned} \tag{2. 18}$$

Since A_l^0 is regular on U_l , $\int_{U_l \cap K} |A_l^0|^4$ is bounded and from (2.15) we deduce that

$$\int_{U_l \cap K} |A_l^0 + A^p|^4 \leq C \quad \text{indep. of } p \quad . \tag{2. 19}$$

Thus from (2. 18) and (2. 19) it follows that φ_l^p weakly converges in $W_{loc}^{1,2}(U_l, \mathbb{C})$ (up to a subsequence) and this is also the case for $\tilde{A}_l^p = A_l^0 + A^p$. By lower semi continuity of the $W^{1,2}$ it is clear that the limit $(\varphi_{l,\varepsilon}, A_l^0 + A_\varepsilon^p)_{l=1 \dots k+1}$ minimizes G_ε .

The fact that equations (2. 14) and (2. 15) are verified in $\mathcal{D}'(\mathbb{R}^3 \setminus \{x_1, \dots, x_n\})$ and that, moreover, $|\varphi_\varepsilon|^2, J_\varepsilon$ and h_ε are regular in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ is proved in [10]. Let us just prove that (2. 15) is verified in $\mathcal{D}'(\mathbb{R}^3)$.

Let Ω be a bounded domain of \mathbb{R}^3 , since $h_\varepsilon \in L^2(\Omega)$, $d^*h_\varepsilon \in W^{-1,2}(\Omega)$, moreover, since $\int_{\mathbb{R}^3} (1 - |\varphi_\varepsilon|^2)^2 < +\infty$, $|\varphi_\varepsilon| \in L^4(\Omega)$ and this implies that

the function J_ε on \mathbb{R}^3 defined by $J_\varepsilon \equiv (i\varphi_\varepsilon, d_{A_\varepsilon}\varphi_\varepsilon)$ in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ is in $L^{4/3}(\Omega) \subset W^{-1,2}(\Omega)$.

Thus

$$d^*h_\varepsilon + J_\varepsilon \in W^{-1,2}(\Omega) \quad ,$$

and $\text{supp}(d^*h_\varepsilon + J_\varepsilon) \subset \{x_1, \dots, x_n\}$, so this implies (2. 15). ■

2.3. AN UPPER BOUND FOR THE ENERGY

In this section we prove the following lemma

LEMMA 2.1. For any $\alpha > 0$ we have

$$\min_V G_\varepsilon \leq 2\pi(L + \alpha)Ln\frac{1}{\varepsilon} + K(\alpha) \quad , \tag{2. 20}$$

where L is the minimal connection between the (x_i, d_i) and K only depends on α , not on ε . ■

REMARK 2.6. Contrary to the 2-dim. case, the upper bound for the energy requires the introduction of an external parameter α . This is not because the dimension 3 leads to new technical difficulties, this weaker upper-bound, Compare to the two dimensional one (see proposition 2.5 [4]), is intrinsically linked to the dimension 3 for any minimal connection between the (x_i, d_i) containing a segment with integer multiplicity different from 1.

Consider for instance the case of a dipole (P, N) of charge $d > 1$ ie $((x_i, d_i) = ((P, d), (N, -d))$. The minimal connection is the segment $[P, N]$ with the integer multiplicity d , that is $\mathbb{L} = d[[P, N]]$ and $L = d|P - N|$. The singular set tends to concentrate along this segment when $\varepsilon \mapsto 0$. Suppose this configuration is adopted for any $\varepsilon > 0$. Then the energy would be greater than $2\pi d^2|P - N| \log \frac{1}{\varepsilon}$! Thus at the stage ε , what we call the bad set prefer to decompose d distinct tubes between P and N which are close to the segment $[P, N]$ (but not too much so that the interaction energy between them is at most of order $\approx \log \log \varepsilon$).

REMARK 2.7. If (x_i, d_i) admits a minimal connection which contains no segment of integer multiplicity different from 1, then Lemma 2.1 holds for $\alpha = 0$. It seems that this is a minimal connection that the singular set would prefer to adopt.

Proof of Lemma 2.1.

Let $\alpha > 0$ and $\varepsilon > 0$. We construct (φ_l, \tilde{A}_l) in V verifying (2. 20). First of all we construct $|\varphi|$.

For $1 \leq i \leq k$ consider n_i the multiplicity of (P_i, N_i) in the minimal connection that we consider and for $1 \leq s \leq n_i$ introduce $L_{i,s}$ to be the following segments.

If $n_i = 1$, take $L_{i,s} = L_{i,1} = [P_i, N_i]$. and if $n_i \neq 1$ we identify the perpendicular plane to (P_i, N_i) passing by O_i , the center of $[P_i, N_i]$, with \mathbb{C} and denote by $(a_{i,s})_{s=1, \dots, n_i}$ the n_i -th roots of 1 multiplied by α . (we suppose $\alpha < \frac{r}{4}$ in such a way that $a_{i,s} \in C_i(\frac{r}{2})$) and also for technical reasons which

will be clear later). Take $L_{i,s} = [P_i, a_{i,s}] \cup [a_{i,s}, N_i]$.

Consider \tilde{L} the following union of segments

$$\tilde{L} = \cup_{i=1}^k \cup_{s=1}^{n_i} L_{i,s}$$

Let χ be the following cut-off function on \mathbb{R}_+

$$\chi(x) = x \text{ for } x \leq 1 \quad \text{and} \quad \chi(x) = 1 \text{ for } x \geq 1.$$

Define $|\varphi|(x) = \chi\left(\frac{\text{dist}(x, \tilde{L})}{\varepsilon}\right)$. We have

$$\int_{\mathbb{R}^3} |\nabla|\varphi||^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 \leq \int_{\text{supp}(1-|\varphi|)} \frac{C}{\varepsilon^2} \leq C \quad , \quad (2.21)$$

where C does not depend on ε .

We choose

$$\forall l \leq k + 1 \quad \tilde{A}_l = A_l^0 \quad \text{in } U_l \quad . \quad (2.22)$$

We claim that

$$\sum_{l=1 \dots k+1} \int_{U_l} |A_l^0|^2 < +\infty \quad . \quad (2.23)$$

It suffices to prove that, if one considers C_1 the union of the two cones of vertices $p = (0, 0 + 1)$ and $N = (0, 0, -1)$ and of base the horizontal unit disk of center 0 and if we take

$$A_1^0 = \left[\frac{1}{2} \left(\frac{1-z}{|X-P|} + \frac{z+1}{|X-N|} \right) - 1 \right] \frac{xdy - ydx}{(x^2 + y^2)} \quad \text{in } C_1$$

$$A_2^0 = \left[\frac{1}{2} \left(\frac{1-z}{|X-P|} + \frac{z+1}{|X-N|} \right) \right] \frac{xdy - ydx}{x^2 + y^2} \quad \text{in } \mathbb{R}^3 \setminus C_1 \quad ,$$

where $X = (x, y, z)$. Then we have

$$\int_{C_1} |A_1^0|^2 + \int_{\mathbb{R}^3 \setminus C_1} |A_2^0|^2 < +\infty \quad .$$

This is verified by a direct computation.

Let $l \leq k$, consider an orthonormal basis $(\vec{i}_l, \vec{j}_l, \vec{k}_l)$ such that $\vec{k}_l = \frac{P_l - N_l}{|P_l - N_l|}$. Each horizontal plane is identified with \mathbb{C} .

If $n_l = 1$, take

$$\varphi_l(x) = |\varphi|(x)e^{-i\theta_l} \quad \text{in } U_l \quad (2.24)$$

where θ_l is the polar horizontal angle of x in cylindrical coordinates associated with $(\vec{i}_l, \vec{j}_l, \vec{k}_l)$.

If $n_l > 1$. For $0 \leq z_0 \leq |P_l - O_l| = D_l$ (resp. $0 \geq z_0 \geq -|P_l - O_l|$), denote by $a_{l,s}(z_0)$ the point $a_{l,s}(z_0) = \{z = z_0\} \cap [P_l, a_{l,s}]$ (resp. $a_{l,s}(z_0) = \{z = z_0\} \cap [a_{l,s}, N_l]$). Take

$$\varphi_l(x, y, z) = |\varphi|(x, y, z) \prod_{s=1}^{n_l} \left(\frac{\xi - \lambda \tilde{a}_{l,s}(z)}{|\xi - \lambda \tilde{a}_{l,s}(z)|} \right)^{-1} \quad \text{in } U_l \quad , \quad (2.25)$$

where $\xi = x + iy$, $\tilde{a}_{l,s}(z)$ is the projection on $\{z = 0\} \simeq \mathbb{C}$ of $a_{l,s}(z)$ and λ is an interpolation function between 0 and 1 verifying $\lambda \equiv 1$ in $C_l(\frac{r}{4})$, $\lambda \equiv 0$ in $C_l(\frac{r}{2})$ and $\lambda = 0$ outside. Precisely we take

$$\text{for } 0 \leq z \leq D_l \quad \lambda(x, y, z) = \lambda(|\xi|, z) = 2 - \frac{4D_l}{D_l - z} \frac{|\xi|}{r} \quad ,$$

and

$$\lambda(|\xi|, -z) = \lambda(|\xi|, z) \quad .$$

From (2. 24) and (2. 25) we deduce that

$$\varphi_{k+1} = |\varphi| \quad \text{in } U_{k+1} \quad .$$

If $n_l = 1$, we have from (2. 24)

$$\begin{aligned} \int_{U_l} |\nabla \varphi_l|^2 &= \int_{U_l} |\nabla |\varphi||^2 + \int_{U_l} |\varphi|^2 \frac{1}{\rho^2} \rho d\rho d\theta dz \\ &\leq C + 2\pi \times 2 \int_0^{D_l - \varepsilon} \int_\varepsilon^{r \frac{D_l - z}{b_l}} \frac{1}{\rho} d\rho dz \\ &\leq C + 2\pi(2D_l) \log \frac{1}{\varepsilon} + 4\pi \int_0^{D_l - \varepsilon} \log \left[\left(1 - \frac{z}{D_l}\right) r \right] dz \\ &\leq 2\pi |P_l - N_l| \log \frac{1}{\varepsilon} + C \quad , \end{aligned} \tag{2. 26}$$

where C does not depend on ε .

If $n_l > 1$ we write

$$\int_{U_l} |\nabla \varphi_l|^2 = \int_{C_l(\frac{r}{4})} |\nabla \varphi_l|^2 + \int_{U_l \setminus C_l(\frac{r}{4})} |\nabla \varphi_l|^2 \quad . \tag{2. 27}$$

One verifies that the second integral on the right hand side of (2. 27) can be bounded independently of ε . Indeed, in $U_l \setminus (C_l(\frac{r}{4}) \cup B_\varepsilon(P_l) \cup B_\varepsilon(N_l))$ $|\varphi| = 1$ and

$$\begin{aligned} \left| \nabla \prod_{s=1}^{n_l} \left(\frac{\xi - \lambda(|\xi|, z) \tilde{a}_{l,s}(z)}{|\xi - \lambda(|\xi|, z) \tilde{a}_{l,s}(z)|} \right) \right| &\leq C \sum_{s=1}^{n_l} \frac{1}{|\xi - \lambda \tilde{a}_{l,s}(z)|} \\ &\leq C \frac{n_l}{\frac{r}{4} - \alpha} . \end{aligned}$$

Recall that the α that we consider are strictly less than $r/4$. The first integral of (2. 27) can be bounded in the following way

$$\begin{aligned} \int_{C_l(\frac{r}{4})} |\nabla \varphi_l|^2 &\leq \int_{C_l(\frac{r}{4})} |\nabla |\varphi_l||^2 + C \\ &+ 2 \int_0^{D_l} dz \int_{B_2(z) \setminus UB(a_{l,s}(z), \varepsilon)} \left| \nabla \prod_{s=1}^{n_l} \left(\frac{|\xi - \tilde{a}_{l,s}(z)|}{|\xi - \tilde{a}_{l,s}(z)|} \right) \right|^2 \end{aligned} \tag{2. 28}$$

where $B_2(z)$ denotes the horizontal disk of radius $r \frac{D_l - z}{D_l}$ and center the point $(0, 0, z)$ in $(\vec{i}_l, \vec{j}_l, \vec{k}_l)$.

We have

$$\left| \frac{\partial}{\partial z} \prod_{s=1}^{n_l} \left(\frac{\xi - \tilde{a}_{l,s}(z)}{|\xi - \tilde{a}_{l,s}(z)|} \right) \right| \leq \sum_{s=1}^{n_l} \left| \frac{\partial \tilde{a}_{l,s}}{\partial z} \right| \times \frac{C}{|\xi - \tilde{a}_{l,s}(z)|} \tag{2. 29}$$

but

$$\left| \frac{\partial \tilde{a}_{l,s}}{\partial z} \right| \leq C\alpha \text{ and } \int_{B_2(z) \setminus B(a_{l,s}(z), \varepsilon)} \frac{1}{|\xi - \tilde{a}_{l,s}(z)|^2} \leq C \log \frac{1}{\varepsilon}$$

thus we have

$$\int_0^{D_l} dz \int_{B_2(z) \setminus B(a_{l,s}(z), \varepsilon)} \left| \frac{\partial}{\partial z} \prod_{s=1}^{n_l} \left(\frac{\xi - \tilde{a}_{l,s}(z)}{|\xi - \tilde{a}_{l,s}(z)|} \right) \right|^2 \leq C\alpha^2 \log \frac{1}{\varepsilon} \tag{2. 30}$$

Concerning the horizontal derivatives in the second integral of the right hand side of (2. 28) we use an upper bound given by (4.38) in [7]. Finally one verifies easily that the energy of $|\varphi_l|$ is uniformly bounded, thus we deduce

$$\int_{C_l(\frac{r}{4})} |\nabla \varphi_l|^2 \leq 2\pi n_l (|P_l - N_l| + c\alpha^2) \log \frac{1}{\varepsilon} + K(\alpha) \tag{2. 31}$$

Combining (2. 21), (2. 23), (2. 26) and (2. 31) we get the desired result.

2.4. PRELIMINARY ESTIMATES

We establish, first, the L^∞ bound for $|\varphi_\varepsilon|$.

LEMMA 2.2. Let $(\varphi_\varepsilon, \tilde{A}_\varepsilon) \in V$ be a minimizer of G_ε then we have

$$|\varphi_\varepsilon| \leq 1 \quad \text{in } \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}. \tag{2. 32}$$

■

Proof of Lemma 2.2.

Suppose $\omega = \{x \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\} / |\varphi_\varepsilon|(x) > 1\}$ is not empty. Define $\tilde{\varphi}_\varepsilon$ in the following way

$$\tilde{\varphi}_\varepsilon = \varphi_\varepsilon \quad \text{if } |\varphi_\varepsilon| \leq 1$$

and

$$\tilde{\varphi}_\varepsilon = \varphi_\varepsilon / |\varphi_\varepsilon| \quad \text{in } \omega \tag{2. 32}$$

It is clear that

$$\int_{\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}} (1 - |\tilde{\varphi}_\varepsilon|^2)^2 < \int_{\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}} (1 - |\varphi_\varepsilon|^2)^2 \tag{2. 33}$$

Moreover since

$$\forall x \in \omega \quad |\nabla_{\tilde{A}_\varepsilon} \tilde{\varphi}_\varepsilon|^2(x) = \left| \nabla_{\tilde{A}_\varepsilon} \frac{\varphi_\varepsilon}{|\varphi_\varepsilon|} \right|^2,$$

and

$$\forall x \in \omega \quad |\nabla_{\tilde{A}_\varepsilon} \varphi_\varepsilon|^2(x) = |\nabla |\varphi_\varepsilon||^2(x) + \left| \nabla_{\tilde{A}_\varepsilon} \frac{\varphi_\varepsilon}{|\varphi_\varepsilon|} \right|^2,$$

we have

$$\forall x \in \omega \quad |\nabla_{\tilde{A}_\varepsilon} \tilde{\varphi}_\varepsilon|(x) \leq |\nabla_{\tilde{A}_\varepsilon} \varphi_\varepsilon|(x) \quad . \quad (2. 34)$$

The previous inequality combined with (2. 33) imply $G_\varepsilon(\tilde{\varphi}_\varepsilon, \tilde{A}_\varepsilon) < G_\varepsilon(\varphi_\varepsilon, \tilde{A}_\varepsilon)$, contradiction. ■

We establish, now, preliminary estimates for A_ε and h_ε

LEMMA 2.3. Let $(\varphi_\varepsilon, \tilde{A}_\varepsilon) \in V$ be a minimizer of $G_\varepsilon, h_\varepsilon = d\tilde{A}_\varepsilon - h^0$ and A_ε defined by (2. 6), we have for ε sufficiently small

$$\left(\int_{\mathbb{R}^3} |A_\varepsilon|^6 \right)^{1/3} + \int_{\mathbb{R}^3} |\nabla A_\varepsilon|^2 \leq C \int_{\mathbb{R}^3} h_\varepsilon^2 \leq C \log \frac{1}{\varepsilon} \quad , \quad (2. 35)$$

$$\int_{\mathbb{R}^3} |\nabla h_\varepsilon|^2 \leq \int_{\mathbb{R}^3} |\nabla^2 A_\varepsilon|^2 \leq C \log \frac{1}{\varepsilon}, \quad (2. 36)$$

$$\forall r \leq 1 \quad \forall x \in \mathbb{R}^3 \quad \frac{1}{r^2} \int_{B_r(x)} h_\varepsilon^2 \leq C \log \frac{1}{\varepsilon}, \quad (2. 37)$$

$$\forall r \leq 1 \quad \forall x \in \mathbb{R}^3 \quad \frac{1}{r} \int_{\partial B_r(x)} h_\varepsilon^2 \leq C \log \frac{1}{\varepsilon} \quad (2. 38)$$

and

$$\forall x, y \in \mathbb{R}^3 \text{ s.t. } |x - y| \leq 1, \quad (2. 39)$$

$$|A_\varepsilon(x) - A_\varepsilon(y)| \leq C \left(\log \frac{1}{\varepsilon} \right)^{1/2} |x - y|^{1/2}.$$
■

Proof of Lemma 2.3.

(2. 35) comes from (2. 20) and classical estimates on Calderon-Zygmund operators.

From (2. 15) we know that

$$-d^* h_\varepsilon = J_\varepsilon \quad \text{in } \mathcal{D}'(\mathbb{R}^3)$$

In the other hand, by definition, we have

$$A_\varepsilon = -\frac{1}{4\pi} \frac{1}{|x|} * d^* h_\varepsilon = +\frac{1}{4\pi} \frac{1}{|x|} * J_\varepsilon$$

From (2. 20) and (2. 32) we have $J_\varepsilon \in L^2(\mathbb{R}^3)$ and

$$\int_{\mathbb{R}^3} |J_\varepsilon|^2 = \int_{\mathbb{R}^3} |(i\varphi, \nabla_{\tilde{A}} \varphi)|^2 \leq C \log \frac{1}{\varepsilon}. \quad (2. 40)$$

In view of (2. 40), we get, by classical estimates on the Riesz potential

$$\int_{\mathbb{R}^3} |\nabla^2 A_\varepsilon|^2 \leq C \int_{\mathbb{R}^3} J_\varepsilon^2 \leq C \log \frac{1}{\varepsilon}.$$

This proves (2. 36).

From (2. 36) we deduce $\|h_\varepsilon\|_{W^{1,2}(B_1(x))}^2 \leq C \log \frac{1}{\varepsilon}$ for any $x \in \mathbb{R}^3$. Thus by Sobolev injection we have

$$\left(\int_{B_1(x)} h_\varepsilon^6 \right)^{1/6} \leq C \left(\log \frac{1}{\varepsilon} \right)^{1/2} .$$

Hölder’s inequality on $B_r(x)$ leads to

$$\int_{B_r(x)} h^2 \leq |B_r(x)|^{\frac{2}{3}} \left(\int_{B_r(x)} h_\varepsilon^6 \right)^{1/3} \leq Cr^2 \log \frac{1}{\varepsilon} .$$

This is (2. 37). Since $\int_{B_r(x)} |\nabla h_\varepsilon|^2 + \frac{1}{r^2} \int_{B_r(x)} |h_\varepsilon^2| \leq C \log \frac{1}{\varepsilon}$, taking the trace of h on $\partial B_r(x)$ we get (2. 38).

Finally (2. 39) is a consequence of (2. 35), (2. 36) , which both imply $\|A_\varepsilon\|_{W^{2,2}(B_1(x))} \leq C (\log \frac{1}{\varepsilon})^{1/2}$ and Sobolev injections. ■

The last preliminary estimate is the L^∞ bound (in $\frac{1}{\varepsilon}$) of $|\nabla_{A_\varepsilon} \varphi_\varepsilon|$. This estimate with (2. 36) and also (2. 32) are the only preliminary estimates which are optimal. We have the following Lemma

LEMMA 2.4. There exists C independent of ε such that for any x verifying $\text{dist}(x, \{x_1, \dots, x_n\}) \geq C\varepsilon$ we have

$$|\nabla_{\tilde{A}_\varepsilon} \varphi_\varepsilon|(x) \leq \frac{C}{\varepsilon} . \tag{2. 41}$$
■

Proof of Lemma 2.4.

Let $x \in \mathbb{R}^3 \setminus \cup_{i=1}^n B(x_i, 2\varepsilon)$.

There exists $\beta > 0$ depending only on r , the diameter of the cones $C_i(r) (i \leq k)$, and $l \leq k$ such that $B(x, \beta\varepsilon) \subset U_l \cap \mathbb{R}^3 \setminus \cup_{i=1}^n B(x_i, \varepsilon)$. We will only work in $B(x, \beta\varepsilon)$. Take the Coulomb gauge in U_l given by (2. 8) and change a little bit the trivialisation on $B(x, \beta\varepsilon)$ such that the expression of \tilde{A} over $B(x, \beta\varepsilon)$ becomes

$$\forall y \in B(x, \beta\varepsilon) \quad \tilde{A}_l(y) = A_l^0(y) + A(y) - A(x) . \tag{2. 42}$$

We just have to multiply all the fibers over $B(x, \beta\varepsilon)$ by $e^{-iA(x)y}$. Note that (2. 2) implies

$$\forall y \in B(x, \beta\varepsilon) \quad |A_l^0|(y) \leq \frac{C}{\varepsilon} . \tag{2. 43}$$

Combining (2. 39), (2. 42) and (2. 43) we have

$$\forall y \in B(x, \beta\varepsilon) \quad |\tilde{A}_l|(y) \leq \frac{C}{\varepsilon} . \tag{2. 44}$$

φ_l verifies the following equation in $B(x, \beta\varepsilon)$

$$-\Delta \varphi_l = \frac{1}{\varepsilon^2} \varphi_l (1 - |\varphi_l|^2) - \tilde{A}_l^2 \varphi_l - 2i \tilde{A}_l \cdot \nabla \varphi_l .$$

Consider a dilation of rate ε and denote

$$\bar{\varphi}_l(y) = \varphi_l(\varepsilon y) \quad \text{and} \quad \bar{A}_l(y) = \varepsilon \tilde{A}_l(\varepsilon y) \quad .$$

We have

$$-\Delta \bar{\varphi}_l = \bar{\varphi}_l(1 - |\bar{\varphi}_l|^2) - \bar{A}_l^2 \bar{\varphi}_l - 2i\bar{A}_l \cdot \nabla \bar{\varphi}_l \quad \text{in } B(x, \beta) \quad . \quad (2. 45)$$

Let $\bar{f}_l = \bar{\varphi}_l(1 - |\bar{\varphi}_l|^2) - \bar{A}_l^2 \bar{\varphi}_l$, from (2. 44) we have

$$\forall y \in B(x, \beta) \quad |\bar{f}_l(y)| \leq C \quad . \quad (2. 46)$$

Since $d^* \bar{A}_l = 0$ on $B(x, \beta)$, we write $i\bar{A}_l \cdot \nabla \bar{\varphi}_l$ in the following form

$$2i\bar{A}_l \cdot \nabla \bar{\varphi}_l = d^* (2i\bar{\varphi}_l \bar{A}_l) = \sum_{k=1}^3 \frac{\partial}{\partial k} (i\bar{\varphi}_l \bar{A}_l^k) \quad .$$

Let U and V be the solutions of the following equations

$$\begin{cases} -\Delta U = 0 & \text{in } B(x, \beta) \\ U = \bar{\varphi}_l & \text{on } \partial B(x, \beta) \quad , \end{cases} \quad (2. 47)$$

and

$$\begin{cases} -\Delta V = \bar{f}_l - 2d^*(i\bar{\varphi}_l \bar{A}_l) & \text{in } B(x, \beta) \\ V = 0 & \text{on } \partial B(x, \beta) \quad . \end{cases} \quad (2. 48)$$

We clearly have $\bar{\varphi}_l = U + V$ in $B(x, \beta)$. Since $\|\bar{\varphi}_l\|_{L^\infty} \leq 1$, from elliptic estimates we have,

$$|\nabla U|(y) \leq C \quad \text{in } B(x, \beta/2) \quad . \quad (2. 49)$$

Moreover, since $\varphi_l \bar{A}_l$ is bounded in $B(x, \beta)$, $d^*(i\varphi_l \bar{A}_l) \in W^{-1,q}(B(x, \beta))$ for all $q > 1$ and we have

$$\forall 1 < p < +\infty \quad \int_{B(x, \beta)} |\nabla V|^p \leq C_p \quad . \quad (2. 50)$$

Combining (2. 49) and (2. 50) we get

$$\forall 1 < p < +\infty \quad \int_{B(x, \beta/2)} |\nabla \bar{\varphi}_l|^p < C_p \quad . \quad (2. 51)$$

We consider now U_1 and V_1 the solutions of the following equations

$$\begin{cases} -\Delta U_1 = 0 & \text{in } B(x, \beta/2) \\ U_1 = \bar{\varphi}_l & \text{on } \partial B(x, \beta/2) \quad , \end{cases} \quad (2. 52)$$

and

$$\begin{cases} -\Delta V_1 = \bar{f}_l - 2i\bar{A}_l \cdot \nabla \bar{\varphi}_l & \text{in } B(x, \beta/2) \\ V_1 = 0 & \text{on } \partial B(x, \beta/2) \quad . \end{cases} \quad (2. 53)$$

Clearly we have $\bar{\varphi}_l = U_1 + V_1$ and since the right hand side of (2. 53) is bounded in $L^p \quad \forall p < +\infty$ we get

$$\forall 1 < p < +\infty \quad \|\bar{\varphi}_l\|_{W^{2,p}(B(x, \beta/4))} \leq C_p \quad ,$$

and this leads $|\nabla \bar{\varphi}_l(x)| < C$ by Sobolev embeddings.

Thus we have

$$|\nabla\bar{\varphi}_l(x) - i\tilde{A}_l\bar{\varphi}_l(x)| < C \quad \text{indep. of } \varepsilon \quad ,$$

which is the desired result. ■

2.5. THE MAXWELL STRESS-ENERGY TENSOR AND THE ALMOST MONOTONICITY FORMULA

Because of the renormalisation, by subtracting h_0^2 in the energy, the conservation law of the Maxwell stress tensor given in [10] (formula (2.1), chapter 2.2) is modified by a term depending on h_0 which could be bad for us when we are close to the x_i (that could make the monotonicity formula blow-up). That is the reason why we give explicitly the computation of this conservation law.

Consider $x_0 \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$. There exists $\rho > 0$ and $l \leq k + 1$ such that $B_\rho(x_0) \subset U_l$.

Let (φ, \tilde{A}) be a minimizer of G_ε (we omit here the subscript). Consider $\xi \in C_0^\infty(B_\rho(x_0), \mathbb{R}^3)$ and the following perturbation $(\varphi_t, \tilde{A}_t) \in V$ of (φ, \tilde{A})

$$\forall i \leq k + 1 \quad \begin{cases} \varphi_i^t = \varphi_i & \text{in } U_i \cap \mathbb{R}^3 \setminus B_\rho(x_0) \\ \tilde{A}_i^t = \tilde{A}_i & \text{in } U_i \cap \mathbb{R}^3 \setminus B_\rho(x_0) \end{cases}$$

and

$$\begin{cases} \varphi_l^t(x) = \varphi_l(x + t\xi) & \text{in } B_\rho(x_0) \\ A_l^t(x) = A_l(x + t\xi) & \text{in } B_\rho(x_0). \end{cases}$$

We do not write the index l anymore since all the perturbation is in U_l .

We have

$$\nabla_{\tilde{A}^t} \varphi^t(x) = \nabla_{\tilde{A}} \varphi(x + t\xi) + t \nabla \varphi(x + t\xi) \cdot \nabla \xi(x) \quad \text{in } B_\rho(x_0) \quad , \quad (2.54)$$

and

$$\begin{aligned} h_{kl}^t &= \tilde{h}_{kl}^t - h_{kl}^0 = \frac{\partial \tilde{A}_l^t}{\partial k} - \frac{\partial \tilde{A}_k^t}{\partial l} - h_{kl}^0 \\ &= \tilde{h}_{kl}(x + t\xi) - h_{kl}^0(x) \\ &+ t \left[\frac{\partial \tilde{A}_l}{\partial i}(x + t\xi) \frac{\partial \xi^i}{\partial k} - \frac{\partial \tilde{A}_k}{\partial i}(x + t\xi) \frac{\partial \xi^i}{\partial l} \right] \text{ in } B_\rho(x_0) \quad . \end{aligned} \quad (2.55)$$

Since we still subtract $|h^0|^2$ to the energy and not $|h^{0,t}|^2$, we will have some change in our stress-energy tensor compared to the classical one.

Since (φ, \tilde{A}) is a minimizer we have

$$\frac{d}{dt} \left[\int_{\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}} |\nabla_{\tilde{A}} \varphi^t|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + \frac{1}{2} \sum_{kl} h_{kl}^t h_{kl}^t \right] \Big|_{t=0} = 0 \quad .$$

After computations we get

$$\begin{aligned} \sum_{jkl} \int_{B_\rho(x_0)} \frac{\partial}{\partial j} \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + \frac{1}{2} h_{kl} \right] \xi^j \\ - 2 \frac{\partial}{\partial k} \left(\nabla_{\tilde{A}_k} \varphi \frac{\partial \varphi}{\partial j} \right) \xi^j \\ - \left[\frac{\partial}{\partial k} \left(h_{kl} \frac{\partial \tilde{A}_l}{\partial j} \right) - \frac{\partial}{\partial l} \left(h_{kl} \frac{\partial \tilde{A}_k}{\partial j} \right) \right] \xi^j \\ + h_{kl} \frac{\partial}{\partial j} h_{kl}^0 \xi^j = 0 \end{aligned} \tag{2. 56}$$

(2. 56) is true for any $\xi \in C_0^\infty(B_\rho(x_0))$, thus we have

$$\begin{aligned} \forall j = 1, 2, 3 \quad \frac{\partial}{\partial j} \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + \frac{1}{2} \sum_{kl} h_{kl} h_{kl} \right] \\ - 2 \sum_k \frac{\partial}{\partial k} (\nabla_{\tilde{A}_k} \varphi, \nabla_{\tilde{A}_j} \varphi) - 2 \sum_k \frac{\partial}{\partial k} (\nabla_{\tilde{A}_k} \varphi, i\varphi \tilde{A}_j) \\ - \sum_{kl} \frac{\partial}{\partial k} (h_{kl} \tilde{h}_{jl}) + \sum_{kl} \frac{\partial}{\partial l} (h_{kl} \tilde{h}_{jk}) \\ - \sum_{kl} \frac{\partial}{\partial k} \left(h_{kl} \frac{\partial \tilde{A}_j}{\partial l} \right) + \sum_{kl} \frac{\partial}{\partial l} \left(h_{kl} \frac{\partial \tilde{A}_j}{\partial k} \right) \\ + h_{kl} \frac{\partial h_{kl}^0}{\partial j} = 0 \quad \text{in } \mathcal{D}'(B_\rho(x_0)) \quad . \end{aligned}$$

We have

$$\frac{\partial}{\partial k} \left[(\nabla_{\tilde{A}_k} \varphi, i\varphi) \tilde{A}_j \right] = \frac{\partial}{\partial k} (\nabla_{\tilde{A}_k} \varphi, i\varphi) \tilde{A}_j + (\nabla_{\tilde{A}_k} \varphi, i\varphi) \frac{\partial \tilde{A}_j}{\partial k} \quad .$$

Since the second Higgs equation (2. 15) is

$$\forall k = 1, 2, 3 \quad \sum_l \frac{\partial}{\partial l} h_{lk} + (i\varphi, \nabla_{\tilde{A}_k} \varphi) = 0$$

and since $\sum_{kl} \frac{\partial^2}{\partial l \partial k} h_{lk} = 0$ (because $h_{lk} = -h_{kl}$) we have

$$\sum_k \frac{\partial}{\partial k} \left[(\nabla_{\tilde{A}_k} \varphi, i\varphi) \tilde{A}_j \right] = - \sum_{kl} \frac{\partial h_{lk}}{\partial l} \frac{\partial \tilde{A}_j}{\partial k} \quad . \tag{2. 57}$$

Note also that

$$\sum_{kl} - \frac{\partial}{\partial k} \left(h_{kl} \frac{\partial \tilde{A}_j}{\partial l} \right) + \frac{\partial}{\partial l} \left(h_{kl} \frac{\partial \tilde{A}_j}{\partial k} \right) = 2 \sum_{kl} \frac{\partial h_{kl}}{\partial l} \frac{\partial \tilde{A}_j}{\partial k} \quad . \tag{2. 58}$$

Combining (2. 57), (2. 57) and (2. 58) we have

$$\begin{aligned}
\forall j = 1, 2, 3 \quad & \frac{\partial}{\partial j} \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + \frac{1}{2} \sum_{kl} h_{kl} h_{kl} \right] \\
& - 2 \sum_k \frac{\partial}{\partial k} (\nabla_{\tilde{A}_k} \varphi, \nabla_{\tilde{A}_j} \varphi) - 2 \sum_{kl} \frac{\partial}{\partial k} (h_{kl} h_{jl}) \\
& - 2 \sum_{kl} \frac{\partial}{\partial k} (h_{kl} h_{jl}^0) + \sum_{kl} h_{kl} \frac{\partial h_{kl}^0}{\partial j} = 0 \text{ in } \mathcal{D}'(B_\rho(x_0)).
\end{aligned} \tag{2. 59}$$

The previous identity is gauge invariant, thus it holds on all of $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$. This is the conservation law verified by the stress tensor perturbed by the term we were looking for

$$\sum_{kl} -2 \frac{\partial}{\partial k} (h_{kl} h_{kl}^0) + h_{kl} \frac{\partial h_{kl}^0}{\partial j} .$$

Take $x_0 \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ (x_0 will be taken, in the next computation, as the centre of coordinates), for any $r < \text{dist}(x_0, \{x_1, \dots, x_n\})$ we multiply (2. 59) by x_j and integrate on $B_r(0)$, to give

$$\begin{aligned}
& - 3 \int_{B_r(0)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \\
& + \int_{\partial B_r(0)} (|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2)(x \cdot \nu) \\
& + 2 \int_{B_r(0)} |\nabla_{\tilde{A}} \varphi|^2 - 2 \int_{\partial B_r(0)} r (\nabla_{\tilde{A}} \varphi \cdot \nu, \nabla_{\tilde{A}} \varphi \cdot \nu) \\
& + 4 \int_{B_r(0)} |h|^2 - 2r \int_{\partial B_r(0)} h_{kl} h_{jl} \nu^k \nu^j \\
& - 2 \int_{B_r(0)} \sum_{jkl} x_j \frac{\partial}{\partial k} (h_{kl} h_{jl}^0) + \int_{B_r(0)} \sum_{jkl} x_j \frac{\partial}{\partial j} h_{kl}^0 h_{kl} = 0 \quad ,
\end{aligned} \tag{2. 60}$$

where ν denotes the outward unit normal of $\partial B_r(0)$. Denote by $h \wedge \nu$ the following 1-form $h \wedge \nu_l = h_{kl} \nu^k$.

We divide (2. 60) by r^2 and we get

$$\begin{aligned}
 & \frac{d}{dr} \left[\frac{1}{r} \int_{B_r(0)} |\nabla_{\bar{A}} \varphi|^2 + \frac{3}{2\varepsilon^2} (|\varphi|^2 - 1)^2 + 3|h|^2 \right] \\
 &= \frac{1}{r} \int_{\partial B_r(0)} 2|\nabla_{\bar{A}} \varphi \cdot \nu|^2 + \frac{(1 - |\varphi|^2)^2}{\varepsilon^2} + 2|h|^2 + 2|h \wedge \nu|^2 \\
 &- \frac{4}{r^2} \int_{B_r(0)} h^2 + \frac{2}{r^2} \int_{B_r(0)} \sum_{jkl} x_j \frac{\partial}{\partial k} (h_{kl} h_{jl}^0) \\
 &- \frac{1}{r^2} \int_{B_r(0)} \sum_{jkl} x_j \frac{\partial h_{kl}^0}{\partial j} h_{kl} \quad .
 \end{aligned} \tag{2. 61}$$

This looks like a monotonicity formula modulo the three last integrals in the right hand side of the equality. The first one $-\frac{4}{r^2} \int_{B_r(0)} h^2$ is not so bad because of the estimate (2. 37). We will see that it has a very small influence. The last two look very bad since they contain derivatives on h^0 which blow-up close to the x_i . In fact, we will see that there is some compensation between the two last terms. Extracting all the quantities of (2. 61) which contain derivatives of h_0 , we get

$$\begin{aligned}
 I &= \int_{B_r(0)} 2 \sum_{jkl} x_j h_{kl} \frac{\partial h_{jl}^0}{\partial k} - \sum_{jkl} x_j h_{kl} \frac{\partial h_{kl}^0}{\partial j} \\
 &= \int_{B_r(0)} 2 \sum_{k<l} h_{kl} \left(\sum_j x_j \frac{\partial h_{jl}^0}{\partial k} - x_j \frac{\partial h_{jk}^0}{\partial l} \right) \\
 &\quad - \sum_{k<l} h_{kl} \left(\sum_j x_j \frac{\partial h_{kl}^0}{\partial j} - \sum_j x_j \frac{\partial h_{lk}^0}{\partial j} \right)
 \end{aligned}$$

which is equal to

$$I = \int_{B_r(0)} 2 \sum_{k<l} h_{kl} \left(\sum_j x_j \left[\frac{\partial h_{jl}^0}{\partial k} + \frac{\partial h_{kj}^0}{\partial l} + \frac{\partial h_{lk}^0}{\partial j} \right] \right) .$$

Taking $k < l$ and j_0 such that $j_0 \neq k$ and $j_0 \neq l$, we have

$$\begin{aligned} \sum_j x_j \left[\frac{\partial h_{jl}^0}{\partial k} + \frac{\partial h_{kj}^0}{\partial l} + \frac{\partial h_{lk}^0}{\partial j} \right] &= x_k \left[\frac{\partial h_{lk}^0}{\partial k} + \frac{\partial h_{kk}^0}{\partial l} + \frac{\partial h_{kl}^0}{\partial k} \right] \\ &+ x_l \left[\frac{\partial h_{ll}^0}{\partial k} + \frac{\partial h_{lk}^0}{\partial l} + \frac{\partial h_{kl}^0}{\partial l} \right] \\ &+ x_{j_0} \left[\frac{\partial h_{j_0 l}^0}{\partial k} + \frac{\partial h_{k j_0}^0}{\partial l} + \frac{\partial h_{lk}^0}{\partial j_0} \right] \\ &= 0 + 0 + *dh^0 \times x_{j_0} = 0 \quad . \end{aligned}$$

Thus $I = 0$. Note also that $\sum \frac{\partial}{\partial k} h_{kl} = -(i\varphi, \nabla_{\tilde{A}_l} \varphi)$ and (2. 61) becomes

$$\begin{aligned} &\frac{d}{dr} \left[\frac{1}{r} \int_{B_r(x_0)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{3}{2\varepsilon^2} (1 - |\varphi|^2)^2 + 3|h|^2 \right] \\ &= \frac{1}{r} \int_{\partial B_r(x_0)} 2|\nabla_{\tilde{A}} \varphi \cdot \nu|^2 + \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2 + 2|h|^2 + 2|h \wedge \nu|^2 \quad (2. 62) \\ &- \frac{4}{r^2} \int_{B_r(x_0)} h^2 - \frac{2}{r^2} \int_{B_r(x_0)} \sum_{jl} h_{jl}^0 x^j (i\varphi, \nabla_{\tilde{A}_l} \varphi) \quad . \end{aligned}$$

We will deduce from (2. 62) the following Lemma

LEMMA 2.5. For ε sufficiently small

$$\begin{aligned} \forall x_0 \in \mathbb{R}^3 \quad \forall r > 0 \\ \frac{1}{r} \int_{B_r(x_0)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \leq C \log \frac{1}{\varepsilon} \end{aligned} \quad (2. 63)$$

where C does not depend on ε . ■

REMARK 2.8. In view of Lemma 2.1 and (2. 62) Lemma 2.5 is straightforward far away from the x_i but when we are close to them there are two difficulties. Firstly the last term in (2. 62) has to be controlled and secondly (2. 62) is true only for $r < \text{dist}(x_0, \{x_1, \dots, x_n\})$.

Proof of Lemma 2.5.

It suffices to prove (2. 63) for r small. Let $r_0 > 0$ be small enough such that the $B(x_i, r_0)$ are disjoint.

We only have to consider the case where x_0 belongs to a $B(x_i, r_0)$, and first we study the case $x_0 = x_i$.

x_i is taken as the center of coordinates.

Multiply the three components of (2. 59) by the corresponding coordinates

and sum the three terms that we obtain. We get

$$\begin{aligned} & \sum_{j=1}^3 x_j \frac{\partial}{\partial j} \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \right] \\ & - 2x_j \sum_{k=1}^3 \frac{\partial}{\partial k} (\nabla_{\tilde{A}_k} \varphi, \nabla_{\tilde{A}_j} \varphi) - 2x_j \sum_{kl} \frac{\partial}{\partial k} (h_{kl} h_{jl}) \\ & - 2x_j \sum_{kl} \frac{\partial}{\partial k} (h_{kl} h_{jl}^0) + x_j \sum_{kl} h_{kl} \frac{\partial h_{kl}^0}{\partial j} = 0 \quad \text{in } \mathcal{D}'(B_{r_0}(0) \setminus \{0\}). \end{aligned} \tag{2.64}$$

We saw that

$$\begin{aligned} & \sum_{j=1}^3 -2x_j \sum_{kl} \frac{\partial}{\partial k} (h_{kl} h_{jl}^0) + x_j \sum_{kl} h_{kl} \frac{\partial h_{kl}^0}{\partial j} \\ & = 2 \sum_{j=1}^3 x_j \sum_{l=1}^3 (i\varphi, \nabla_{\tilde{A}_l} \varphi) h_{jl}^0, \end{aligned}$$

and (2.64) can be written in the following form

$$\begin{aligned} & \sum_{j=1}^3 \frac{\partial}{\partial j} \left[x_j \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \right] \right] \\ & - 2 \sum_{j,k=1}^3 \frac{\partial}{\partial k} \left[x_j \left[(\nabla_{\tilde{A}_k} \varphi, \nabla_{\tilde{A}_j} \varphi) + h_{kl} h_{jl} \right] \right] \\ & - \left[|\nabla_{\tilde{A}} \varphi|^2 + \frac{3}{2\varepsilon^2} (1 - |\varphi|^2)^2 - |h|^2 \right] \\ & + 2 \sum_{l,j=1}^3 x_j h_{jl}^0 (i\varphi, \nabla_{\tilde{A}_l} \varphi) = 0 \quad \text{in } \mathcal{D}'(B_{r_0}(0) \setminus \{0\}) \end{aligned} \tag{2.65}$$

Remark that the left-hand side of (2.65) is a distribution on all of $B_{r_0}(0)$, we have L^1 functions, derivatives of L^1 functions and the last term is also in L^1 because $|x| \times |h^0| \simeq_0 \frac{c}{|x|}$ is in $L^p(B_{r_0}(0))$ for $p < 3$ and $(i\varphi, \nabla_{\tilde{A}} \varphi) \in L^2(B_{r_0}(0))$.

We claim that the left-hand side of (2.65) that we call D , whose support is included in $\{0\}$ is equal to 0 in $\mathcal{D}'(B_{r_0}(0))$. It can be written in the following form

$$D = \sum_{j=1}^3 \frac{\partial}{\partial j} [x_j G] + \sum_{j,k} \frac{\partial}{\partial k} [x_j H_{kj}] + F, \tag{2.66}$$

where G, H_{kj} and F belong to $L^1(B_{r_0}(0))$.

Since $\text{Supp } D \subset \{0\}$ in view of (2.66), D is at most a Dirac at 0. Take any $\psi \in C_0^\infty(B(0,1))$ and for any $0 < \rho < r_0$ denote by ψ_ρ the function

$$\psi_\rho(x) = \psi\left(\frac{x}{\rho}\right).$$

$$\begin{aligned} |\langle D, \psi_\rho \rangle| &= \left| - \int_{B_\rho(0)} \sum_j \frac{1}{\rho} \frac{\partial \psi}{\partial j} \left(\frac{x}{\rho}\right) x_j G \right. \\ &\quad \left. - \int_{B_\rho(0)} \sum_{jk} \frac{x_j}{\rho} \frac{\partial \psi}{\partial k} \left(\frac{x}{\rho}\right) H_{kj} + \int_{B_\rho(0)} \psi_\rho F \right|. \end{aligned}$$

Since $\frac{|x|}{\rho} |\nabla \psi|(\frac{x}{\rho})$ is uniformly bounded for $|x| < \rho$ independently of ρ it is clear that $|\langle D, \psi_\rho \rangle| \rightarrow 0$ as $\rho \rightarrow 0$ and we get that $D = 0$ in $\mathcal{D}'(B_{r_0}(0))$. Integrating D on $B_r(0)$ for any $r < r_0$, we get

$$\begin{aligned} &\frac{d}{dr} \left[\frac{1}{r} \int_{B_r} |\nabla_{\tilde{A}} \varphi|^2 + \frac{3}{2\varepsilon^2} (|\varphi|^2 - 1)^2 + 3|h|^2 \right] \\ &= \frac{1}{r} \int_{\partial B_r(0)} 2|\nabla_{\tilde{A}} \varphi \cdot \nu|^2 + \frac{1}{\varepsilon^2} (|\varphi|^2 - 1)^2 + 2|h|^2 + 2|h \wedge \nu|^2 \quad (2. 67) \\ &\quad - \frac{4}{r^2} \int_{B_r(0)} h^2 - \frac{2}{r^2} \int_{B_r(0)} \sum_l (i\varphi, \nabla_{\tilde{A}_l} \varphi) \sum_j x_j h_{jl}^0. \end{aligned}$$

The part of h^0 which diverges close to the pole 0, is in view of (2), $-\frac{1}{2}d_i * d(\frac{1}{|x|})$ and since $\sum_j x_j h_{jl}^0 = x \wedge *h^0$ and $x \wedge \nabla \frac{1}{|x|} = 0$ we have that

$$\left| \sum_{j=1}^3 x_j h_{jl}^0 \right| \leq C|x| \quad \text{in } B_r(0) \quad (2. 68)$$

We use the following notation, $E_r(x) = \int_{B_r(0)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{3}{2\varepsilon^2} (|\varphi|^2 - 1)^2 + 3|h|^2$ (2. 67) implies

$$\frac{d}{dr} \left(\frac{E_r}{r} \right) \geq -\frac{4}{r^2} \int_{B_r(0)} h^2 - \frac{2}{r} \int_{B_r(0)} \sum_l (i\varphi, \nabla_{\tilde{A}_l} \varphi) \sum_j x_j h_{jl}^0$$

From (2. 37) and (2. 68) we have

$$\forall r \leq r_0 \quad \left| \frac{4}{r^2} \int_{B_r(0)} h^2 + \frac{2}{r^2} \int_{B_r(0)} \sum_l (i\varphi, \nabla_{\tilde{A}_l} \varphi) \sum_j x_j h_{jl}^0 \right| \leq C \log \frac{1}{\varepsilon} .$$

Thus $\frac{d}{dr} \left(\frac{E_r}{r} \right) \geq -C \log \frac{1}{\varepsilon}$ and from Lemma 2.1 we have $\frac{E_{r_0}}{r_0} \leq \frac{C}{r_0} \log \frac{1}{\varepsilon}$. The two previous inequalities imply the result for $x_0 = x_i$.

Take now $x_0 \in B_{r_0}(x_i) \setminus \{x_i\}$. Let $d_{x_0} = |x_0 - x_i|$. We just have established above

$$\frac{1}{2d_{x_0}} \int_{B_{2d_{x_0}}(x_i)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \leq C \log \frac{1}{\varepsilon}.$$

Since $B_{\frac{d_{x_0}}{2}}(x_0) \subset B_{2d_{x_0}}(x_i)$ the previous inequality implies

$$\frac{2}{d_{x_0}} \int_{B_{\frac{d_{x_0}}{2}}(x_0)} |\nabla_{\tilde{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2) + |h|^2 \leq C \log \frac{1}{\varepsilon} \quad (2. 69)$$

For $r < \frac{d_{x_0}}{2}$ we have

$$\frac{d}{dr} \left[\frac{E_r(x_0)}{r} \right] \geq - \left| \frac{4}{r^2} \int_{B_r(x_0)} h^2 \right| - \left| \frac{2}{r^2} \int_{B_r(x_0)} |\nabla_{\tilde{A}}\varphi| \times |x \wedge *h^0| \right| \quad (2. 70)$$

We use the following bound $|x \wedge *h^0| \leq c \frac{|x|}{d_{x_0}^2}$ and (2. 70) implies

$$\frac{d}{dr} \left[\frac{E_r(x_0)}{r} \right] \geq -C \log \frac{1}{\varepsilon} - C \frac{r}{d_{x_0}^2} \left(\frac{E_r(x_0)}{r} \right)^{1/2}, \quad (2. 71)$$

where we have used (2. 37) and Cauchy-Schwarz inequality. In particular, (2. 71) implies

$$\frac{d}{dr} \left[\frac{E_r(x_0)}{r} + Cr \log \frac{1}{\varepsilon} \right] \leq -\frac{Cr}{d_{x_0}^2} \left[\frac{E_r(x_0)}{r} + Cr \log \frac{1}{\varepsilon} \right]^{1/2} \quad (2. 72)$$

Integrating (2. 70) between $\frac{d_{x_0}}{2}$ and any $\rho < \frac{d_{x_0}}{2}$, the bound (2. 69) yields

$$\forall \rho \leq \frac{d_{x_0}}{2}, \quad \frac{E_\rho(x_0)}{\rho} \geq C \log \frac{1}{\varepsilon} \quad ,$$

and the lemma is proved for any $x_0 \in \cup_{i=1}^n B(x_i, r_0)$. On the remains of \mathbb{R}^3 , as we mentioned above, the proof of (2. 63) is much more straightforward. ■

3. THE η -COMPACTNESS LEMMA

This part is devoted to the proof of the following lemma.

LEMMA 3.1. There exists $\eta > 0, \nu$ and $\lambda > 1$ such that for any ε sufficiently small, for any $(\varphi, \tilde{A}) \in V$ minimizer of G_ε , for any $x_0 \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$ and for any $\lambda\varepsilon \leq \rho \leq \min(\nu, \eta d_{x_0})$, where $d_{x_0} = \text{dist}(x_0, \{x_1, \dots, x_n\})$, if

$$\frac{1}{\rho} \int_{B_\rho(x_0)} |\nabla_{\tilde{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 + |h|^2 \leq \eta \log\left(\frac{\rho}{\varepsilon}\right)$$

then

$$|\varphi| \geq \frac{1}{2} \quad \text{in} \quad B_{\rho/2}(x_0)$$

■

Proof of Lemma 3.1.

First of all, we introduce some useful notations. Denote

$$\begin{aligned} E_r(x_0) &= \int_{B_r(x_0)} |\nabla_{\tilde{A}}\varphi|^2 + \frac{3}{2\varepsilon^2}(1 - |\varphi|^2)^2 + 3|h|^2, \\ I_r(x_0) &= \int_{\partial B_r(x_0)} |\nabla_{\tilde{A}}\varphi|^2 + \frac{3}{2\varepsilon^2}(1 - |\varphi|^2)^2 + 3|h|^2 = \frac{d}{dr} E_r(x_0), \\ F_r(x_0) &= \int_{B_r(x_0)} 2|\nabla_{\tilde{A}}\varphi \cdot \nu|^2 + \frac{1}{\varepsilon^2}(1 - |\varphi|^2)^2 + 2|h|^2 + 2|h \wedge \nu|^2 \end{aligned}$$

$$J_r(x_0) = \int_{\partial B_r(x_0)} 2|\nabla_{\tilde{A}}\varphi \cdot \nu|^2 + \frac{1}{\varepsilon^2}(1 - |\varphi|^2)^2 + 2|h|^2 + 2|h \wedge \nu|^2 = \frac{dF_r(x_0)}{dr}$$

and

$$R_r(x_0) = -\frac{4}{r^2} \int_{B_r(x_0)} h^2 - \frac{2}{r^2} \int_{B_r(x_0)} \sum_{jl} h_{jl}^0 x^j (i\varphi, \nabla_{\tilde{A}_l} \varphi) \quad .$$

Using those notations (2. 62) becomes

$$\frac{d}{dr} \left[\frac{E_r(x_0)}{r} \right] = \frac{J_r(x_0)}{r} + R_r(x_0) \quad . \tag{3. 1}$$

As in the proof of Lemma 2.5 we use the bound

$$|R_r(x_0)| \leq C \log \frac{1}{\varepsilon} + C \frac{r}{d_{x_0}^2} \left(\frac{E_r(x_0)}{r} \right)^{1/2} \quad . \tag{3. 2}$$

Suppose for some $\lambda\varepsilon < \rho < \frac{d_{x_0}}{2}$ we have $E_\rho/\rho \leq \eta \log(\rho/\varepsilon)$ ($\lambda > 1$ will be fixed later but we already consider couples (η, λ) such that $\log \lambda > \frac{1}{\eta}$). As in the proof of Lemma 2.5, using (3. 1) and (3. 2), we prove that $\forall r < \rho \quad \frac{E_r(x_0)}{r} \leq C\eta \log(\frac{\rho}{\varepsilon})$ and combined with (3. 2) this gives

$$\forall r < \rho \quad |R_r(x_0)| \leq C \log \frac{1}{\varepsilon} + C \frac{r}{d_{x_0}^2} \left(\eta \log \left(\frac{\rho}{\varepsilon} \right) \right)^{1/2} \quad . \tag{3. 3}$$

Integrating (3. 2) between ρ and $\lambda\varepsilon$ (λ will be fixed later) we have

$$\int_{\lambda\varepsilon}^{\rho} \frac{J_r(x_0)}{r} dr \leq \eta \log \left(\frac{\eta}{\varepsilon} \right) + C\rho \log \frac{1}{\varepsilon} + C \frac{\rho^2}{d_{x_0}^2} \left(\eta \log \left(\frac{\rho}{\varepsilon} \right) \right)^{1/2} \quad . \tag{3. 4}$$

It is clear that there exists $\nu > 0$ independent of ε for ε small enough such that, if $\rho < \nu$ and $\rho < \frac{d_{x_0}}{2}$, we have

$$C\rho \log \frac{1}{\varepsilon} + C \frac{\rho^2}{d_{x_0}^2} (\eta \log(\frac{\rho}{\varepsilon}))^{1/2} \leq \eta \log \left(\frac{\rho}{\varepsilon} \right)$$

(λ has also to be chosen sufficiently large (independently of ε) such that $c(\eta \log \lambda)^{1/2} \leq \eta \log \lambda$).

Finally, for $\lambda\varepsilon \leq \rho \leq \min [d_{x_0}/2, \nu]$, we have

$$\int_{\lambda\varepsilon}^{\rho} \frac{J_r(x_0)}{r} dr \leq 2\eta \log \left(\frac{\rho}{\varepsilon} \right) \tag{3. 5}$$

Thus we have also

$$\int_{\lambda\varepsilon}^{\rho} \frac{J_r(x_0)}{r} + \frac{J_{r/2}(x_0)}{r} dr \leq C\eta \log \left(\frac{\rho}{\varepsilon} \right) \quad .$$

Moreover

$$\begin{aligned} \int_{\lambda\varepsilon}^{\rho} \frac{J_r(x_0)}{r} &= \int_{\lambda\varepsilon}^{\rho} \frac{1}{r} \frac{d}{dr} F_r(x_0) = \left[\frac{F_\rho(x_0)}{\rho} - \frac{F_{\lambda\varepsilon}(x_0)}{\lambda\varepsilon} \right] + \\ &+ \int_{\lambda\varepsilon}^{\rho} \frac{F_r(x_0)}{r^2} dr \end{aligned}$$

and this implies

$$\frac{F_{\lambda\varepsilon}(x_0)}{\lambda\varepsilon} + C\eta \log\left(\frac{\rho}{\varepsilon}\right) \geq \int_{\lambda\varepsilon}^{\rho} \frac{F_r(x_0)}{r^2} dr,$$

but $F_{\lambda\varepsilon}(x_0)/\lambda\varepsilon \leq E_{\lambda\varepsilon}(x_0)/\lambda\varepsilon \leq C\eta \log(\frac{\rho}{\varepsilon})$, and this yields

$$\int_{\lambda\varepsilon}^{\rho} \frac{J_r(x_0)}{r} + \frac{J_{r/2}(x_0)}{r} + \frac{F_r(x_0)}{r^2} dr \leq C\eta \log\left(\frac{\rho}{\varepsilon}\right) \quad (3.6)$$

Applying the mean value formula in (3.6) we get the existence of $r_0 \in [\lambda\varepsilon, \rho]$ such that

$$J_{r_0}(x_0) + J_{r_0/2}(x_0) + \frac{F_{r_0}}{r_0}(x_0) \leq C\eta \quad (3.7)$$

where C is independent on ε and η . Our aim is to prove that (3.7) implies that

$$\frac{E_{r_0/2}(x_0)}{r_0/2} \leq f(\eta) \quad \text{where} \quad f(\eta) \rightarrow 0 \text{ as } \eta \rightarrow 0 \quad (3.8)$$

And (3.8) will imply (see below at the end of the proof of the Lemma),

$$\begin{aligned} \frac{E_{\varepsilon}(x_0)}{\varepsilon} &\leq f(\eta) + \int_{\varepsilon}^{r_0/2} |R_r(x_0)| \\ &\leq g(\eta) \quad \text{with } g \rightarrow 0 \text{ as } \eta \rightarrow 0 \end{aligned} \quad (3.9)$$

This yields in particular

$$\frac{1}{|B_{\varepsilon}(x_0)|} \int_{B_{\varepsilon}(x_0)} (1 - |\varphi|^2)^2 \leq g(\eta) \quad (3.10)$$

and since $|\nabla|\varphi|^2| \leq |\nabla|\varphi|| \leq |\nabla_{\bar{A}}\varphi| \leq C/\varepsilon$, for $g(\eta)$ sufficiently small (3.10) implies $|\varphi| \geq \frac{1}{2}$ in $B_{\lambda\varepsilon}(x_0)$ and in particular $|\varphi|(x_0) \geq \frac{1}{2}$, which is the desired result.

The difficulty now is to obtain (3.8) from (3.7). (3.7) says that the total energy is small except possibly the tangential projections of $\nabla_{\bar{A}}\varphi$ on all the spheres in $B_{r_0}(x_0)$ of center x_0 .

First of all note that we can take $r_0 \in [\lambda\varepsilon, \frac{\eta}{Ln\frac{1}{\varepsilon}}]$. This is clear if $\rho < \frac{\eta}{\log\frac{1}{\varepsilon}}$, otherwise, since ρ verifies

$$\frac{E_{\rho}(x_0)}{\rho} \leq \eta \log\left(\frac{\rho}{\varepsilon}\right) \quad ,$$

as we mentioned above we have

$$\forall r < \rho \quad \frac{E_r(x_0)}{r} \leq C\eta \log\left(\frac{\rho}{\varepsilon}\right) .$$

Take $r = \frac{\eta}{\log\frac{1}{\varepsilon}} = r'$ and we have

$$\begin{aligned} \frac{E_{r'}(x_0)}{r'} &\leq C\eta \log\left(\frac{r'}{\varepsilon}\right) + C\eta \log\left(\frac{\rho}{r'}\right) \\ &\leq C\eta \log\left(\frac{r'}{\varepsilon}\right) + C\eta \log\left[\frac{1}{\eta} \log\frac{1}{\varepsilon}\right] \end{aligned}$$

but $\log\left[\frac{1}{\eta} \log \frac{1}{\varepsilon}\right] < \log\left(\frac{\eta}{\varepsilon \log \frac{1}{\varepsilon}}\right)$ (for ε sufficiently small), thus we have $E_{r'}(x_0)/r' \leq C\eta \log(r'/\varepsilon)$ and we get (3. 7) for $r_0 \in [\lambda\varepsilon, \eta/\log \frac{1}{\varepsilon}]$.

Since $r_0 < \eta \operatorname{dist}(x_0, \{x_1, \dots, x_n\})$ we can clearly include $B_{r_0}(x_0)$ in some U_l . We trivialize E over $B_{r_0}(x_0)$ and consider the particular representative of (φ, \hat{A}) in $B_{r_0}(x_0) \subset U_l$ given by Proposition 2.1 modulo the multiplication of the fiber over y by $e^{-iA(x_0) \cdot y}$ for each y in $B_{r_0}(x_0)$, and we have

$$\tilde{A}_l(y) = A_l^0(y) + A(y) - A(x_0) \quad .$$

Because of (2. 39) we have

$$\begin{aligned} |\tilde{A}_l(y)| &\leq |A_l^0(y)| + C \frac{\eta^{1/2}}{(Ln \frac{1}{\varepsilon})^{1/2}} (\log 1/\varepsilon)^{1/2} \\ &\leq \frac{C}{d_{x_0}} + C\eta^{1/2} \quad . \end{aligned} \tag{3. 11}$$

We change the scale in view of working in a unit ball and multiply by $\frac{1}{r_0}$. Denote in $B(x_0, 1)$

$$\begin{aligned} \bar{\varphi}(x) &= \varphi(r_0x) \\ \hat{A}(x) &= r_0\tilde{A}(r_0x) = r_0A_l^0(r_0x) + r_0A(r_0x) \\ \bar{A}(x) &= r_0A(r_0x) \\ \text{and } \bar{h}(x) &= d\bar{A}(x) = r_0^2dA(r_0x) = r_0^2h(r_0x) \quad . \end{aligned}$$

Thus we have, in particular,

$$\nabla_{\hat{A}} \bar{\varphi}(x) = r_0 \nabla_{\bar{A}} \varphi(r_0x)$$

and

$$d^* \bar{h}(x) = r_0^3 d^* h(r_0x) \quad ,$$

moreover $(\bar{\varphi}, \hat{A})$ minimizes the functional

$$\bar{G}_{\varepsilon, r_0}(\bar{\varphi}, \hat{A}) = \int_{B_1} |\nabla_{\hat{A}} \bar{\varphi}|^2(x) + \frac{1}{2} \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2(x) + \frac{1}{r_0^2} |\bar{h}|^2(x) dx$$

and verifies the equations

$$\begin{cases} \nabla_{\hat{A}}^2 \bar{\varphi} &= \left(\frac{r_0}{\varepsilon}\right)^2 \bar{\varphi}(1 - |\bar{\varphi}|^2) \\ -\frac{d^* \bar{h}}{r_0^2} &= (i\bar{\varphi}, \nabla_{\hat{A}} \bar{\varphi}) \quad . \end{cases}$$

(3. 12) implies, in particular,

$$(i\bar{\varphi}, \nabla_{\hat{A}}^2 \bar{\varphi})(x) = 0 \tag{3. 12}$$

where

$$\begin{aligned} \nabla_{\hat{A}}^2 \bar{\varphi} &= \sum_{k=1}^3 \left(\frac{\partial}{\partial k} - i\hat{A}_k\right) \left(\frac{\partial \bar{\varphi}}{\partial k} - i\hat{A}_k \bar{\varphi}\right) \\ &= \Delta \bar{\varphi} - 2i\hat{A} \cdot \nabla \bar{\varphi} - id^* \hat{A} \bar{\varphi} - |\bar{A}|^2 \bar{\varphi} \quad . \end{aligned}$$

But we have chosen a gauge such that $d^*\hat{A} = 0$ and (3. 12) becomes

$$(i\bar{\varphi}, \Delta\bar{\varphi}) = \hat{A} \cdot \nabla |\bar{\varphi}|^2 \quad . \quad (3. 13)$$

For $\sigma \in [\frac{1}{2}, 1]$, denote by Δ_σ the Laplace Beltrami operator on $\partial B_\sigma \subset \mathbb{R}^3$ and by d_τ and $(*)$ we denote, respectively, the exterior derivative and the Hodge operator on $\partial B_\sigma \subset \mathbb{R}^3$. By definition we have

$$\Delta_\sigma = d_\tau(*)d_\tau(*) + (*)d_\tau(*)d_\tau \quad \text{on } \partial B_\sigma \quad .$$

If f is a 0-form in B_1 , for $\sigma \in [\frac{1}{2}, 1]$ we have

$$\Delta f = \Delta_\sigma f + \frac{\partial^2 f}{\partial r^2} + \frac{2}{r} \frac{\partial f}{\partial r} \quad .$$

In particular we have

$$\begin{aligned} \Delta \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) &= \Delta_\sigma \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) + \frac{\partial^2}{\partial \sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \\ &+ \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \quad . \end{aligned} \quad (3. 14)$$

Using Lemma A.3 and (3. 14) we have,

$$\begin{aligned} \Delta \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) &= \frac{\partial}{\partial \sigma} (i\bar{\varphi}, \Delta_\sigma \bar{\varphi}) + 2(*)d_\tau \left[\left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_\tau \bar{\varphi} \right) \right] \\ &+ \frac{2}{\sigma} (i\bar{\varphi}, \Delta_\sigma \bar{\varphi}) + \frac{\partial^2}{\partial \sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \\ &+ \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) . \end{aligned} \quad (3. 15)$$

Using the Laplace Beltrami operator (3. 13) is equivalent to

$$(i\bar{\varphi}, \Delta_\sigma \bar{\varphi}) = \hat{A} \cdot \nabla |\bar{\varphi}|^2 - \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - \frac{2}{\sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) . \quad (3. 16)$$

Combining (3. 15) and (3. 16) we get

$$\begin{aligned} \Delta \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) &= 2(*)d_\tau \left[\left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_\tau \bar{\varphi} \right) \right] + \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \\ &- \frac{\partial^2}{\partial \sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) + \frac{2}{\sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \\ &+ \frac{2}{\sigma} (\hat{A} \cdot \nabla |\bar{\varphi}|^2) - \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - \frac{4}{\sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \\ &+ \frac{\partial^2}{\partial \sigma^2} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) + \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \quad . \end{aligned} \quad (3. 17)$$

The cancelation of the two maximal order terms $\frac{\partial^2}{\partial \sigma^2}(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma})$, in the right hand side of (3. 17), is a crucial point of the proof. We have

$$\begin{aligned} \Delta \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) &= 2(*)d_{\top} \left[\left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_{\top} \bar{\varphi} \right) \right] - \frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} \left[\sigma \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \right] \\ &+ \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nabla |\bar{\varphi}|^2) + \frac{2}{\sigma} \hat{A} \cdot \nabla |\bar{\varphi}|^2 \quad . \end{aligned} \tag{3. 18}$$

Denote by T_1 the following set $T_1 = B(x_0, 1) \setminus B(x_0, 1/2)$. Let V be the 1-form solution of

$$\begin{cases} \Delta V = \left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_{\top} \bar{\varphi} \right) & \text{in } T_1 \\ V = 0 & \text{on } \partial T_1 \quad . \end{cases} \tag{3. 19}$$

By classical elliptic estimates we have $\forall q < 3/2$

$$\begin{aligned} \left(\int_{T_1} |\nabla V|^q \right)^{1/q} &\leq C_q \left\| \left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_{\top} \bar{\varphi} \right) \right\|_{L^1} \\ &\leq C_q \left(\int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} \left(\int_{T_1} |d_{\top} \bar{\varphi}|^2 \right)^{1/2} \quad . \end{aligned} \tag{3. 20}$$

Moreover

$$d_{\top} \Delta V = d_{\top} \left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_{\top} \bar{\varphi} \right) \quad ,$$

thus

$$(*)d_{\top} \left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*)d_{\top} \bar{\varphi} \right) = \langle d_{\top} \Delta V, \omega \rangle \quad , \tag{3. 21}$$

where $\omega = *d\sigma$. Applying Lemma A.4 we have

$$\begin{aligned} \langle d_{\top} \Delta V, \omega \rangle &= \langle \Delta d_{\top} V, \omega \rangle + \frac{4}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \\ &- \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \langle dV, \omega \rangle \\ &= \langle \Delta[\langle dV, \omega \rangle \omega], \omega \rangle + \frac{4}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \\ &- \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \langle dV, \omega \rangle \quad . \end{aligned}$$

Applying the first formula of Lemma A.4 we have

$$\begin{aligned}
 \langle d_{\tau} \Delta V, \omega \rangle &= \Delta \langle dV, \omega \rangle \\
 &- \frac{2}{\sigma^2} \langle dV, \omega \rangle + \frac{4}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \\
 &- \frac{2}{\sigma} \frac{\partial}{\partial \sigma} \langle dV, \omega \rangle \\
 &= \Delta \langle dV, \omega \rangle + \frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \quad .
 \end{aligned}
 \tag{3. 22}$$

Combining (3. 21) and (3. 22) we get

$$\Delta \langle dV, \omega \rangle = (*) d_{\tau} \left(\frac{\partial \bar{\varphi}}{\partial \sigma}, i(*) d_{\tau} \bar{\varphi} \right) - \frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \quad . \tag{3. 23}$$

Subtracting (3. 23) and (3. 18)

$$\begin{aligned}
 \Delta \left[\left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - 2 \langle dV, \omega \rangle \right] &= -\frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} \left[\left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \right] \\
 + \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nabla |\bar{\varphi}|^2) + \frac{2}{\sigma} \hat{A} \cdot \nabla |\varphi|^2 + \frac{4}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \quad .
 \end{aligned}
 \tag{3. 24}$$

Let H_1 be the solution of

$$\begin{cases}
 \Delta H_1 &= -\frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} \left[\sigma \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) \right] + \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nabla |\bar{\varphi}|^2) \\
 &+ \frac{2}{\sigma} \hat{A} \cdot \nabla |\bar{\varphi}|^2 + \frac{4}{\sigma^2} \frac{\partial}{\partial \sigma} [\sigma \langle dV, \omega \rangle] \\
 H_1 = 0 &\text{on } \partial T_1 \quad .
 \end{cases}
 \tag{3. 25}$$

Remark that, since $d^* \hat{A} = 0$, we have $\hat{A} \cdot \nabla |\bar{\varphi}|^2 = d^*(\hat{A} |\bar{\varphi}|^2)$ and

$$\frac{\partial}{\partial \sigma} d^*(\hat{A} |\bar{\varphi}|^2) - d^* \left(\frac{\partial}{\partial \sigma} (\hat{A} |\bar{\varphi}|^2) \right) = -\frac{1}{\sigma} d^*(\hat{A} |\bar{\varphi}|^2) + \frac{1}{\sigma} \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nu |\bar{\varphi}|^2)$$

Thus

$$\begin{cases}
 \Delta H_1 = -\frac{2}{\sigma^2} \frac{\partial}{\partial \sigma} \left[\sigma \left[\left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - 2 \langle dV, \omega \rangle \right] \right] + \\
 + d^* \left(\frac{\partial}{\partial \sigma} (\hat{A} |\bar{\varphi}|^2) \right) + \frac{1}{\sigma} d^*(\hat{A} \cdot \nu |\bar{\varphi}|^2) + \frac{1}{\sigma} \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nu |\bar{\varphi}|^2) \text{ in } T_1 \\
 H_1 = 0 \quad \text{on } \partial T_1 \quad .
 \end{cases}
 \tag{3. 26}$$

By classical elliptic estimates we have $\forall p < \frac{3}{2}$

$$\begin{aligned} \left(\int_{T_1} |\nabla H_1|^p \right)^{1/p} &\leq C_p \left[\left(\int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^p \right)^{1/p} + \left(\int_{T_1} |\nabla v|^p \right)^{1/p} \right. \\ &\quad \left. + \left(\int_{T_1} \left| \frac{\partial}{\partial \sigma} (\hat{A}|\bar{\varphi}|^2) \right|^p \right)^{1/p} + \left(\int_{T_1} |\hat{A}|\bar{\varphi}|^2|^p \right)^{1/p} \right] \end{aligned} \tag{3. 27}$$

Let H_2 be the solution of

$$\begin{cases} \Delta H_2 = 0 & \text{in } T_1 \\ H_2 = (i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma}) - 2 \langle dV, \omega \rangle \\ = (i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma}) - 2(*)d_{\tau}V = (i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma}) & \text{on } \partial T_1 \end{cases} \tag{3. 28}$$

Let T'_1 be the following subannulus of T_1 , $T'_1 = B_{7/8}(x_0) \setminus B_{5/8}(x_0)$. By classical results on harmonic functions we have

$$\left(\int_{T'_1} |\nabla H_2|^2 \right)^{1/2} \leq C \left(\int_{\partial T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} \tag{3. 29}$$

Combining (3. 24), (3. 25) and (3. 28), it is clear that

$$\left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) = 2 \langle dV, \omega \rangle + H_1 + H_2 \quad \text{in } T_1 \tag{3. 30}$$

The conclusion at this point is that we have been able to write $(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma})$ (which includes the radial derivative) using tangential derivatives: $\langle dV, \omega \rangle$ plus some remaining terms H_1 and H_2 which are not so bad (we will see this below).

For $1 < p < \frac{3}{2}$ and $1 < q < \frac{3}{2}$ fixed and to be made precise later, applying the mean value formula, we can find $t \in [\frac{5}{8}, \frac{7}{8}]$ such that, for C sufficiently large, we have simultaneously

$$\begin{aligned} \left(\int_{\partial B_t} |\nabla H_1|^p \right)^{\frac{1}{p}} &\leq C \left(\int_{T_1} |\nabla H_1|^p \right)^{1/p}, \\ \left(\int_{\partial B_t} |\nabla H_2|^2 \right)^{1/2} &\leq C \left(\int_{T'_1} |\nabla H_2|^2 \right)^{1/2} \\ \left(\int_{\partial B_t} |\nabla V|^q \right)^{1/q} &\leq C \left(\int_{T_1} |\nabla V|^q \right)^{1/q}, \\ \left(\int_{\partial B_t} |\nabla_{\tau} \bar{\varphi}|^2 \right)^{1/2} &\leq C \left(\int_{T_1} |\nabla_{\tau} \bar{\varphi}|^2 \right)^{1/2}, \\ &\int_{\partial B_t} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 + \frac{1}{2} \left(\frac{r_0}{\varepsilon} \right)^2 (1 - |\bar{\varphi}|^2)^2 \\ &\leq C \int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 + \frac{1}{2} \left(\frac{r_0}{\varepsilon} \right)^2 (1 - |\bar{\varphi}|^2)^2 \end{aligned} \tag{3. 31}$$

and

$$\left(\int_{\partial B_t} |\nabla \bar{A}|^p \right)^{1/p} \leq C \left(\int_{T_1} |\nabla \bar{A}|^p \right)^{1/p}.$$

We have also chosen t such that

$$\begin{aligned} \int_{\partial B_{tr_0}} \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2 &= \int_{\partial B_t} \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 \\ &\leq C \int_{T_1} \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 = \frac{1}{r_0} \int_{B_{r_0}(x_0) \setminus B_{r_0/2}(x_0)} \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2, \end{aligned}$$

and combined with (3. 7) the previous inequality gives

$$\int_{\partial B_{tr_0}} \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2 \leq C\eta \quad . \tag{3. 32}$$

Since $\|\nabla_{\bar{A}}\varphi\|_{\infty} \leq C/\varepsilon$ in $B_{r_0}(x_0)$ (from Lemma 2.4), we have, in particular, $\|\nabla|\varphi|\|_{\infty} \leq C/\varepsilon$ and, in view of (3. 32), for η sufficiently small we have $|\varphi| \geq \frac{1}{2}$ on $\partial B_{tr_0}(x_0)$ that is $|\bar{\varphi}| \geq \frac{1}{2}$ on $\partial B_t(x_0)$.

Since $H_{dR}^1(\partial B_t(x_0)) = 0$ there exists $\bar{\psi} \in C^\infty(\partial B_\rho, \mathbb{R})$ such that

$$e^{i\bar{\psi}} = \frac{\bar{\varphi}}{|\bar{\varphi}|} \quad \text{on } \partial B_\rho(x_0) \quad \text{and} \quad \int_{\partial B_t} \bar{\psi} \in [0, 2\pi) \quad .$$

We have

$$d_\top \bar{\psi} = \left(i \frac{\bar{\varphi}}{|\bar{\varphi}|}, d_\top \left(\frac{\bar{\varphi}}{|\bar{\varphi}|} \right) \right) = \frac{1}{|\bar{\varphi}|^2} (i\bar{\varphi}, d_\top \bar{\varphi}) \tag{3. 33}$$

On the other hand one verifies that for any $a \in C^\infty(\partial B_t, \mathbb{C})$ we have

$$(*)d_\top(*) (ia, d_\top a) = (ia, \Delta_t a) \quad . \tag{3. 34}$$

Applying (3. 34) to $\bar{\varphi}$ we have, using (3. 16),

$$\begin{aligned} (*)d_\top(*) (i\bar{\varphi}, d_\top \bar{\varphi}) &= (i\bar{\varphi}, \Delta_\rho \bar{\varphi}) \\ &= -\frac{\partial}{\partial \sigma} \left(i\bar{\varphi} \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - \frac{2}{\sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) + d^*(\hat{A}|\varphi|^2) \end{aligned} \tag{3. 35}$$

We will replace, in (3. 35), $d^*(\hat{A}|\varphi|^2)$ by $d_\top^{(*)}(\hat{A}|\bar{\varphi}|^2) + \frac{\partial}{\partial \sigma}(\hat{A}.\nu|\bar{\varphi}|^2) + \frac{2}{\sigma}\hat{A}.\nu|\bar{\varphi}|^2$ and combining (3. 33) and (3. 35) we obtain

$$\begin{aligned} \Delta_t \bar{\psi} &= -\frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) - \frac{2}{\sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) + d_\top^{(*)}(\hat{A}|\varphi|^2) + \frac{\partial}{\partial \sigma}(\hat{A}.\nu|\bar{\varphi}|^2) \\ &\quad + \frac{2}{\sigma}\hat{A}.\nu|\bar{\varphi}|^2 + d_\top^{(*)} \left(\left(\frac{1}{|\bar{\varphi}|^2} - 1 \right) (i\bar{\varphi}, d_\top \bar{\varphi}) \right) \quad \text{on } \partial B_t(x_0) \quad . \end{aligned} \tag{3. 36}$$

$\frac{\partial}{\partial \sigma}(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma})$ looks a priori bad in the right hand side of (3. 36) in view of estimating $\bar{\psi}$ but using (3. 30) we have, using also Lemma A.1,

$$\begin{aligned} \frac{\partial}{\partial \sigma} \left(i\bar{\varphi}, \frac{\partial \bar{\varphi}}{\partial \sigma} \right) &= 2 \frac{\partial}{\partial \sigma} \langle dV, \omega \rangle + \frac{\partial H_1}{\partial \sigma} + \frac{\partial H_2}{\partial \sigma} \\ &= 2 \langle \frac{\partial}{\partial \sigma} dV, \omega \rangle + \frac{\partial H_1}{\partial \sigma} + \frac{\partial H_2}{\partial \sigma} \\ &= 2(*)d_{\top} \frac{\partial V}{\partial \sigma} - \frac{2}{\sigma} (*)d_{\top} V + \frac{\partial H_1}{\partial \sigma} + \frac{\partial H_2}{\partial \sigma} . \end{aligned} \tag{3. 37}$$

Taking $q = \frac{4}{3}$ and $p = \frac{4}{3}$, classical elliptic estimates yield the following upper bound for $\int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\psi}|^{4/3}$.

$$\begin{aligned} \left(\int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\psi}|^{\frac{4}{3}} \right)^{3/4} &\leq C \left[\left(\int_{\partial B_t(x_0)} |\nabla V|^{\frac{4}{3}} \right)^{3/4} \right. \\ &+ \left(\int_{\partial B_t(x_0)} |\nabla H_1|^{\frac{4}{3}} \right)^{3/4} + \left(\int_{\partial B_t(x_0)} |\nabla H_2|^2 \right)^{1/2} \\ &+ \left(\int_{\partial B_t(x_0)} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} + \left(\int_{\partial B_t(x_0)} |\hat{A}|\bar{\varphi}|^2 \right)^{4/3} \Big)^{3/4} \\ &+ \left(\int_{\partial B_t(x_0)} \left| \frac{\partial}{\partial \sigma} (\hat{A} \cdot \nu |\bar{\varphi}|^2) \right|^{4/3} \right)^{3/4} \\ &+ \left. \left(\int_{\partial B_t(x_0)} (1 - |\bar{\varphi}|^2)^{4/3} |\nabla_{\top} \bar{\varphi}|^{4/3} \right)^{3/4} \right] . \end{aligned} \tag{3. 38}$$

Combining (3. 20), (3. 27), (3. 29) and (3. 31), (3. 38) implies

$$\begin{aligned} \left(\int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\psi}|^{\frac{4}{3}} \right)^{3/4} &\leq \\ &\leq C \left[\left(\int_{T_1} |\nabla_{\top} \bar{\varphi}|^2 \right)^{1/2} \left(\int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} \right. \\ &+ \left(\int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} + \left(\int_{\partial T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right)^{1/2} \\ &+ \left(\int_{T_1} |\nabla \hat{A}|^2 \right)^{1/2} + \left(\int_{T_1} (1 - |\bar{\varphi}|^2)^4 \right)^{1/4} \left(\int_{T_1} |\nabla_{\top} \bar{\varphi}|^2 \right)^{\frac{1}{2}} + \eta \Big] , \end{aligned} \tag{3. 39}$$

where we have also used the bound $\|\hat{A}\|_\infty \leq \frac{r_0}{dx_0} + r_0\eta^{1/2} \leq C\eta$ established above (see (3. 11)). Let $\psi_t = \int_{\partial B_t(x_0)} \bar{\psi} \in [0, 2\pi)$ by Poincaré inequality and Sobolev embedding we have

$$\|\bar{\psi} - \bar{\psi}_t\|_{H^{1/2}(\partial B_t(x_0))} \leq C \left(\int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\psi}|^{4/3} \right)^{3/4} . \tag{3. 40}$$

Let $\tilde{\psi}$ be the harmonic extension of $\bar{\psi}$ in $B_t(x_0)$. We have from (3. 40) and classical results on harmonic functions,

$$\left(\int_{B_t(x_0)} |\nabla \tilde{\psi}|^2 \right)^{1/2} \leq C \left(\int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\psi}|^{4/3} \right)^{3/4} . \tag{3. 41}$$

Let $\bar{\xi} : B_t(x_0) \rightarrow \mathbb{R}_+$ such that

$$\begin{cases} \bar{\xi} = |\bar{\varphi}| & \text{on } \partial B_t(x_0) \\ \bar{\xi} \text{ is a minimizer of } \int_{\partial B_t(x_0)} |\nabla \bar{\xi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi})^2 \end{cases} . \tag{3. 42}$$

Clearly $\bar{\xi}$ is the solution of

$$\begin{cases} -\left(\frac{\varepsilon}{r_0}\right)^2 \Delta \bar{\xi} + \bar{\xi} = 1 & \text{in } B_t(x_0) \\ \bar{\xi} = |\bar{\varphi}| & \text{on } \partial B_t(x_0) \end{cases} . \tag{3. 43}$$

By the maximum principle we have

$$\max_{B_t(x_0)} (1 - \bar{\xi}) \leq \max_{\partial B_t(x_0)} (1 - |\bar{\varphi}|) \leq \frac{1}{2} . \tag{3. 44}$$

Multiply (3. 43) by $1 - \bar{\xi}$ and integrate on $B_t(x_0)$, we have

$$\int_{B_t(x_0)} |\nabla \bar{\xi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi})^2 = \int_{\partial B_t(x_0)} (1 - \bar{\xi}) \frac{\partial}{\partial \nu} (1 - \bar{\xi}) ,$$

and in view of (3. 44) this implies

$$\begin{aligned} & \int_{B_t(x_0)} |\nabla \bar{\xi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi})^2 \\ & \leq C \left(\int_{\partial B_t(x_0)} (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} \left(\int_{\partial B_t(x_0)} \left| \frac{\partial \bar{\xi}}{\partial \nu} \right|^2 \right)^{1/2} . \end{aligned} \tag{3. 45}$$

Multiply (3. 43) by $\sum_{i=1}^3 x_i \frac{\partial(1-\bar{\xi})}{\partial x_i}$ and integrate on $B_t(x_0)$ we get the following Pohozaev identity

$$\begin{aligned} \frac{1}{2} t \int_{\partial B_t(x_0)} |\nabla_{\top} \bar{\xi}|^2 &= \frac{1}{2} t \int_{\partial B_t(x_0)} \left| \frac{\partial \bar{\xi}}{\partial \nu} \right|^2 + \frac{1}{2} \int_{B_t(x_0)} |\nabla \bar{\xi}|^2 + \\ &+ \frac{3}{2} \int_{B_t(x_0)} \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi})^2 - \frac{t}{2} \int_{\partial B_t(x_0)} \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi})^2 . \end{aligned}$$

The Pohozaev identity and (3. 45) imply

$$\int_{\partial B_t(x_0)} \left| \frac{\partial \bar{\xi}}{\partial \nu} \right|^2 \leq C \int_{\partial B_t(x_0)} |\nabla_{\top} |\bar{\varphi}||^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 \quad . \quad (3. 46)$$

(3. 46) combined with (3. 45), give

$$\begin{aligned} & \int_{\partial B_t(x_0)} |\nabla \bar{\xi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - \bar{\xi}^2)^2 \\ & \leq C \left(\int_{\partial B_t} (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} \left(\int_{\partial B_t} |\nabla_{\top} |\bar{\varphi}||^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} \quad (3. 47) \\ & \leq C \left(\int_{T_1} (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} \left(\int_{T_1} |\nabla \bar{\varphi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} . \end{aligned}$$

Let $\bar{W} = \bar{\xi} e^{i\bar{\psi}}$ in $B_t(x_0)$, since $(\bar{\varphi}, \hat{A})$ minimizes $\bar{G}_{\varepsilon, r_0}$, we have $\bar{G}_{\varepsilon, r_0}(\bar{\varphi}, \hat{A}) \leq \bar{G}_{\varepsilon, r_0}(\bar{W}, \hat{A})$ and this implies

$$\begin{aligned} & \int_{B_t(x_0)} |\nabla_{\hat{A}} \bar{\varphi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 + \frac{1}{r_0^2} |\bar{h}|^2 \\ & \leq \int_{B_t(x_0)} 2|\nabla \bar{\psi}|^2 + |\nabla \bar{\xi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\xi}|^2)^2 + 2|\hat{A}|^2 + \frac{1}{r_0^2} |\bar{h}|^2 \end{aligned} \quad (3. 48)$$

Combining (3. 39), (3. 41), (3. 47) and (3. 48) we have

$$\begin{aligned} & \int_{B_t(x_0)} |\nabla_{\hat{A}} \bar{\varphi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 (1 - |\bar{\varphi}|^2)^2 + \frac{1}{r_0^2} |\bar{h}|^2 \\ & \leq C \left[\left(\int_{T_1} |\nabla_{\top} \bar{\varphi}|^2 \right) \times \left(\int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right) + \int_{T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 + \int_{\partial T_1} \left| \frac{\partial \bar{\varphi}}{\partial \sigma} \right|^2 \right. \\ & \quad + \left(\int_{T_1} (1 - |\bar{\varphi}|^2)^2 \right)^{1/2} \left(\int_{T_1} |\nabla \bar{\varphi}|^2 + \left(\frac{r_0}{\varepsilon}\right)^2 + (1 - |\bar{\varphi}|^2)^2 \right) \\ & \quad \left. + \int_{T_1} |\nabla \hat{A}|^2 + \eta^2 + \int_{B_t(x_0)} \frac{1}{r_0^2} |\bar{h}|^2 \right] \end{aligned}$$

Going back to the usual scale we have

$$\begin{aligned}
 \frac{1}{r_0} E_{tr_0}(x_0) &\leq C \left[\frac{1}{r_0} \int_{B_{r_0}} |\nabla\varphi|^2 \times \frac{1}{r_0} \int_{B_{r_0}} \left| \frac{\partial\varphi}{\partial r} \right|^2 + \frac{1}{r_0} \int_{B_{r_0}} \left| \frac{\partial\varphi}{\partial r} \right|^2 \right. \\
 &+ \int_{\partial B_{r_0} \cup \partial B_{r_0/2}} \left| \frac{\partial\varphi}{\partial r} \right|^2 + r_0 \int_{B_{r_0}} |\nabla\tilde{A}|^2 \\
 &+ \left(\frac{1}{r_0^3} \int_{B_{r_0}} (1 - |\varphi|^2)^2 \right)^{1/2} \left(\frac{1}{r_0} \int_{B_{r_0}} |\nabla\varphi|^2 + \frac{1}{\varepsilon^2} (1 - |\varphi|^2)^2 \right)^2 \\
 &\left. + \eta^2 + \frac{1}{r_0} \int_{B_{tr_0}} |h|^2 \right] . \tag{3.49}
 \end{aligned}$$

Since $|\tilde{A}| \leq \frac{1}{d_{x_0}} + C\eta^{1/2}$ (see 3. 11), we have

$$\frac{1}{r_0} \int_{B_{r_0}(x_0)} |\tilde{A}|^2 \leq \left(\frac{r_0}{d_{x_0}} \right)^2 + C r_0^2 \eta^{1/2} \leq C\eta \quad . \tag{3.50}$$

Moreover $|\nabla\tilde{A}|^2 \leq 2|\nabla A_0|^2 + 2|\nabla A|^2$ and we have $|\nabla A_0|^2 \leq C/d_{x_0}^4$ and $\int_{\mathbb{R}^3} |\nabla A|^2 \leq C \log \frac{1}{\varepsilon}$, thus

$$r_0 \int_{B_{r_0}} |\nabla\tilde{A}|^2 \leq C \left(\frac{r_0^4}{d_{x_0}^4} + r_0 \log 1/\varepsilon \right) \leq C\eta \tag{3.51}$$

Finally combining (3. 48), (3. 49), (3. 50) and (3. 51) we have

$$\begin{aligned}
 \frac{1}{r_0} E_{tr_0}(x_0) &\leq \frac{E_{r_0}}{r_0} \times C\eta + C\eta + \left(\left(\frac{\varepsilon}{r_0} \right)^2 C\eta \right)^{1/2} \frac{E_{r_0}}{r_0} + \eta^2 \\
 &\leq C\eta^{1/2} \left(\frac{E_{r_0}}{r_0} + 1 \right) . \tag{3.52}
 \end{aligned}$$

In the other hand we have by (3. 1)

$$\frac{E_{r_0}}{r_0} - \frac{E_{tr_0}}{tr_0} = \int_{tr_0}^{r_0} \frac{J_r(x_0)}{r} + \int_{tr_0}^{r_0} R_r(x_0) \quad , \tag{3.53}$$

we have also

$$\int_{tr_0}^{r_0} \frac{J_r(x_0)}{r} \leq \frac{1}{tr_0} \int_{tr_0}^{r_0} J_r(x_0) \leq \frac{F_{r_0}(x_0)}{tr_0} \leq \frac{C}{t} \eta \tag{3.54}$$

and furthermore, since we could have chosen $r_0 \leq \frac{\eta}{\log 1/\varepsilon}$ we have

$$\begin{aligned}
 \int_{tr_0}^{r_0} R_r(x_0) &\leq C r_0 \log \frac{1}{\varepsilon} + C \frac{r_0^2}{d_{x_0}^2} \left(\frac{E_{r_0}(x_0)}{r_0} \right)^{1/2} \\
 &\leq C\eta + C\eta \left(\frac{E_{r_0}(x_0)}{r_0} \right)^{1/2} \tag{3.55}
 \end{aligned}$$

Combining (3. 52), (3. 53), (3. 54) and (3. 55) we easily get that for η sufficiently small, independently of r_0 , we have

$$\frac{E_{r_0}}{r_0}(x_0) \leq f(\eta) \longrightarrow 0 \quad \text{as} \quad \eta \rightarrow 0$$

This implies

$$\begin{aligned} \frac{E_\varepsilon}{\varepsilon}(x_0) &\leq f(\eta) + \int_\varepsilon^{r_0} R_r(x_0) \\ &\leq f(\eta) + Cr_0 \log \frac{1}{\varepsilon} + C\eta(f(\eta))^{1/2} \quad , \end{aligned}$$

and since we could have chosen $r_0 \leq c \frac{\eta}{\log 1/\varepsilon}$ we have

$$\frac{1}{\varepsilon^3} \int_{B_\varepsilon(x_0)} (1 - |\bar{\varphi}|^2)^2 \leq \frac{E_\varepsilon}{\varepsilon} \leq g(\eta) \longrightarrow 0 \quad \text{as} \quad \eta \rightarrow 0 \quad (3. 56)$$

and since $|\nabla|\varphi|| \leq |\nabla_{\tilde{A}}\varphi| \leq C/\varepsilon$ for η sufficiently small independent of ε (3. 56) ensures $|\varphi|(x_0) \geq \frac{1}{2}$. This proves the lemma. ■

4. LOCATING THE SINGULARITIES

As in [4], the bad set or the locus of the singularities is by definition the place where $|\varphi_\varepsilon| \leq \frac{1}{2}$ and where we are going to have a loss of compactness (as we will see later). The Lemma 3.1 says that a point x_0 where $|\varphi_\varepsilon|(x_0) < \frac{1}{2}$ produces energy in the ball $B_\rho(x_0)$ greater than $\rho \times \eta \log(\frac{\rho}{\varepsilon})$. In view of the global bound given by Lemma 2.1 we can easily obtain the three dimensional equivalent of Proposition 4.2 of [4], that is, we can put the bad set in $\frac{c(\alpha)}{\varepsilon^\alpha}$ 3-balls of radius ε^α for $\alpha < 1$, where $c(\alpha)$ only depends on α (this is equivalent to say, in dimension 2, that we can put the bad set in $c(\alpha)$ 2-balls of radius ε^α). Contrary to the dimension 2 (see [4], theorem 4.1 and [3] Lemma 4.1) we are not able to get the same result for $\alpha = 1$, that is to give the location of the singularities at the scale ε but as we will see it in part V this is not necessary in view of getting global $W^{1,p}$ estimates.

We have the following technical Lemma, part 4 is devoted to the proof of this lemma.

LEMMA 4.1. Given any $0 < \alpha < 1$ and $\alpha < \gamma < 1$ there exist $N_\alpha \in \mathbb{N}$ and $(R_i)_{1 \leq i \leq N_\alpha}$ a finite sequence of disjoint rectangle sets whose boundary are unions of faces parallel to the planes $x O y, y O z$ and $x O z$ such that

- 1 $-\exists, \mu > 0, \forall 1 \leq i \leq N_\alpha, \forall x \in R_i$ (resp. $\forall y \in \mathbb{R}^3 \setminus R_i$) there exists a cube of edge $\mu\varepsilon^\alpha$ included in R_i (resp. in $\mathbb{R}^3 \setminus R_i$) which contains x (resp. y) and R_i is contained in a cube of edge $2\varepsilon^\alpha$ (μ does not depend on ε). Those cubes have also faces parallel to the planes $x O y, y O z$ and $x O z$.
- 2 $-N_\alpha \leq \frac{c(\alpha)}{\varepsilon^\alpha}$, where $c(\alpha)$ only depends on α and $\{(x_i, d_i)\}_{i=1, \dots, n}$.
- 3 $-\forall 1 \leq i \leq N_\alpha \int_{\partial R_i} |\nabla_{\tilde{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 + |h|^2 \leq C \log \frac{1}{\varepsilon}$ where C only depends on $\{(x_l, d_l)\}_{l=1, \dots, n}$.

4 - $\forall 1 \leq i \leq N_\alpha$ there exists a subset S_i of ∂R_i (possibly empty) which can be contained in $n(\alpha, \gamma)$ balls of radius $\leq \varepsilon^\gamma$ such that

$$|\varphi| \geq \frac{1}{2} \quad \text{on } \partial R_i \setminus S_i$$

where $n(\alpha, \gamma)$ only depends on α, γ and $\{(x_i, d_i)\}_{i=1\dots n}$.

5 - for any $1 \leq i \leq N_\alpha$ such that $\forall k \in \{1, \dots, n\} x_k \notin R_i$, for any regular closed curve in $\partial R_i \setminus S_i$ and for any trivialisation of E over R_i , the degree of $\varphi/|\varphi|$ along this curve, which does not depend on the trivialisation chosen over R_i , is bounded by an integer $d(\alpha, \gamma)$ which only depends on α, γ and $\{(x_i, d_i)\}_{i=1\dots n}$.

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$$|\varphi| \geq \frac{1}{2} \quad \text{in } \mathbb{R}^3 \setminus \bigcup_{i=1}^{N_k} \bar{R}_i \cup_{k=1}^n C\left(\frac{\varepsilon^\alpha}{\eta}, x_k\right)$$

where η is given by Lemma 3.1 and $C\left(\frac{\varepsilon^\alpha}{\eta}, x_k\right)$ is the cube of centre x_k of edge $\frac{\varepsilon^\alpha}{\eta}$ parallel to Ox, Oy and Oz . ■

Proof of Lemma 4.1.

Consider the lattice $\varepsilon^\alpha \mathbb{Z}^3$ in \mathbb{R}^3 , and denote by $C(\rho, z)$ the cube of centre z , of edge ρ and whose faces are parallel to $x O y, y O z$ and $x O z$. We will only consider the points z_i of the lattice such that $C(2\varepsilon^\alpha, z_i)$ is not included in $\bigcup_{k=1}^n C\left(\frac{\varepsilon^\alpha}{\eta}, x_k\right)$. Let $J_\alpha \subset \varepsilon^\alpha \mathbb{Z}$ be this set. Let $(z_i)_{i \in I_\alpha}$ be those which verifies

$$\frac{1}{2\varepsilon^\alpha} \int_{C(2\varepsilon^\alpha, z_i)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 > \eta \log\left(\frac{2\varepsilon^\alpha}{\varepsilon}\right)$$

where η is the constant given by Lemma 3.1. We have

$$\sum_{i \in I_\alpha} \int_{C(2\varepsilon^\alpha, z_i)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 > \#I_\alpha 2\varepsilon^\alpha \eta (1 - \alpha) \log \frac{1}{\varepsilon} .$$

Since any point of \mathbb{R}^3 is at most covered by 4^3 of those cubes we have

$$\int_{\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 > \#I_\alpha C(1 - \alpha) \varepsilon^\alpha \log \frac{1}{\varepsilon} \tag{4. 1}$$

Combining (4. 1) and the global upper bound given by Lemma 2.1 we have

$$\#I_\alpha < \frac{C}{(1 - \alpha) \varepsilon^\alpha} . \tag{4. 2}$$

Moreover let $j \notin I_\alpha$, since

$$\frac{1}{2\varepsilon^\alpha} \int_{C(2\varepsilon^\alpha, z_j)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \leq \eta \log\left(\frac{\varepsilon^\alpha}{\varepsilon}\right) ,$$

by Lemma 3.1, $|\varphi| \geq 1/2$ in $C(\varepsilon^\alpha, z_j)$. In the other hand

$$\mathbb{R}^3 \setminus \bigcup_{i \in I_\alpha} C(\varepsilon^\alpha, z_i) \cup_{k=1}^n C\left(x_k, \frac{\varepsilon^\alpha}{\eta}\right) \subset \bigcup_{j \in J_\alpha \setminus I_\alpha} C(\varepsilon^\alpha, z_j)$$

thus

$$|\varphi| \geq \frac{1}{2} \quad \text{in } \mathbb{R}^3 \setminus \bigcup_{i \in I_\alpha} C(\varepsilon^\alpha, z_i) \cup_{k=1}^n C\left(x_k, \frac{\varepsilon^\alpha}{\eta}\right) \quad . \quad (4.3)$$

We change now the covering of the bad set, $\{C(2\varepsilon^\alpha, z_i)\}_{i \in I_\alpha}$ into a covering of cubes $\{C(\rho_i, z_i)\}_{i \in I_\alpha}$ with edges a little bit smaller than $2\varepsilon^\alpha$ (ie $(2 - \delta)\varepsilon^\alpha \leq \rho_i \leq 2\varepsilon^\alpha$, δ will be fixed later) in view of ensuring conditions 3, 4 and 5 of the lemma on the boundaries of the different connected components of $\mathbb{R}^3 \setminus \bigcup_{i \in I_\alpha} \partial C(\rho_i, z_i)$ which will be rectangles verifying 1.

Let $0 < \delta < 1$ chosen small enough at the end of the proof. Let $i \in I_\alpha$, from Lemma 2.5 there exists C independent of ε such that

$$\frac{1}{2\varepsilon^\alpha} \int_{C(2\varepsilon^\alpha, z_i)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \leq C \log \frac{1}{\varepsilon} \quad . \quad (4.4)$$

Let $\alpha < \beta < 1$ and consider the lattice of size ε^β (i.e. $\varepsilon^\beta \mathbb{Z}^3$) included in $C(2\varepsilon^\alpha, z_i)$. Denote it by $(y_j)_{j \in J_\alpha^\beta(z_i)}$. Denote by $(y_j)_{j \in I_\alpha^\beta(z_i)}$ the points of this lattice such that

$$\frac{1}{2\varepsilon^\beta} \int_{C(2\varepsilon^\beta, y_j)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 > \eta \log \left(\frac{2\varepsilon^\beta}{\varepsilon} \right) \quad .$$

We have

$$\sum_{j \in I_\alpha^\beta(z_i)} \int_{C(2\varepsilon^\beta, y_j)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 > \#I_\alpha^\beta(z_i) 2\varepsilon^\beta (1 - \beta) \times \log \frac{1}{\varepsilon}.$$

Since a point of $C(2\varepsilon^\alpha, z_i)$ can be covered by, at most, 4^3 cubes $C(2\varepsilon^\beta, y_j)$ for $j \in J_\alpha^\beta(z_i)$ we have

$$\begin{aligned} \int_{C(z_i, 2\varepsilon^\alpha)} |\nabla_{\tilde{A}} \varphi|^2 + \frac{1}{2\varepsilon^2} (1 - |\varphi|^2)^2 + |h|^2 \\ > C \#I_\alpha^\beta(z_i) 2\varepsilon^\beta (1 - \beta) \log \frac{1}{\varepsilon} \quad . \end{aligned} \quad (4.5)$$

Combining (4.4) and (4.5) we get

$$\#I_\alpha^\beta(z_i) < \frac{C}{(1 - \beta) \varepsilon^{\beta - \alpha}} \quad . \quad (4.6)$$

For $j \in J_\alpha^\beta(z_i) \setminus I_\alpha^\beta(z_i)$ we have, by Lemma 3.1,

$$|\varphi| \geq \frac{1}{2} \quad \text{in } C(\varepsilon^\beta, y_j) \quad . \quad (4.7)$$

Moreover

$$C(2\varepsilon^\alpha, z_i) \setminus \left(\bigcup_{j \in I_\alpha^\beta(z_i)} C(\varepsilon^\beta, y_j) \right) \subset \bigcup_{j \in J_\alpha^\beta(z_i) \setminus I_\alpha^\beta(z_i)} C(\varepsilon^\beta, y_j) \quad ,$$

thus

$$|\varphi| \geq \frac{1}{2} \quad \text{in } C(2\varepsilon^\alpha, z_i) \setminus \bigcup_{j \in I_\alpha^\beta(z_i)} C(\varepsilon^\beta, y_j) \quad . \quad (4.8)$$

Using the mean value formula we will choose $\rho \in ((2 - \delta)\varepsilon^\alpha; 2\varepsilon^\alpha)$ such that $\partial C(\rho, z_i)$ intersects a number of cubes $C(2\varepsilon^\beta, y_j)$, for $j \in I_\alpha^\beta(z_i)$, bounded independently of i, ε and such that

$$\int_{\partial C(\rho, z_i)} |\nabla_{\bar{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 + |h|^2 \leq K \log \frac{1}{\varepsilon} \quad , \quad (4. 9)$$

where K only depends on δ , and $\{(x_i, d_i)\}_{i=1, \dots, n}$.

Let $\mathcal{N}_\rho(z_i)$ be the number of indices $j \in I_\alpha^\beta(z_i)$ such that

$$C(2\varepsilon^\beta, y_j) \cap \partial C(\rho, z_i) \neq \emptyset \quad .$$

If $|\rho_1 - \rho_2| > 2\varepsilon^\beta \mathcal{N}_{\rho_1}(z_i)$ and $\mathcal{N}_{\rho_2}(z_i)$ count different indices thus, using (4. 6) we have

$$\begin{aligned} \int_{(2-\delta)\varepsilon^\alpha}^{2\varepsilon^\alpha} \mathcal{N}_\rho(z_i) \, d\rho &\leq \sum_{p=0}^{\left[\frac{\delta\varepsilon^\alpha}{2\varepsilon^\beta}\right]+1} \int_{2\varepsilon^\alpha-2(p+1)\varepsilon^\beta}^{2\varepsilon^\alpha-2p\varepsilon^\beta} \mathcal{N}_p(z_i) \, d\rho \\ &\leq 2\varepsilon^\beta \#I_\alpha^\beta(z_i) \leq \frac{C}{(1-\beta)}\varepsilon^\alpha \end{aligned} \quad (4. 10)$$

In the other hand, (4. 4) gives

$$\frac{\delta}{\varepsilon^\alpha} \int_{(2-\delta)\varepsilon^\alpha}^{2\varepsilon^\alpha} d\rho \int_{\partial C(\rho, z_i)} |\nabla_{\bar{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 + |h|^2 \leq K \log \frac{1}{\varepsilon}. \quad (4. 11)$$

Applying the mean value formula, simultaneously for (4. 10) and (4. 11) we deduce the existence of C independent of i and ε and the existence of $\rho_i \in ((2 - \delta)\varepsilon^\alpha, 2\varepsilon^\alpha)$ such that

$$\left\{ \begin{array}{l} \mathcal{N}_{\rho_i}(z_i) \leq C(\alpha, \beta, \delta) \\ \int_{\partial C(\rho_i, z_i)} |\nabla_{\bar{A}}\varphi|^2 + \frac{1}{2\varepsilon^2}(1 - |\varphi|^2)^2 + |h|^2 \leq C(\alpha, \beta, \delta) \log \frac{1}{\varepsilon} \end{array} \right. \quad (4. 12)$$

We claim now that we can put the bad set ($|\varphi| \leq 1/2$) on $\partial C(\rho_i, z_i)$ in a union of cubes of edges bounded by $\varepsilon^{\frac{\alpha+\beta}{2}}$ such that

1) the number of those cubes is still bounded by $\mathcal{N}_{\rho_i}(z_i) \leq C(\alpha, \beta, \delta)$

2) in any trivialisation of E over $C(2\varepsilon^\alpha, x_i)$, the degree of $\varphi/|\varphi|$ along the intersection of $\partial C(\rho_i, z_i)$ and the union of those cubes is bounded independently of ε .

Let $j \in I_\alpha^\beta(z_i)$ such that $C(2\varepsilon^\beta, y_j) \cap \partial C(\rho_i, z_i) \neq \emptyset$.

We claim that there exists $k \in [1; 2\mathcal{N}_{\rho_i}(z_i) + 2] \cap \mathbb{N}$ such that

$$|\varphi| \geq \frac{1}{2} \quad \text{in} \quad C\left(\varepsilon^{\beta-\frac{k}{2}\frac{\beta-\alpha}{(2\mathcal{N}_{\rho_i}+2)}}, y_j\right) \setminus C\left(\varepsilon^{\beta-\frac{(k-1)}{2}\frac{\beta-\alpha}{(2\mathcal{N}_{\rho_i}+2)}}, y_j\right) \cap \partial C(\rho_i, z_i) \quad . \quad (4. 13)$$

Indeed any cube $C(2\varepsilon^\beta, y_{j'})$ for $j' \in I_\alpha^\beta(z_i)$ intersects at most two of the following annuli (because of its size)

$$C\left(\varepsilon^{\beta-\frac{k}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}, y_j\right) \setminus C\left(\varepsilon^{\beta-\frac{(k-1)}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}, y_j\right) \text{ for } k \in \{1, 2, \dots, 2N_{\rho_i} + 2\} .$$

In fact, Since the number of the cubes $C(2\varepsilon^\beta, y_{j'})$ ($j' \in I_\alpha^\beta(z_i)$) which intersect $\partial C(\rho_i, x_i)$ is bounded by N_{ρ_i} there exists necessarily one of those annuli which intersects none of those cubes.

Let $k_0 \in \{1, 2, \dots, 2N_{\rho_i} + 2\}$ such that (4. 13) is verified.

For $r \in \left[\varepsilon^{\beta-\frac{(k_0-1)}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}, \varepsilon^{\beta-\frac{k_0}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}} \right]$, we denote by Γ_r the following lipshitz closed curve in $\partial C(x_i, \rho_i)$

$$\Gamma_r = \partial(C(r, y_j) \cap \partial C(\rho_i, z_i))$$

and Σ_r the surface in $\partial C(\rho_i, z_i)$ that Γ_r bounds

$$\Sigma_r = C(r, y_j) \cap \partial C(\rho_i, z_i).$$

We have

$$\begin{aligned} \oint_{\Gamma_r} \frac{1}{|\varphi|^2} (i\varphi, \nabla_{\tilde{A}}\varphi) &= \oint_{\Gamma_r} (i\frac{\varphi}{|\varphi|}, \nabla \frac{\varphi}{|\varphi|}) - \oint_{\Gamma_r} \tilde{A} \\ &= 2\pi d - \int_{\Sigma_r} *\tilde{h}.\nu \quad , \end{aligned} \tag{4. 14}$$

where d is the degree of $\varphi/|\varphi|$ on Γ_r which is independent of

$r \in \left[\varepsilon^{\beta-\frac{(k_0-1)}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}, \varepsilon^{\beta-\frac{k_0}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}} \right]$ and where ν is the outward unit normal of $\partial C(\rho_i, z_i)$.

Since $|\varphi| \geq \frac{1}{2}$ on Γ_r and $|\Gamma_r| \leq C r$, (4. 14) implies, using Cauchy-Schwartz inequality, that

$$\frac{(2\pi d)^2}{r} \leq C \left[\int_{\Gamma_r} |\nabla_{\tilde{A}}\varphi|^2 + \frac{1}{r} \left(\int_{\Sigma_r} |\tilde{h}| \right)^2 \right]. \tag{4. 15}$$

Moreover we have $|\tilde{h}| \leq |h_0| + |h|$ and

$$|h_0| \leq \frac{C}{\text{dist}^2\{y_j; \{x_i, \dots, x_n\}\}} \leq \frac{C}{\varepsilon^{2\alpha}} . \tag{4. 16}$$

In the other hand, by (2. 38), we have

$$\left(\int_{\Sigma_r} |h| \right)^2 \leq Cr^2 \int_{\Sigma_r} |h|^2 \leq Cr^3 \log \frac{1}{\varepsilon} . \tag{4. 17}$$

Combining (4. 15), (4. 16), (4. 17) and integrating with respect to r we have

$$\int_{\varepsilon^{\beta-\frac{(k_0-1)}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}}^{\varepsilon^{\beta-\frac{k_0}{2}\frac{\beta-\alpha}{(2N_{\rho_i}+2)}}} \frac{(2\pi d)^2 - C\varepsilon^{4(\beta-\alpha)}}{r} dr \leq C \int_{\partial C(\rho_i, z_i)} |\nabla_{\tilde{A}}\varphi|^2 \tag{4. 18}$$

In view of (4. 12) this implies

$$\frac{\beta - \alpha}{2} \times \frac{1}{2\mathcal{N}_{\rho_i} + 2} 4\pi^2(d^2 - C\varepsilon^{4(\beta-\alpha)}) \leq C(\alpha, \beta, \delta) \quad (4. 19)$$

Thus d is bounded independently of i and ε and we have put the bad set on $\partial C(\rho_i, z_i)$ in at most \mathcal{N}_{ρ_i} cubes of edge less than $\varepsilon^{\frac{\beta+\alpha}{2}}$ such that the degree of $\varphi/|\varphi|$ in any trivialisation of E over $C(\rho_i, x_i)$ on the boundary of the intersection of those cubes with $\partial C(\rho_i, z_i)$ is uniformly bounded independently of i and ε .

The connected components of $\mathbb{R}^3 \setminus \cup_{i \in I_\alpha} \partial C(\rho_i, z_i)$ are sets which boundary is a union of faces parallel to xOy , yOz and zOx and for which conditions 2, 3, 4, 5 and 6 are verified for $\gamma = \frac{\beta+\alpha}{2}$. But there is no reason for condition 1 to be verified by this sequence of connected components. Some small changes have to be done in the proof to get 1.

They appear here at the end of the proof in order not to overcomplicate the beginning.

Let $C(\rho, y)$ and $C(\rho', y')$ be two cubes with faces parallel to xOy , yOz and zOx and denote by

$$d(C(\rho, y); C(\rho', y'))$$

the minimum among all the distances between the faces of $\partial C(\rho, y)$ and the faces of $\partial C(\rho', y')$ which are parallel.

We are looking for a function $\lambda : \varepsilon^\alpha \mathbb{Z}^3 \rightarrow [2; 2 + \frac{1}{2}]$ such that

$$\forall z, z' \in \varepsilon^\alpha \mathbb{Z}^3 \quad y \neq y' \quad (4. 20)$$

$$d(C(\lambda(z)\varepsilon^\alpha, z); C(\lambda(z')\varepsilon^\alpha, z')) \geq \nu \varepsilon^\alpha$$

for some $\nu > 0$ independent of ε

Suppose such $\lambda : \varepsilon^\alpha \mathbb{Z}^3 \rightarrow [2, 2 + \frac{1}{2}]$ exists. Then in the previous proof, $\forall i \in I_\alpha$ choose ρ_i , by the mean value arguments developed above, between $\lambda(z_i)\varepsilon^\alpha$ and $(\lambda(z_i) - \delta)\varepsilon^\alpha$ (where δ is independent of i and ε) instead of choosing ρ_i between $2\varepsilon^\alpha$ and $(2 - \delta)\varepsilon^\alpha$. At the end of the proof, one comes exactly to the same conclusions, that is for $\mathbb{R}^3 \setminus \cup_{i \in I_\alpha} \partial C(\rho_i; z_i)$ the different connected components verify, 2, 3, 4, 5 and 6, moreover if one have chosen $\delta < \nu/4$ (where ν is given by (4. 20)) since, in taking $\partial C(\rho_i, z_i)$ instead of taking $\partial C(\lambda(z_i)\varepsilon^\alpha, z_i)$, smaller than $\frac{\nu}{4}\varepsilon^\alpha$ we have

$$\forall i, i' \in I_\alpha \quad i \neq i' \quad d(C(\rho_i, z_i); C(\rho_{i'}, z_{i'})) \geq (\nu - 2\frac{\nu}{4})\varepsilon^\alpha = \frac{\nu}{2}\varepsilon^\alpha \quad (4. 21)$$

and

$$\forall i \in I_\alpha, \quad \forall i' \notin I_\alpha \quad (4. 22)$$

$$d(C(\rho_i, z_i); C(\lambda(z_{i'})\varepsilon^\alpha, z_{i'})) \geq (\nu - \frac{\nu}{4})\varepsilon^\alpha = \frac{3\nu}{4}\varepsilon^\alpha.$$

(4. 21) and (4. 22) clearly ensure that the connected components of $\mathbb{R}^3 \setminus \cup_{i \in I_\alpha} \partial C(\rho_i, z_i)$ verify condition 1 for $\mu = \nu/2$ and the Lemma is proved.

Thus, we just have to prove that $\lambda \varepsilon^\alpha \mathbb{Z}^3 \rightarrow [2, 2 + \frac{1}{2}]$ verifying (4. 20) exists.

Take $\bar{\lambda} \varepsilon^\alpha \{1, 2, \dots, 5\}^3 \rightarrow [2; 2 + \frac{1}{4}]$ any injective function, and $\lambda(\varepsilon^\alpha(n_1, n_2, n_3)) = \bar{\lambda}(\varepsilon^\alpha(\bar{n}_1; \bar{n}_2; \bar{n}_3))$ where \bar{n}_l is the representative of n_l in $\mathbb{Z}/5\mathbb{Z}$ belonging to $\{1, 2, \dots, 5\}$.

We claim that (4. 20) holds for such λ

Letting $y, y' \in \varepsilon^\alpha \mathbb{Z}$, two situations may occur

first case : $d(C(2\varepsilon^\alpha, z); C(2\varepsilon^\alpha, z')) \geq \varepsilon^\alpha$

by replacing $2\varepsilon^\alpha$ by $\lambda\varepsilon^\alpha$, since $\lambda \in [2, 2 + \frac{1}{4}]$, we move each face of $C(2\varepsilon^\alpha, z)$ and $C(2\varepsilon^\alpha, z')$ parallelly along a distance less than $\frac{\varepsilon^\alpha}{4}$ thus

$$d(C(\lambda(z)\varepsilon^\alpha, z); C(\lambda(z')\varepsilon^\alpha, z')) \geq \left(1 - \frac{2}{4}\right)\varepsilon^\alpha = \frac{\varepsilon^\alpha}{2}$$

second case : $d(C(2\varepsilon^\alpha, y), C(2\varepsilon^\alpha, y')) = 0$

The parallel faces of $\overline{C(2\varepsilon^\alpha, z)}$ and $\overline{C(2\varepsilon^\alpha, z')}$ are either far from at least ε^α or touch themselves. For those which are far from at least ε^α , replacing $2\varepsilon^\alpha$ by $\lambda\varepsilon^\alpha$, as above, we keep the distance between themselves greater than $\varepsilon^\alpha/2$. For the parallel faces of $\overline{C(2\varepsilon^\alpha, z)}$ and $\overline{C(2\varepsilon^\alpha, z')}$ which touch themselves we have $\max_{l=1,2,3} |z_l - z'_l| = 2\varepsilon^\alpha$ and λ has been chosen such that, for points which are so close, $|\lambda(z) - \lambda(z')| \geq \nu$ where

$$\nu = \min\{|\bar{\lambda}(z) - \bar{\lambda}(z')|, z \neq z' \quad z, z' \in \varepsilon^\alpha \{1, 2, \dots, 5\}^3\} \quad .$$

Thus in replacing $2\varepsilon^\alpha$ by $\lambda\varepsilon^\alpha$ the two parallel faces become distant by at least $\nu\varepsilon^\alpha$. ■

5. $W^{1,p}$ ESTIMATES AND WEAK CONVERGENCE FOR THE MAGNETIC FIELD

This part is devoted to the proof of theorem 5.1 stated below. First, we establish $W^{1,p}$ estimates for the renormalised magnetic field $h_\varepsilon = \tilde{h}_\varepsilon - h_0$. Precisely we have the following lemma

LEMMA 5.1. Let $(\varphi_\varepsilon, \tilde{A}_\varepsilon)$ be a minimizer of G_ε among V and $h_\varepsilon = d\tilde{A}_\varepsilon - h_0$, we have

$$\forall 1 < p < \frac{3}{2} \quad \exists C_p \text{ indep. of } \varepsilon \quad \text{s.t.} \quad \int_{\mathbb{R}^3} |\nabla h_\varepsilon|^p \leq C_p$$

■

Proof of Lemma 5.1.

Let $0 < \alpha < 1$ and $\alpha < \gamma < 1$ be fixed later. Let $(R_i)_{1 \leq i \leq N_\alpha}$ be the union of disjoint rectangle sets given by Lemma 4.1 and denote T_{ε^α} the following set

$$T_{\varepsilon^\alpha} = \cup_{i=1}^{N_\alpha} \bar{R}_i \cup_{l=1}^n \bar{B}(x_l, \varepsilon^\alpha/\eta) \quad .$$

In $\mathbb{R}^3 \setminus T_{\varepsilon^\alpha}$ we have $|\varphi| \geq \frac{1}{2}$. Taking the exterior derivative of (2. 15) we have

$$-\Delta \tilde{h} + \tilde{h} = d \left[\left(1 - \frac{1}{|\varphi|^2}\right) (i\varphi, d_{\tilde{A}}\varphi) \right] \quad \text{in } \mathbb{R}^3 \setminus T_{\varepsilon^\alpha} \quad . \quad (5. 1)$$

$T_{\varepsilon\alpha}$ can be included in a big ball B_{R_ε} of finite radius $R_\varepsilon < +\infty$ and we have

$$\begin{aligned} \forall 1 < p < 2 \quad & \int_{\mathbb{R}^3 \setminus B_{R_\varepsilon}} \left(1 - \frac{1}{|\varphi|^2}\right)^p |(i\varphi, \nabla_{\tilde{A}}\varphi)|^p \leq \\ & \leq C \left[\int_{\mathbb{R}^3 \setminus B_{R_\varepsilon}} (1 - |\varphi|^2)^{\frac{2p}{2-p}} \right]^{1-\frac{p}{2}} \times \left[\int_{\mathbb{R}^3 \setminus B_{R_\varepsilon}} |\nabla_{\tilde{A}}\varphi|^2 \right]^{\frac{p}{2}} \quad (5.2) \\ & \leq C \left[\int_{\mathbb{R}^3} (1 - |\varphi|^2)^2 \right]^{1-\frac{p}{2}} \times \left[\int_{\mathbb{R}^3} |\nabla_{\tilde{A}}\varphi|^2 \right]^{p/2} < +\infty \end{aligned}$$

(5. 1) in $\mathbb{R}^3 \setminus B_{R_\varepsilon}$ and (5. 2) yields by classical estimates on Bessel potential that

$$\forall 1 < p \leq 2 \quad \int_{\mathbb{R}^3 \setminus B_{R_\varepsilon}} |\nabla \tilde{h}|^p < +\infty \quad . \quad (5.3)$$

Moreover one easily verifies that $\nabla h_0 \in L^p(\mathbb{R}^3 \setminus B_{R_\varepsilon})$ for any $p > 1$. Thus

$$\forall 1 < p \leq 2 \quad \int_{\mathbb{R}^3 \setminus B_{R_\varepsilon}} |\nabla h|^p < +\infty \quad , \quad (5.4)$$

and since, from Lemma II.3, we have $h \in W^{1,2}(\mathbb{R}^3)$ we have

$$\forall 1 < p \leq 2 \quad \int_{\mathbb{R}^3} |\nabla h|^p < +\infty \quad . \quad (5.5)$$

This implies, since $d^*\tilde{h} = d^*h$ in $\mathcal{D}'(\mathbb{R}^3)$,

$$\forall 1 < p \leq 2 \quad \int_{\mathbb{R}^3} |d^*\tilde{h}|^p < +\infty \quad . \quad (5.6)$$

We are going to establish a bound independent of ε for $\int_{\mathbb{R}^3} |d^*\tilde{h}|^p$ for $1 < p < \frac{3}{2}$.

Let $0 < \beta < 1$ be fixed later and K be the 1-form on \mathbb{R}^3 solution of

$$-\Delta K + K = \frac{d^*h}{|d^*h|^\beta} \quad \text{in } \mathcal{D}'(\mathbb{R}^3)$$

and $k = dK$ the two-form which verifies

$$-\Delta k + k = d \left(\frac{d^*h}{|d^*h|^\beta} \right) \quad \text{in } \mathcal{D}'(\mathbb{R}^3) \quad . \quad (5.7)$$

We have

$$\forall x \in \mathbb{R}^3 \quad K(x) = c \int_{\mathbb{R}^3} \frac{d^*h}{|d^*h|^\beta}(y) \frac{e^{-|x-y|}}{|x-y|} dy \quad . \quad (5.8)$$

By classical estimates on Bessel potential we have for any q such that $1 < q(1 - \beta) \leq 2$

$$\int_{\mathbb{R}^3} |\nabla k|^q \leq C_q \int_{\mathbb{R}^3} |d^*h|^{q(1-\beta)} \quad (5.9)$$

We will only work with $q > 3$. Let $\delta = 1 - 3/q$. Sobolev injections give

$$\forall x, y \in \mathbb{R}^3 \quad |k(x) - k(y)| \leq C_q \left(\int_{\mathbb{R}^3} |\nabla k|^q \right)^{1/q} |x - y|^\delta \quad . \quad (5.10)$$

Moreover

$$\begin{aligned} \forall x \in \mathbb{R}^3 \quad |k(x)| &= |dK(x)| \\ &\leq C \int_{\mathbb{R}^3} \frac{e^{-|x-y|}}{|x-y|} \left(\frac{1}{|x-y|} + 1 \right) |d^*h|^{1-\beta}(y) dy \end{aligned}$$

and since $q > 3$, $q' = \frac{q}{q-1} < \frac{3}{2}$ and $\forall x \in \mathbb{R}^3$ we have, $\forall x \in \mathbb{R}^3$,

$$\begin{aligned} |k(x)| &\leq \left(\int_{\mathbb{R}^3} \frac{e^{-q'|x-y|}}{|x-y|^{q'}} \left(\frac{1}{|x-y|^{q'}} + 1 \right) \right)^{\frac{1}{q'}} \times \left(\int_{\mathbb{R}^3} \frac{|d^*h|^{q(1-\beta)}}{(1+|d^*h|^\beta)^q} \right)^{1/q} \\ &\leq C_q \left(\int_{\mathbb{R}^3} |d^*h|^{q(1-\beta)} \right)^{1/q} \end{aligned} \tag{5.11}$$

Multiply (5. 1) by k and integrate on $\Omega_{\varepsilon\alpha} = \mathbb{R}^3 \setminus T_{\varepsilon\alpha}$ we have

$$\int_{\Omega_{\varepsilon\alpha}} -(k.dd^*\tilde{h} + k.d^*d\tilde{h}) + k.\tilde{h} = \int_{\Omega_{\varepsilon\alpha}} k.d \left[\left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, d_{\tilde{A}}\varphi) \right] .$$

Integrating by parts on both sides of the equality and using the fact that $d\tilde{h} \equiv 0$ in $\Omega_{\varepsilon\alpha}$ we have

$$\begin{aligned} & - \int_{\partial T_{\varepsilon\alpha}} k.d^*\tilde{h} \wedge \nu + \int_{\Omega_{\varepsilon\alpha}} d^*k.d^*\tilde{h} + k.d^*d\tilde{h} = \\ & - \int_{\Omega_{\varepsilon\alpha}} d^*k. \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, d_{\tilde{A}}\varphi) + \int_{\partial T_{\varepsilon\alpha}} k. \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, d_{\tilde{A}}\varphi) \wedge \nu \end{aligned} \tag{5.12}$$

where ν is the 1-form associated to the inward unit normal of $\partial T_{\varepsilon\alpha}$.

Multiplying (5. 7) by \tilde{h} and integrating on all of \mathbb{R}^3 we have

$$\int_{\mathbb{R}^3} -(\tilde{h}.dd^*k + \tilde{h}.d^*dk) + k.\tilde{h} = \int_{\mathbb{R}^3} \tilde{h}.d \left(\frac{d^*\tilde{h}}{|d^*\tilde{h}|^\beta} \right) .$$

Integrating by parts on both sides of the equality and using the fact that $dk \equiv 0$ in \mathbb{R}^3 we have

$$\int_{\mathbb{R}^3} d^*\tilde{h}.d^*k + \tilde{h}.k = - \int_{\mathbb{R}^3} |d^*\tilde{h}|^{2-\beta} . \tag{5.13}$$

Combining (5. 12) and (5. 13) we have

$$\begin{aligned} \int_{\mathbb{R}^3} |d^*\tilde{h}|^{2-\beta} &= - \int_{\partial T_{\varepsilon\alpha}} k.d^*\tilde{h} \wedge \nu - \int_{T_{\varepsilon\alpha}} d^*\tilde{h}.d^*k + \tilde{h}.k \\ & - \int_{\partial T_{\varepsilon\alpha}} k. \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, d_{\tilde{A}}\varphi) \wedge \nu + \\ & + \int_{\Omega_{\varepsilon\alpha}} d^*k. \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, d_{\tilde{A}}\varphi), \end{aligned} \tag{5.14}$$

We are going to establish some upper bounds for the four terms of the right hand side of (5. 14).

a) A bound for the fourth term of the r.h.s. of (5. 14).

$$\begin{aligned}
 & \left| \int_{\Omega_{\varepsilon\alpha}} d^*k \cdot \left(1 - \frac{1}{|\varphi|^2}\right) (i\varphi, d_{\tilde{A}}\varphi) \right| \\
 & \leq C \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \times \left(\int_{\mathbb{R}^3} (1 - |\varphi|^2)^{q'} |\nabla_{\tilde{A}}\varphi|^{q'} \right)^{1/q'} \\
 & \leq C \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \times \left[\int_{\mathbb{R}^3} (1 - |\varphi|^2)^2 \right]^{\frac{1}{q'} - \frac{1}{2}} \left[\int_{\mathbb{R}^3} |\nabla_{\tilde{A}}\varphi|^2 \right]^{\frac{1}{2}} \tag{5. 15} \\
 & \leq C\varepsilon^{\frac{2}{q'}-1} \left(\log \frac{1}{\varepsilon} \right)^{\frac{1}{q'}} \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q}
 \end{aligned}$$

where $\frac{1}{q'} = 1 - \frac{1}{q}$.

b) A bound for the third term of the r.h.s. of (5. 14).

$$\left| \int_{\partial T_{\varepsilon\alpha}} k \cdot \left(1 - \frac{1}{|\varphi|^2}\right) (i\varphi, d_{\tilde{A}}\varphi) \wedge \nu \right| \leq \|k\|_{L^\infty(\mathbb{R}^3)} \times \sum_{i=1}^{N_{\alpha+n}} \int_{\partial R_i} |\nabla_{\tilde{A}}\varphi|(1 - |\varphi|^2)$$

(recall $R_{N_{\alpha+l}} = C(\varepsilon^\alpha/\eta, x_l)$ for $1 \leq l \leq n$).

We bound $\int_{\partial R_i} |\nabla_{\tilde{A}}\varphi|(1 - |\varphi|^2)$ in the following way

$$\begin{aligned}
 \int_{\partial R_i} |\nabla_{\tilde{A}}\varphi|(1 - |\varphi|^2) & \leq \left(\int_{\partial R_i} |\nabla_{\tilde{A}}\varphi|^2 \right)^{1/2} \left(\int_{\partial R_i} (1 - |\varphi|^2)^2 \right)^{1/2} \\
 & \leq C\varepsilon \log \frac{1}{\varepsilon}
 \end{aligned}$$

and since $N_{\alpha} + n \leq \frac{C(\alpha)}{\varepsilon^\alpha}$ we have, using also (5. 11),

$$\left| \int_{\partial T_{\varepsilon\alpha}} k \cdot \left(1 - \frac{1}{|\varphi|^2}\right) (i\varphi, d_{\tilde{A}}\varphi) \right| \leq C\varepsilon^{1-\alpha} \log \frac{1}{\varepsilon} \left(\int_{\mathbb{R}^3} |d^*h|^{q(1-\beta)} \right)^{1/q} \tag{5. 16}$$

c) A bound for the second term of the r.h.s. of (5. 14).

We have $d^*\tilde{h} = d^*h$, thus

$$\begin{aligned}
 \left| \int_{T_{\varepsilon\alpha}} d^*h \cdot d^*k \right| & \leq \left(\int_{T_{\varepsilon\alpha}} |d^*h|^2 \right)^{1/2} \left(\int_{T_{\varepsilon\alpha}} |d^*k|^2 \right)^{1/2} \\
 & \leq \left(\int_{T_{\varepsilon\alpha}} |\nabla_{\tilde{A}}\varphi|^2 \right)^{1/2} |T_{\varepsilon\alpha}|^{1-\frac{2}{q}} \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \tag{5. 17} \\
 & \leq \varepsilon^{\alpha(2-\frac{4}{q})} \left(\log \frac{1}{\varepsilon} \right)^{1/2} \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q}.
 \end{aligned}$$

On the other hand $\int_{T_{\varepsilon^\alpha}} \tilde{h}.k = \int_{T_{\varepsilon^\alpha}} h_0.k + \int_{T_{\varepsilon^\alpha}} h.k$ and we have

$$\left| \int_{T_{\varepsilon^\alpha}} h_0.k \right| \leq \|k\|_\infty \left| \int_{T_{\varepsilon^\alpha}} |h_0| \right| \quad . \tag{5. 18}$$

Let $R < +\infty$ such that $\forall l \leq n \quad x_l \in B_R(0)$. Clearly $h_0 \in L^\infty(\mathbb{R}^3 \setminus B_r(0))$, moreover $h_0 \in L^r(B_R(0))$ for $1 \leq r < \frac{3}{2}$. Thus

$$\int_{T_{\varepsilon^\alpha}} |h_0| = \int_{T_{\varepsilon^\alpha} \cap B_R(0)} |h_0| + \int_{T_{\varepsilon^\alpha} \cap [\mathbb{R}^3 \setminus B_R(0)]} |h_0|$$

and we have

$$\int_{T_{\varepsilon^\alpha}} |h_0| \leq |T_{\varepsilon^\alpha}| \|h_0\|_{L^\infty(\mathbb{R}^3 \setminus B_R(0))} + |T_{\varepsilon^\alpha}|^{1-\frac{1}{r}} \|h_0\|_{L^r(B_R(0))} \quad .$$

Combining (5. 11), (5. 18) and the previous inequality we get for any $1 < r < \frac{3}{2}$

$$\left| \int_{T_{\varepsilon^\alpha}} h_0.k \right| \leq C\varepsilon^{2\alpha(1-\frac{1}{r})} \left(\int_{\mathbb{R}^3} |d^*h|^q \right)^{1/q} \quad . \tag{5. 19}$$

Moreover, using also (5. 11),

$$\begin{aligned} \left| \int_{T_{\varepsilon^\alpha}} h.k \right| &\leq \|k\|_\infty |T_{\varepsilon^\alpha}|^{1/2} \left(\int_{T_{\varepsilon^\alpha}} h^2 \right)^{1/2} \\ &\leq C\varepsilon^\alpha \left(\log \frac{1}{\varepsilon} \right)^{1/2} \left(\int_{\mathbb{R}^3} |d^*h|^q \right)^{1/q} \end{aligned} \tag{5. 20}$$

d) A bound for the first term of the r.h.s. of (5. 14).

If ν denotes the inward normal for each $(R_i)_{i=1, \dots, N_\alpha+n}$ we have by cancellation of ν and $-\nu$ on $\partial R_i \cap \partial R_j$ ($i \neq j$)

$$\begin{aligned} \int_{\partial T_{\varepsilon^\alpha}} k.d^*\tilde{h} \wedge \nu &= - \int_{\partial(\cup_{i=1}^{N_\alpha+n} R_i)} k.(i\varphi, d_{\tilde{A}}\varphi) \wedge \nu \\ &= \sum_{i=1}^{N_\alpha+n} - \int_{\partial R_i} k.(i\varphi, d_{\tilde{A}}\varphi) \wedge \nu \end{aligned} \tag{5. 21}$$

Denote by \bar{k}_i for $i = 1, \dots, N_\alpha + n$ the mean value of k on R_i and by \bar{k} the following approximation of k in T_{ε^α}

$$\bar{k} = \sum_{i=1}^{N_\alpha+n} \bar{k}_i \chi(R_i)$$

where $\chi(R_i)$ is the characteristic function of R_i equal to 1 in R_i and 0 elsewhere. Remark that from (5. 10) we deduce that

$$\|k - \bar{k}\|_{L^\infty(T_{\varepsilon^\alpha})} \leq C_q \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \varepsilon^{\alpha\delta} \tag{5. 22}$$

where $\delta = 1 - \frac{3}{q}$. Moreover we have by (5. 21)

$$\begin{aligned} \int_{\partial T_{\varepsilon\alpha}} k.d^*\bar{h} \wedge \nu &= \sum_{i=1}^{N_\alpha+n} - \int_{\partial R_i} k.(i\varphi, d_{\bar{A}}\varphi) \wedge \nu \\ &= \sum_{i=1}^{N_\alpha+n} - \int_{\partial R_i} \bar{k}.(i\varphi, d_{\bar{A}}\varphi) \wedge \nu + \int_{\partial R_i} (\bar{k} - k).(i\varphi, d_{\bar{A}}\varphi) \wedge \nu \end{aligned} \tag{5. 23}$$

The second term in the right hand side of (5. 23) can be bounded in the following way, using (5. 22) and condition 3 of Lemma 4.1,

$$\begin{aligned} \left| \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} (\bar{k} - k)(i\varphi, d_{\bar{A}}\varphi) \wedge \nu \right| &\leq \|k - \bar{k}\|_{L^\infty(T_{\varepsilon\alpha})} \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} |\nabla_{\bar{A}}\varphi| \\ &\leq C_q \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \varepsilon^{\alpha\delta} \sum_{i=1}^{N_\alpha+n} |\partial R_i|^{1/2} \left(\log \frac{1}{\varepsilon} \right)^{1/2} . \end{aligned}$$

Since for $i = 1, \dots, N_\alpha + n$ we have $|\partial R_i| \leq C\varepsilon^{2\alpha}$ and $N_\alpha + n \leq C(\alpha)/\varepsilon^\alpha$, the previous inequality implies

$$\left| \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} (\bar{k} - k)(i\varphi, d_{\bar{A}}\varphi) \wedge \nu \right| \leq C_q \left(\log \frac{1}{\varepsilon} \right)^{1/2} \varepsilon^{\alpha\delta} \left(\int_{\mathbb{R}^3} |d^*k|^q \right)^{1/q} \tag{5. 24}$$

We establish now an upper bound for the first term of the right hand side of (5. 23). Consider now just the R_i which do not contain any of the $(x_l)_{l=1, \dots, n}$ since there are only n which contain the (x_l) the bound for them in (5. 23) is easy to establish.

We decompose the vector associated to the 1-form $\star\bar{k}_i$ the canonical basis (e_1, e_2, e_3) of \mathbb{R}^3 , that is $\bar{k}_i^1 e_1 + \bar{k}_i^2 e_2 + \bar{k}_i^3 e_3$. Thus we have

$$\int_{\partial R_i} \bar{k}.(i\varphi, d_{\bar{A}}\varphi) \wedge \nu = \sum_{l=1}^3 \bar{k}_i^l \int_{\partial R_i} (i\varphi, \nabla_{\bar{A}}\varphi).\nu \wedge e_l \tag{5. 25}$$

Since R_i is a set with faces perpendicular to e_1, e_2 , and $e_3, \nu \wedge e_1, \nu \wedge e_2$ and $\nu \wedge e_3$ are unit tangent fields on ∂R_i whose integral curves are the intersection of ∂R_i with planes respectively perpendicular to e_1, e_2 and e_3 . $\int_{\partial R_i} (i\varphi, \nabla_{\bar{A}}\varphi).\nu \wedge e_l$ for $l = 1, \dots, 3$ is the integral over all the planes perpendicular to e_l intersecting ∂R_i of the circulation of the current $(i\varphi, \nabla_{\bar{A}}\varphi)$ along the integral curves of $\nu \wedge e_l$ that is the curves realized by the intersection of those planes perpendicular to e_l with ∂R_i . Let \hat{R}_i be the convex envelope of R_i . Write $\hat{R}_i = \Pi_{l=1}^3 [r_l^i, s_l^i]$. Let $\Gamma_l(t)$ be the following curve

$$\Gamma_l(t) = \partial R_i \cap \{x \in \mathbb{R}^3 / x.e_l = t\} \quad .$$

$\Gamma_l(t)$ is oriented by e_l (i.e. in the plane $\{x \in \mathbb{R}^3 / x.e_l = t\}$ oriented by $e_l \Gamma_l(t)$ has the positive orientation). We have

$$\int_{\partial R_i} (i\varphi, \nabla_{\bar{A}}\varphi).\nu \wedge e_l = - \int_{t=r_l^i}^{s_l^i} \oint_{\Gamma_l(t)} (i\varphi, \nabla_{\bar{A}}\varphi) \tag{5. 26}$$

For $i = 1, \dots, N_\alpha$ and $l = 1, 2, 3$ denote $I_l^i = [r_l^i, s_l^i]$ and let \tilde{I}_l^i be the subset of $t \in I_l^i$ such that $\Gamma_l(t)$ does not intersect the balls, given by Lemma 4.1, of radius bounded by ε^γ which contain the bad set on ∂R_i . Since the number of those balls is bounded by $n(\alpha, \gamma)$ which does not depend neither on i nor on ε , the measure of the surface $\Sigma_l^i = \cup_{t \notin \tilde{I}_l^i} \Gamma_l(t)$ is bounded in the following way

$$|\Sigma_l^i| \leq C \varepsilon^\gamma \varepsilon^\alpha \quad . \tag{5. 27}$$

Using the definition of Σ_l^i we have

$$\int_{I_l^i} \oint_{\Gamma_l(t)} (i\varphi, \nabla_{\tilde{A}} \varphi) = \int_{\tilde{I}_l^i} \oint_{\Gamma_l(t)} (i\varphi, \nabla_{\tilde{A}} \varphi) + \int_{\Sigma_l^i} (i\varphi, \nabla_{\tilde{A}} \varphi) \cdot \nu \wedge e_l. \tag{5. 28}$$

Using this decomposition of ∂R_i we bound the first term of the right hand side of (5. 23) in the following way. First of all we have

$$\begin{aligned} & \left| \sum_{i=1}^{N_\alpha} \bar{k}_i^i \int_{\Sigma_l^i} (i\varphi, \nabla_{\tilde{A}} \varphi) \cdot \nu \wedge e_l \right| \\ & \leq C \|k\|_{L^\infty} \left[\sum_{i=1}^{N_\alpha+n} |\Sigma_l^i|^{1/2} \left(\int_{\partial R_i} |\nabla_{\tilde{A}} \varphi|^2 \right)^{1/2} \right] \\ & \leq C \left[\int_{\mathbb{R}^3} |d^* \tilde{h}|^{q(1-\beta)} \right]^{1/q} \frac{1}{\varepsilon^\alpha} \varepsilon^{\frac{\alpha}{2} + \frac{\gamma}{2}} \left(\log \frac{1}{\varepsilon} \right)^{1/2} \\ & \leq C \varepsilon^{\frac{\gamma-\alpha}{2}} \left(\log \frac{1}{\varepsilon} \right)^{1/2} \left[\int_{\mathbb{R}^3} |d^* \tilde{h}|^{q(1-\beta)} \right]^{1/q}. \end{aligned} \tag{5. 29}$$

On the good set on ∂R_i we write

$$\begin{aligned} \int_{\tilde{I}_l^i} \oint_{\Gamma_l(t)} (i\varphi, \nabla_{\tilde{A}} \varphi) &= \int_{\tilde{I}_l^i} \oint_{\Gamma_l(t)} \left(i \frac{\varphi}{|\varphi|}, \nabla_{\tilde{A}} \frac{\varphi}{|\varphi|} \right) + \\ &+ \int_{\tilde{I}_l^i} \oint_{\Gamma_l(t)} \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, \nabla_{\tilde{A}} \varphi) \end{aligned} \tag{5. 30}$$

and we have

$$\begin{aligned} & \left| \sum_{i=1}^{N_\alpha} \bar{k}_i^i \int_{\tilde{I}_l^i} \oint_{\Gamma_l(t)} \left(1 - \frac{1}{|\varphi|^2} \right) (i\varphi, \nabla_{\tilde{A}} \varphi) \right| \\ & \leq C \|k\|_{L^\infty} \left[\sum_{i=1}^{N_\alpha} \left(\int_{\partial R_i} (1 - |\varphi|^2)^2 \right)^{1/2} \left(\int_{\partial R_i} |\nabla_{\tilde{A}} \varphi|^2 \right)^{1/2} \right] \\ & \leq C \left[\int_{\mathbb{R}^3} |d^* h|^{q(1-\beta)} \right]^{1/q} \frac{1}{\varepsilon^\alpha} \varepsilon \log \frac{1}{\varepsilon}. \end{aligned} \tag{5. 31}$$

Finally we have, for any trivialisation of E over R_i ,

$$\begin{aligned} \int_{\tilde{I}_i} \oint_{\Gamma_l(t)} \left(i \frac{\varphi}{|\varphi|}, \nabla_{\tilde{A}} \frac{\varphi}{|\varphi|} \right) &= \int_{\tilde{I}_i} \oint_{\Gamma_l(t)} \left(i \frac{\varphi}{|\varphi|}, \nabla \frac{\varphi}{|\varphi|} \right) - \int_{\tilde{I}_i} \oint_{\Gamma_l(t)} \tilde{A} \\ &= \int_{\tilde{I}_i} 2\pi d_l(t) + \int_{\tilde{I}_i} \int_{S_i^l(t)} * \tilde{h} \cdot \nu \end{aligned} \tag{5.32}$$

where $d_l(t)$ is the degree of $\varphi/|\varphi|$ on $\Gamma_l(t)$ (which is bounded independently of i and ε by condition 5 of Lemma 4.1) and $S_i^l(t)$ is the surface in ∂R_i which bounds $\Gamma_l(t)$ (taking into account the orientation of $\Gamma_l(t)$). ∂R_i is of course oriented by the outward unit normal.

$$\begin{aligned} \left| \int_{S_i^l(t)} * \tilde{h} \cdot \nu \right| &\leq \int_{\partial R_i} |h_0| + \int_{\partial R_i} |h| \\ &\leq C \left[\int_{\partial R_i} \frac{1}{\varepsilon^{2\alpha}} + |\partial R_i|^{1/2} \left(\log \frac{1}{\varepsilon} \right)^{1/2} \right] \leq C \end{aligned} \tag{5.33}$$

Thus combining (5.32) and (5.33) we have

$$\begin{aligned} \left| \sum_{i=1}^{N_\alpha} \tilde{k}_i^l \int_{\tilde{I}_i} \oint_{\Gamma_l(t)} \left(i \frac{\varphi}{|\varphi|}; \nabla_{\tilde{A}} \frac{\varphi}{|\varphi|} \right) \right| &\leq C \|k\|_{L^\infty} \left[\sum_{i=1}^{N_\alpha} |\tilde{I}_i^l| \right] \\ &\leq \left(\int_{\mathbb{R}^3} |d^* \tilde{h}|^{q(1-\beta)} \right)^{1/q} \frac{1}{\varepsilon^\alpha} \times \varepsilon^\alpha \end{aligned} \tag{5.34}$$

Combining (5.23), (5.24), (5.25), (5.29), (5.31) and (5.34) we finally get an upper bound for the first term in the right hand side of (5.14), namely

$$\begin{aligned} \left| \int_{\partial T_{\varepsilon\alpha}} k \cdot d^* h \wedge \nu \right| &\leq \left(1 + \varepsilon^{1-\alpha} \log \frac{1}{\varepsilon} + \varepsilon^{\gamma-\alpha} \left(\log \frac{1}{\varepsilon} \right)^{1/2} + \right. \\ &\quad \left. + \varepsilon^{\alpha\delta} \left(\log \frac{1}{\varepsilon} \right)^{1/2} \right) \times \left(\int_{\mathbb{R}^3} |d^* \tilde{h}|^{q(1-\beta)} \right)^{\frac{1}{q}} \end{aligned} \tag{5.35}$$

Combining (5.14), (5.15), (5.19), (5.20) and (5.35) we obtain for $\gamma > \alpha$, $q > 3$ and ε sufficiently small

$$\int_{\mathbb{R}^3} |d^* \tilde{h}|^{2-\beta} \leq C(q, \gamma, \alpha) \left(\int_{\mathbb{R}^3} |d^* \tilde{h}|^{q(1-\beta)} \right)^{1/q} \tag{5.36}$$

where C does not depend on ε .

Choose β such that $2 - \beta = q(1 - \beta)$. Since (5.36) holds for any $q > 3$, this holds for any $1 > \beta > \frac{1}{2}$. Thus for any $1 < p < \frac{3}{2}$ we have

$$\int_{\mathbb{R}^3} |d^* \tilde{h}|^p \leq C(p, \gamma, \alpha) \left(\int_{\mathbb{R}^3} |d^* \tilde{h}|^p \right)^{\frac{p-1}{p}}, \tag{5.37}$$

and we conclude that $\|d^*h\|_{L^p(\mathbb{R}^3)}$ is bounded independently of ε for $1 < p < \frac{3}{2}$. By definition $A = -\frac{1}{4\pi} \frac{1}{|x|} * d^*h$, thus by classical estimates on the Riesz transform, we have

$$\int_{\mathbb{R}^3} |\nabla h|^p \leq \int_{\mathbb{R}^3} |\nabla^2 A|^p \leq C_p \int_{\mathbb{R}^3} |d^*h|^p \quad .$$

This proves the Lemma. ■

From Lemma 5.1 we deduce the following useful corollary

COROLLARY 5.1. Let $(\varphi_\varepsilon, \tilde{A}_\varepsilon)$ be a minimizer of G_ε among V . Then we have

$$\forall \Omega \subset \mathbb{R}^3 \text{ s. t. } |\Omega| < +\infty \quad \forall 1 < p < \frac{3}{2} \tag{5.38}$$

$$\int_{\Omega} |\nabla_{\tilde{A}} \varphi|^p \leq C(p, |\Omega|) \quad .$$

■

Proof of corollary 5.1

Let $0 < \alpha < 1$ and T_{ε^α} be the set introduced in the proof of Lemma 5.1. Since $|\varphi| \geq 1/2$ in $\mathbb{R}^3 \setminus T_{\varepsilon^\alpha}$ we have

$$|\nabla_{\tilde{A}} \varphi| \leq C(|(i\varphi, \nabla_{\tilde{A}} \varphi)| + |\nabla|\varphi||) \tag{5.39}$$

From Lemma 5.1 and equation (2. 15) we deduce

$$\forall 1 < p < 3/2 \quad \int_{\mathbb{R}^3} |(i\varphi, \nabla_{\tilde{A}} \varphi)|^p \leq C(p) \quad . \tag{5.40}$$

On the other hand for any $p \in (1; 3/2)$

$$\begin{aligned} \int_{T_{\varepsilon^\alpha}} |\nabla_{\tilde{A}} \varphi|^p &\leq |T_{\varepsilon^\alpha}|^{1-\frac{p}{2}} \left(\int_{T_{\varepsilon^\alpha}} |\nabla_{\tilde{A}} \varphi|^2 \right)^{\frac{p}{2}} \\ &\leq C\varepsilon^{\alpha(2-p)} \left(\log \frac{1}{\varepsilon} \right)^{\frac{p}{2}} \rightarrow 0 \text{ as } \varepsilon \rightarrow 0 \end{aligned} \tag{5.41}$$

We bound $\int_{\Omega} |\nabla|\varphi||^p$ using a method developed in [3] (lemma X.13).

Recall equation (2. 14) verified by $|\varphi|^2$ in $\mathbb{R}^3 \setminus \{x_1, \dots, x_n\}$

$$\frac{1}{2} \Delta |\varphi|^2 = -\frac{|\varphi|^2}{\varepsilon^2} (1 - |\varphi|^2) + |\nabla_{\tilde{A}} \varphi|^2 \tag{5.42}$$

Let $\overline{|\varphi|} = \max(|\varphi|; 1 - \frac{1}{\log \frac{1}{\varepsilon}})$ and let

$$K = \left\{ x \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\} \text{ s. t } |\varphi| \geq 1 - \frac{1}{Ln \frac{1}{\varepsilon}} \right\} .$$

Multiply (5. 42) by $1 - \overline{|\varphi|}$ and integrate on $K \setminus \cup_{i=1}^n B(x_i, \varepsilon)$ we have

$$\begin{aligned} \frac{1}{2} \int_{K \setminus \cup_{i=1}^n B(x_i, \varepsilon)} |\nabla|\varphi||^2 &\leq \int_{K \setminus \cup_{i=1}^n B(x_i, \varepsilon)} |\nabla_{\tilde{A}}\varphi|^2 (1 - \overline{|\varphi|}) \\ &+ \sum_{i=1}^n \int_{\partial B(x_i, \varepsilon)} |\nabla|\varphi|| (1 - \overline{|\varphi|}) \quad (5. 43) \\ &\leq C \frac{\log 1/\varepsilon}{\log 1/\varepsilon} + n \frac{C}{\varepsilon} \frac{\varepsilon^2}{\log 1/\varepsilon} \leq C . \end{aligned}$$

On the other hand

$$|\mathbb{R}^3 \setminus K| \times \frac{1}{(\log \frac{1}{\varepsilon})^2} \leq \int_{\mathbb{R}^3} (1 - |\varphi|)^2 \leq C \varepsilon^2 \log \frac{1}{\varepsilon} . \quad (5. 44)$$

Let $p < 2$, we have

$$\begin{aligned} \int_{\mathbb{R}^3 \setminus K} |\nabla|\varphi||^p &\leq |\mathbb{R}^3 \setminus K|^{1-\frac{p}{2}} \times \left(\int_{\mathbb{R}^3} |\nabla_{\tilde{A}}\varphi|^2 \right)^{\frac{p}{2}} \\ &\leq C \varepsilon^{2-p} \log \frac{1}{\varepsilon} \quad (5. 45) \end{aligned}$$

The combination of (5. 39), (5. 40), (5. 41), (5. 43) and (5. 45) implies easily the desired result. ■

We are now in position to state and prove the following theorem which was the aim of the section 5.

THEOREM 5.1. Let ε_n be a sequence of positive numbers tending to 0 and $(\varphi_{\varepsilon_n}, \tilde{A}_{\varepsilon_n})$ a sequence of minimizers of G_{ε_n} . One can extract a subsequence of ε_n still denoted ε_n such that there exists an integer multiplicity 1-dimensional rectifiable current \mathbb{L} verifying

$$L = M(\mathbb{L}) < +\infty, \quad (5. 46)$$

$$\partial \mathbb{L} = \sum_i d_i [x_i] \quad (5. 47)$$

$$\mathcal{H}^1(\text{supp} \mathbb{L}) < +\infty \quad (5. 48)$$

and such that $h_{\varepsilon_n} = \tilde{h}_{\varepsilon_n} - h_0 \rightharpoonup h_*$ in $W_{loc}^{1,p}(\mathbb{R}^3)$ for any $1 < p < \frac{3}{2}$ where h_* verifies

$$-\Delta h_* + h_* + h_0 = -2\pi * \mathbb{L} \quad (5. 49)$$

■

Proof of theorem 5.1.

Let $0 < \alpha < 1$ and $\alpha < \gamma < 1$. Consider the domains $R_i, i \leq n + N_\alpha$, given by Lemma 4.1 using the notation

$$R_{l+N_\alpha} = C(\varepsilon^\alpha/\eta, x_l) \text{ for } l = 1, \dots, n.$$

By the same Lemma the set

$$\left\{ x \in \mathbb{R}^3 \setminus \{x_1, \dots, x_n\} \text{ s.t. } |\varphi| \leq \frac{1}{2} \right\} \cap \bigcup_{i=1 \dots N_\alpha+n} \partial R_i$$

can be contained in $C(\alpha, \gamma)/\varepsilon^\alpha$ distinct balls of radius less than ε^γ . Moreover, since for each I a number of those balls bounded by $C'(\alpha, \gamma)$, independent of ε and i , intersect ∂R_i . Provided we multiply the radius by some coefficient independent of i and ε we can consider that, if such a ball intersects ∂R_i , it has a centre on ∂R_i . Denote by $(a_i^k, r_i^k)_{k=1, \dots, n_i}$ the centres and radii of the balls which intersect ∂R_i (n_i is the number of those balls). Moreover denote by d_i^k the degree of $\varphi/|\varphi|$ in any trivialisation of E over R_i (for $i \leq N_\alpha$) on the closed curve $\partial R_i \cap \partial B(a_i^k, r_i^k)$ oriented by the outward unit normal of ∂R_i . Lemma 4.1 says that d_i^k is uniformly bounded independently of i and ε . For topological reasons we have

$$\forall i \leq N_\alpha \quad \sum_{k=1}^{n_i} d_i^k = 0 \quad . \tag{5. 50}$$

For any i and $k \leq n_i$, $\partial R_i \cap \partial B(a_i^k, r_i^k)$ is a union of parts of circles. Each of these parts of circles belongs to the boundary of exactly two R_j which have a face in common, but for the two R_j the orientation of the face is opposite. Thus if one considers a trivialisation of E over $\bar{B}(a_i^k, r_i^k)$ and if one adds all the degrees of $\varphi/|\varphi|$ on the $\cup_j \partial R_j \cap \partial B(a_i^k, r_i^k)$ for j such that $a_i^k \in \partial B_j$, one will cover exactly two times the different parts of circles of $\cup_j \partial R_j \cap \partial B(a_i^k, r_i^k)$ but in the opposite sense and this implies

$$\forall i \leq N_\alpha \quad \forall k \leq n_i \quad \sum_{\{j \leq N_\alpha+n / \exists k_j \leq n_j \text{ s.t. } a_j^{k_j} = a_i^k\}} d_j^{k_j} = 0 \quad . \tag{5. 51}$$

Let $i \leq N_\alpha$, since (5. 50) holds, one can establish a connection between the (a_i^k, d_i^k) included in R_i , that is, there exists a finite union of on closed regular oriented curves included in R_i with boundary included in $\cup_{i=1}^{n_i} \{a_i^k\}$ such that the 1-dimensional integer multiplicity rectifiable current \mathbb{L}_i^ε associated to this union of curves verify

$$\partial \mathbb{L}_i^\varepsilon = \sum_{k=1}^{n_i} d_i^k [[a_i^k]] \quad . \tag{5. 52}$$

Moreover since R_i verifies the condition 1 of Lemma 4.1, since from the conditions 4 and 5 of the same Lemma d_i^k and n_i are bounded independently of i and ε one can choose \mathbb{L}_i^ε such that

$$L_i^\varepsilon = M(\mathbb{L}_i^\varepsilon) \leq C(\alpha, \gamma) \varepsilon^\alpha \quad . \tag{5. 53}$$

Let

$$\mathbb{L}^\varepsilon = \sum_{i=1}^{N_\alpha} \mathbb{L}_i^\varepsilon, \tag{5. 54}$$

since $|N_\alpha| \leq C(\alpha)/\varepsilon^\alpha$ we have

$$L^\varepsilon = M(\mathbb{L}^\varepsilon) \leq C(\alpha, \gamma) \quad \text{indep. of } \varepsilon \quad . \tag{5. 55}$$

Moreover, from (5. 51) we deduce that

$$\partial\mathbb{L}^\varepsilon = \sum_{l=1}^n \sum_{k=1}^{n_{N_\alpha+l}} d_{N_\alpha+l}^k [|a_{N_\alpha+l}^k|] \tag{5. 56}$$

Contrary to the case $i \leq N_\alpha$ for $l = 1, \dots, n$ we cannot trivialize E over $R_{N_\alpha+l} = C\left(\frac{\varepsilon_\alpha}{\eta}, x_l\right)$ thus (5. 50) does not hold in this case but we have that $h_0/2\pi$ is the Euler class of the restriction of the bundle E over $\partial R_{N_\alpha+l}$ and theorem II.16 of [5] implies

$$\forall l = 1 \dots n \quad \sum_{k=1}^{n_{N_\alpha+l}} d_{N_\alpha+l}^k = \frac{1}{2\pi} \int_{\partial R_{N_\alpha+l}} h_0 = d_l \tag{5. 57}$$

and this imply that

$$\partial\mathbb{L}^\varepsilon \longrightarrow \sum_{l=1}^n d_l [|x_l|] \quad \text{as } \varepsilon \rightarrow 0 \quad . \tag{5. 58}$$

Moreover since $n_{N_\alpha+l}$ and $d_{N_\alpha+l}$ are uniformly bounded independently of ε we have

$$M(\partial\mathbb{L}^\varepsilon) \leq C \quad \text{indep. of } \varepsilon. \tag{5. 59}$$

Let ε_n be a sequence of positive number tending to 0, from (5. 55) and (5. 59), using Federer-Fleming's compactness theorem (see for instance [9] page 97), we deduce the existence of a 1-dimensional, integer multiplicity rectifiable current \mathbb{L} and a subsequence still denoted ε_n such that

$$\mathbb{L}^{\varepsilon_n} \rightarrow \mathbb{L} \quad \text{in } \mathcal{D}_1(\mathbb{R}^3) \quad . \tag{5. 60}$$

Moreover from (5. 58) we deduce that

$$\partial\mathbb{L} = \sum_{l=1}^n d_l [|x_l|] \quad . \tag{5. 61}$$

From Lemma 5.1 we know that for any $1 < p < \frac{3}{2}$, $\int_{\mathbb{R}^3} |\nabla h_{\varepsilon_n}|^p < C$ independently of ε and since $h_{\varepsilon_n} \in L^{\frac{3p}{p-3}}(\mathbb{R}^3)$ we have

$$\forall 1 < p < \frac{3}{2} \quad \int_{\mathbb{R}^3} |h_{\varepsilon_n}|^{\frac{3p}{p-3}} + |\nabla h_{\varepsilon_n}|^p \leq C \quad \text{indep. of } \varepsilon \quad .$$

One can extract a subsequence still denoted ε_n such that (5. 60) holds and

$$\text{for } 1 < p < \frac{3}{2} \quad h_{\varepsilon_n} \rightarrow h_\star \quad \text{weakly in } W_{loc}^{1,p}(\mathbb{R}^3). \tag{5. 62}$$

Let ψ be a C^∞ 1-form on \mathbb{R}^3 with compact support. Our aim, now, is to prove

$$\int_{\mathbb{R}^3} \ast d\psi \cdot d^\ast h_{\varepsilon_n} + \ast \psi \cdot \tilde{h}_{\varepsilon_n} + \langle 2\pi \mathbb{L}^{\varepsilon_n}, \psi \rangle \longrightarrow 0 \tag{5. 63}$$

as n tends to infinity. In view of (5. 60) and (5. 62) that will imply (5. 49).

We use the notation h_n, A_n, φ_n, L^n , for $h_{\varepsilon_n}, A_{\varepsilon_n}, \varphi_{\varepsilon_n}$ and L^{ε_n} .

Recall that on $\mathbb{R}^3 \setminus T_{\varepsilon_n^\alpha}$ we have

$$-dd^\ast h_n + \tilde{h}_n = d \left[\left(1 - \frac{1}{|\varphi_n|^2} \right) (i\varphi_n, d_{\tilde{A}_n} \varphi) \right] \quad . \tag{5. 64}$$

Multiplying (5. 64) by $\ast\psi$ and integrating on $\Omega_{\varepsilon_n^\alpha} = \mathbb{R}^3 \setminus T_{\varepsilon_n^\alpha}$ we get

$$\begin{aligned} & \int_{\Omega_{\varepsilon_n^\alpha}} \ast d\psi \cdot d^\ast h_n + \int_{\partial\Omega_{\varepsilon_n^\alpha}} \ast\psi \cdot d^\ast h_n \wedge \nu + \int_{\Omega_{\varepsilon_n^\alpha}} \ast d\psi \cdot \tilde{h} \\ &= - \int_{\Omega_{\varepsilon_n^\alpha}} \left(1 - \frac{1}{|\varphi_n|^2}\right) d\psi \cdot \ast(i\varphi_n, d_{\tilde{A}_n} \varphi_n) \\ &+ \int_{\partial\Omega_{\varepsilon_n^\alpha}} \left(1 - \frac{1}{|\varphi_n|^2}\right) \ast\psi \cdot (i\varphi_n, d_{\tilde{A}_n} \varphi_n) \wedge \nu \end{aligned} \tag{5. 65}$$

where ν denotes the 1-form associated to the unit inward normal of $\partial T_{\varepsilon_n^\alpha}$. Since $d^\ast h_n$ and \tilde{h}_n converges in $L^p(\text{supp}\psi)$ and since $|T_{\varepsilon_n^\alpha}| = |\mathbb{R}^3 \setminus \Omega_{\varepsilon_n^\alpha}| \rightarrow 0$ we have

$$\int_{\Omega_{\varepsilon_n^\alpha}} \ast d\psi \cdot d^\ast h_n + \int_{\Omega_{\varepsilon_n^\alpha}} \ast\psi \cdot \tilde{h}_n \xrightarrow{n \rightarrow +\infty} \int_{\mathbb{R}^3} \ast d\psi \cdot d^\ast h_\ast + \ast\psi \cdot \tilde{h}_\ast \quad . \tag{5. 66}$$

Moreover $(1 - |\varphi_n|^2) \rightarrow 0$ in $L^q(\mathbb{R}^3)$ ($2 \leq q < \infty$) and $(i\varphi_n, d_{\tilde{A}_n} \varphi_n) = -d^\ast h_n$ is bounded in $L^p(\mathbb{R}^3)$ ($1 < p < \frac{3}{2}$), thus

$$\int_{\Omega_{\varepsilon_n^\alpha}} |1 - |\varphi_n|^2| |d\psi| |(i\varphi_n, d_{\tilde{A}_n} \varphi_n)| \xrightarrow{n \rightarrow +\infty} 0 \quad . \tag{5. 67}$$

Moreover

$$\begin{aligned} & \left| \int_{\partial T_{\varepsilon_n^\alpha}} \left(1 - \frac{1}{|\varphi_n|^2}\right) \ast\psi \cdot (i\varphi_n, d_{\tilde{A}_n} \varphi_n) \wedge \nu \right| \\ & \leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} (1 - |\varphi_n|^2) |\nabla_{\tilde{A}_n} \varphi_n| \\ & \leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha+n} \left(\int_{\partial R_i} |1 - |\varphi_n|^2|^2 \right)^{1/2} \left(\int_{\partial R_i} |\nabla_{\tilde{A}_n} \varphi_n|^2 \right)^{1/2} \\ & \leq C \|\psi\|_\infty \frac{1}{\varepsilon_n^\alpha} \varepsilon_n \log \frac{1}{\varepsilon_n} \xrightarrow{n \rightarrow +\infty} 0 \end{aligned} \tag{5. 68}$$

Let ν denotes the 1-form associated to the unit inward normal to ∂R_i , by cancelation on common faces we have

$$- \int_{\partial T_{\varepsilon_n^\alpha}} \ast\psi \cdot d^\ast h_n \wedge \nu = + \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} \ast\psi \cdot (i\varphi_n, d_{\tilde{A}_n} \varphi_n) \wedge \nu \tag{5. 69}$$

Let $\bar{\psi}_i$ be the mean value of ψ on R_i , we have

$$\|\psi - \bar{\psi}_i\|_{L^\infty(R_i)} \leq C \|\nabla\psi\|_\infty \varepsilon^\alpha \quad . \tag{5. 70}$$

and

$$\begin{aligned} \sum_{i=1}^{N_\alpha+n} \int_{\partial R_i} |\psi - \bar{\psi}_i| |(i\varphi_n, d_{\tilde{A}_n} \varphi_n)| \\ \leq C \|\nabla \psi\|_\infty \sum_{i=1}^{N_\alpha+n} \varepsilon_n^\alpha |\partial R_i|^{1/2} \left(\int_{\partial R_i} |\nabla_{\tilde{A}_n} \varphi|^2 \right)^{1/2} \quad (5.71) \\ \leq C \|\nabla \psi\|_\infty \varepsilon_n^\alpha \left(\log \frac{1}{\varepsilon_n} \right)^{1/2} \longrightarrow 0 \text{ as } n \rightarrow +\infty. \end{aligned}$$

We are led to study closely the quantities

$$I_i = \int_{\partial R_i} * \bar{\psi}_i \cdot (i\varphi_n, d_{\tilde{A}_n} \varphi_n) \wedge \nu \quad \text{for } i \leq N_\alpha$$

Use vector instead of 1-forms (ie $\bar{\psi}_i$ denotes also the associated vector to the 1-form $\bar{\psi}_i$).

$$I_i = \int_{\partial R_i} * \bar{\psi}_i \cdot (i\varphi_n, d_{\tilde{A}_n} \varphi_n) \wedge \nu = \int_{\partial R_i} (i\varphi, \nabla_{\tilde{A}_n} \varphi) \cdot \nu \wedge \bar{\psi}_i$$

Decompose $\bar{\psi}_i$ in the canonical basis of \mathbb{R}^3 , $\bar{\psi}_i = \sum_{l=1}^3 \bar{\psi}_i^l e_l$. We have

$$I_i = \sum_{l=1}^3 \bar{\psi}_i^l \int_{\partial R_i} (i\varphi_n, \nabla_{\tilde{A}_n} \varphi_n) \cdot \nu \wedge e_l \quad (5.72)$$

Let $\hat{R}_i = \prod_{l=1}^3 [r_i^l; s_i^l]$ be the convex envelope of R_i and let $\Gamma_l(t) = \partial R_i \cap \{x \in \mathbb{R}^3 / x \cdot e_l = t\}$ as in the proof of Lemma 5.1 (see (5.25)) we have

$$\int_{\partial R_i} (i\varphi_n, \nabla_{\tilde{A}_n} \varphi_n) \cdot \nu \wedge e_l = - \int_{t=r_i^l}^{s_i^l} \oint_{\Gamma_l(t)} (i\varphi_n, \nabla_{\tilde{A}_n} \varphi_n) \quad (5.73)$$

where $\Gamma_l(t)$ is oriented by e_l (i.e. in the plane $\{x \in \mathbb{R}^3 / x \cdot e_l = t\}$ oriented by e_l , $\Gamma_l(t)$ has the positive orientation). Denote by \tilde{I}_i^l the following interval

$$\tilde{I}_i^l = [r_i^l; s_i^l] \setminus \cup_{k=1}^{n_i} (a_i^{k,l} - \varepsilon_n^\gamma; a_i^{k,l} + \varepsilon_n^\gamma) \quad ,$$

where $a_i^{k,l} = a_i^k \cdot e_l$. On $\Gamma_l(t)$ for $t \in \tilde{I}_i^l$ we have $|\varphi| \geq \frac{1}{2}$. Moreover the surface

$$\Sigma_i^l = \cup_{t \in [r_i^l; s_i^l] \setminus \tilde{I}_i^l} \Gamma_l(t)$$

verifies

$$|\Sigma_i^l| \leq C \varepsilon_n^\alpha \varepsilon_n^\gamma \quad . \quad (5.74)$$

In one hand we have

$$\begin{aligned}
 & \left| \sum_{i=1}^{N_\alpha} \sum_{l=1}^3 \bar{\psi}_i^l \int_{[r_i^l; s_i^l] \setminus \bar{I}_i^l} \oint_{\Gamma_l(t)} (i\varphi_n, \nabla_{\tilde{A}_n} \varphi_n) \right| \\
 & \leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha} |\Sigma_i^l|^{1/2} \left(\int_{\partial R_i} |\nabla_{\tilde{A}_n} \varphi_n|^2 \right)^{1/2} \quad (5.75) \\
 & \leq C \|\psi\|_\infty \varepsilon_n^{\frac{\gamma-\alpha}{2}} \left(\log \frac{1}{\varepsilon_n} \right)^{1/2} \xrightarrow{n \rightarrow \infty} 0 \quad .
 \end{aligned}$$

On the other hand

$$\begin{aligned}
 & \left| \sum_{i=1}^{N_\alpha} \sum_{l=1}^3 \bar{\psi}_i^l \int_{\bar{I}_i^l} \oint_{\Gamma_l(t)} \left(1 - \frac{1}{|\varphi_n|^2} \right) (i\varphi_n, \nabla_{\tilde{A}_n} \varphi_n) \right| \\
 & \leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha} \int_{\partial R_i} |1 - |\varphi_n|^2| |\nabla_{\tilde{A}_n} \varphi_n| \quad (5.76) \\
 & \leq C \|\psi\|_\infty \frac{1}{\varepsilon_n^\alpha} \varepsilon_n \log \frac{1}{\varepsilon_n} \xrightarrow{n \rightarrow \infty} 0 \quad \text{as } n \rightarrow \infty \quad .
 \end{aligned}$$

Thus we are led to study the following quantities

$$\int_{\bar{I}_i^l} \oint_{\Gamma_l(t)} \left(i \frac{\varphi_n}{|\varphi_n|}, \nabla_{\tilde{A}_n} \frac{\varphi_n}{|\varphi_n|} \right) = \int_{\bar{I}_i^l} 2\pi d_l(t) - \int_{\bar{I}_i^l} \oint_{\Gamma_l(t)} \tilde{A}_n \quad , \quad (5.77)$$

where $d_l(t)$ is the degree of $\varphi/|\varphi|$ on $\Gamma_l(t)$ in any trivialisation of E over R_i . Let $S_i^l(t)$ be the surface in ∂R_i which has $\Gamma_l(t)$ as boundary (taking into account the orientation of $\Gamma_l(t)$), we have

$$\int_{\bar{I}_i^l} \oint_{\Gamma_l(t)} \tilde{A}_n = \int_{\bar{I}_i^l} \int_{S_i^l(t)} * \tilde{h}_n \cdot \nu \quad . \quad (5.78)$$

Moreover

$$\left| \sum_{i=1}^{N_\alpha} \sum_{l=1}^3 \bar{\psi}_i^l \int_{S_i^l(t)} * \tilde{h}_n \cdot \nu \right| \leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha} \varepsilon_n^\alpha \int_{\partial R_i} |h_0| + |h_n| \quad . \quad (5.79)$$

By Lemma 2.3 we have

$$\begin{aligned}
 \sum_{i=1}^{N_\alpha} \varepsilon_n^\alpha \int_{\partial R_i} |h_n| & \leq \sum_{i=1}^{N_\alpha} \varepsilon_n^{(2+\frac{1}{2})\alpha} \left(\log \frac{1}{\varepsilon_n} \right)^{1/2} \\
 & \leq C \varepsilon_n^{(1+\frac{1}{2})\alpha} \left(\log \frac{1}{\varepsilon_n} \right)^{1/2} \xrightarrow{n \rightarrow +\infty} 0 \quad .
 \end{aligned} \quad (5.80)$$

On ∂R_i we use the bound $\|h_0\|_{L^\infty(\partial R_i)} \leq \frac{C}{\text{dist}(\partial R_i; \{x_1, \dots, x_n\})^2}$. Thus, using notations of Lemma 4.1, we have

$$\begin{aligned} \sum_{i=1}^{N_\alpha} \varepsilon_n^\alpha \int_{\partial R_i} |h_0| &\leq C \varepsilon_n^{3\alpha} \sum_{i=1}^{N_\alpha} \frac{1}{|\text{dist}(\partial R_i; \{x_1, \dots, x_n\})|^2} \\ &\leq C \varepsilon_n^{3\alpha} \sum_{j \in I_\alpha} \frac{1}{|\text{dist}(z_j; \{x_1, \dots, x_n\})|^2} \end{aligned}$$

where $I_\alpha \subset \varepsilon^\alpha \mathbb{Z}^3$ and z_j are defined in part 4.

We have

$$\forall j \in I_\alpha \quad |\text{dist}(z_j, \{x_1, \dots, x_n\})| \geq \varepsilon_n^\alpha.$$

Let $2\alpha/3 < \alpha' < \alpha$. The number of elements of $\varepsilon_n^\alpha \mathbb{Z}^3$ included in $\cup_{k=1}^n C(\varepsilon_n^{\alpha'}, x_k)$ is bounded by $C \varepsilon_n^{3(\alpha' - \alpha)}$ thus separating the $j \in I_\alpha$ which can be included in $\cup_{k=1}^n C(\varepsilon_n^{\alpha'}, x_k)$ and the ones in the exterior of this set we have

$$\begin{aligned} \sum_{i=1}^{N_\alpha} \varepsilon_n^\alpha \int_{\partial R_i} |h_0| &\leq C \varepsilon_n^{3\alpha} \times \frac{1}{\varepsilon_n^{3(\alpha' - \alpha)}} \times \frac{1}{\varepsilon_n^{2\alpha}} + C \varepsilon_n^{3\alpha} \times \frac{1}{\varepsilon_n^3} \times \frac{1}{\varepsilon_n^{2\alpha'}} \\ &\leq C \varepsilon_n^{3\alpha' - 2\alpha} + C \varepsilon_n^{2(\alpha - \alpha')} \end{aligned} \tag{5. 81}$$

Since we have chosen α' such that $3\alpha' - 2 > 0$ and $\alpha' < \alpha$ the right hand side of (5. 81) tends to zero. Combining the previous remark and (5. 80) we have

$$\left| \sum_{i=1}^{N_\alpha} \sum_{l=1}^3 \psi_i^l \int_{S_i^l(t)} * \tilde{h}_n \cdot \nu \right| \rightarrow 0 \quad \text{as } n \rightarrow 0. \tag{5. 82}$$

Let us compute $\int_{\tilde{I}_i} 2\pi d_l(t)$. First extend $d_l(t)$ on $[r_i^l; s_i^l] \setminus \cup_{k=1}^{n_i} \{a_i^{k,l}\}$ in the following way : we extend $\varphi/|\varphi|$ continuously in a map from $\partial R_i \setminus \cup \{a_i^k\}$ into S^1 and for $t \in [r_i^l; s_i^l] \setminus \cup_{k=1}^{n_i} \{a_i^{k,l}\} \cup \tilde{I}_i^l$ define $d_l(t)$ to be the degree of such an extension on $\Gamma_l(t)$. We clearly have

$$\begin{aligned} \left| \sum_{i=1}^{N_\alpha} \sum_{l=1}^n \bar{\psi}_i^l \int_{[r_i^l; s_i^l] \setminus \tilde{I}_i^l} 2\pi d_l(t) \right| &\leq C \|\psi\|_\infty \sum_{i=1}^{N_\alpha} C \varepsilon_n^\gamma \\ &\leq C \|\psi\|_\infty \varepsilon_n^{\gamma - \alpha} \rightarrow 0 \text{ as } n \rightarrow \infty \end{aligned} \tag{5. 83}$$

For l fixed, suppose we have chosen an indexation such that $a_i^{k,l}$ be increasing in k , we have

$$\begin{aligned}
 2\pi \int_{x[r_i^l; s_i^l]} d_l(t) &= 2\pi (a_i^{2,l} - a_i^{1,l}) \times (-d_i^{1,l}) \\
 &+ 2\pi (a_i^{3,l} - a_i^{2,l}) \times (-d_i^{1,l} - d_i^{2,l}) + \dots \\
 &+ 2\pi (a_i^{n_i,l} - a_i^{n_i-1,l}) \times (-d_i^{1,l} - \dots - d_i^{n_i-1,l}) \\
 &+ 2\pi (s_i^l - a_i^{n_i,l}) \times (-\sum_{k=1}^{n_i} d_i^{k,l}) \quad .
 \end{aligned}
 \tag{5.84}$$

Considering (5. 50) we get

$$2\pi \int_{r_i^l}^{s_i^l} d_l(t) = 2\pi \sum_{k=1}^{n_i} d_i^{k,l} a_i^{k,l} \quad .
 \tag{5.85}$$

On the other hand we have

$$\langle \mathbb{L}_i^{\varepsilon_n}; e_l \rangle = \sum_{k=1}^{n_i} d_i^{k,l} a_i^{k,l} \quad .
 \tag{5.86}$$

Thus

$$\sum_{i=1}^{N_\alpha} \sum_{l=1}^3 \bar{\psi}_i^l \quad 2\pi \int_{r_i^l}^{s_i^l} d_l(t) = \sum_{i=1}^{N_\alpha} \langle \mathbb{L}_i^{\varepsilon_n}; \bar{\psi}_i \rangle
 \tag{5.87}$$

Moreover

$$\begin{aligned}
 \langle \mathbb{L}^{\varepsilon_n}; \psi \rangle &= \sum_{i=1}^{N_\alpha} \langle \mathbb{L}_i^{\varepsilon_n}; \psi \rangle \\
 &= \sum_{i=1}^{N_\alpha} \langle \mathbb{L}_i^{\varepsilon_n}; \bar{\psi}_i \rangle + \sum_{i=1}^{N_\alpha} \langle \mathbb{L}_i^\varepsilon; \psi - \bar{\psi}_i \rangle,
 \end{aligned}
 \tag{5.88}$$

and we have

$$\begin{aligned}
 \left| \sum_{i=1}^{N_\alpha} \langle \mathbb{L}_i^{\varepsilon_n}, \psi - \bar{\psi}_i \rangle \right| &\leq \sum_{i=1}^{N_\alpha} M(\mathbb{L}_i^{\varepsilon_n}) \|\nabla \psi\|_\infty \varepsilon_n^\alpha \\
 &\leq M(\mathbb{L}) \|\nabla \psi\|_\infty \varepsilon_n^\alpha \longrightarrow 0 \quad \text{as } n \rightarrow 0
 \end{aligned}
 \tag{5.89}$$

Combining (5. 69), (5. 71), (5. 72), (5. 73), (5. 75), (5. 76), (5. 82), (5. 83), (5. 87) and (5. 88) we get

$$- \int_{\partial T_{\varepsilon_n}^\alpha} * \psi . d^* h_n \wedge \nu + \langle \mathbb{L}^{\varepsilon_n}; \psi \rangle \longrightarrow 0 \quad \text{as } n \rightarrow +\infty
 \tag{5.90}$$

Combining (5. 65), (5. 66), (5. 67), (5. 68) and (5. 90) we prove (5. 63) which implies (5. 49).

It remains to prove that $\mathcal{H}^1(\text{supp} \mathbb{L}) < +\infty$. Let r be a positive number,

(we take $r < \eta$ where η is the constant given by Lemma 3.1). Consider the lattice $r\mathbb{Z}^3$, as in the proof of Lemma 4.1 the number of $x \in r\mathbb{Z}^3$ such that

$$\frac{1}{2r} \int_{\partial C(2r,x)} |\nabla_{\tilde{A}_n} \varphi_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 + |h_n|^2 \geq \eta \log\left(\frac{2r}{\varepsilon_n}\right)$$

is bounded by C/r . This implies that $\text{supp}(\mathbb{L}^{\varepsilon_n})$ can be included in $\frac{C}{r}$ cubes of edge $2r$. Let $R > 0$, we can extract from ε_n a subsequence such that the bad cubes $C(2r, z)$ for $z \in r\mathbb{Z}^3 \cap B_R(0)$ are always the same, thus we are sure that $\text{supp}\mathbb{L} \cap B_R(0)$ is included in those cubes. Thus $\forall r, R > 0$ $\text{supp}\mathbb{L} \cap B_R(0)$ is included in C/r cubes where C is independent of R and r , this implies

$$\mathcal{H}^1(\text{supp}\mathbb{L}) < +\infty.$$

■

6. STRONG CONVERGENCE OF THE MAGNETIC FIELD AWAY FROM THE LIMITING SINGULAR SET

In this part we prove the strong convergence in $H^1_{loc}(\mathbb{R}^3 \setminus \text{supp}\mathbb{L})$ of the magnetic field h_{ε_n} for the subsequence given by theorem 5.1. As in [2] and [3], with some more work, it should be possible to prove the convergence away from the limiting singular set in a stronger norm, $C^k (k \geq 0)$ for instance.

THEOREM 6.1. Let h_{ε_n} be the sequence given by theorem 5.1 and let \mathbb{L} be the rectifiable current in the same theorem. We have

$$h_{\varepsilon_n} \rightarrow h_\star \text{ strongly in } H^1_{loc}(\mathbb{R}^3 \setminus \text{supp}\mathbb{L}) \quad . \quad (6. 1)$$

■

Proof of theorem 6.1.

We still use the notation $A_n, h_n, \varphi_n \dots$ for $A_{\varepsilon_n}, h_{\varepsilon_n}, \varphi_{\varepsilon_n} \dots$. Let $x_0 \in \mathbb{R}^3 \setminus \text{supp}\mathbb{L}$ and r small enough such that $B_r(x_0) \subset U_l$ for some $l \in \{1, \dots, k+1\}$ (we use the notation of part 3) and such that $r < \min(\eta; d_{x_0}/\eta)$ where η is given by Lemma 3.1. This is possible since $\{x_1, \dots, x_n\} \subset \text{supp}\mathbb{L}$. Remark that we have $B_r(x_0) \cap \text{supp}\mathbb{L} = \emptyset$.

We will prove that h_n converges strongly in $H^1(B_{r/4}(x_0))$ to h_\star given by (5. 49).

Since the limit h_\star is well identified (= is unique) and since we have weak convergence in $W^{1,p} (1 < p < 3/2)$ $h_n \rightharpoonup h_\star$, it is sufficient to prove the existence of one subsequence of h_n which strongly converges in H^1 to prove the strong convergence of the complete sequence.

We can trivialize E over $B_r(x_0)$ and since $B_r(x_0)$ is included in some U_l we can take the Coulomb gauge given by proposition 2.1 ie $\tilde{A} = A_l^0 + A_n$. Since h_n weakly converges in $W^{1,p}(B_r(x_0))$, A_n strongly converges in $W^{1,p}(B_r(x_0))$ and because of corollary 5.1 (5. 38) we can extract a subsequence such that φ_n weakly converges in $W^{1,p}(B_r(x_0)) (1 < p < 3/2)$ [in fact it is not difficult to see that the complete sequence φ_n converges but, because of the

remark above, we do not take care of it]. Let φ_* be the limit of φ_n , since $\int_{B_r(x_0)} (1 - |\varphi_n|^2)^2 \rightarrow 0$, $|\varphi_*| \equiv 1$ and passing to the limit in (2. 15) we have

$$-d^*h_* = (i\varphi_*, d\varphi_*) - A_l^0 - A_*. \tag{6. 2}$$

Moreover since $\text{supp}\mathbb{L} \cap B_r(x_0) = \emptyset$, (5. 49) implies

$$-dd^*h_* + dA_* + dA^0 = 0 \quad \text{in } \mathcal{D}'(B_r(x_0)) \quad . \tag{6. 3}$$

Combining (6. 2) and (6. 3) we get

$$d(i\varphi_*, d\varphi_*) = 0 \quad \text{in } \mathcal{D}'(B_r(x_0)) \quad . \tag{6. 4}$$

Moreover since $d^*A_l^0 = 0$ and $d^*A_* = 0$ in $(B_r(x_0))$ we have also

$$d^*(i\varphi_*, d\varphi_*) = 0 \quad \text{in } \mathcal{D}'(B_r(x_0)) \quad . \tag{6. 5}$$

From (6. 4) and (6. 5) we deduce the existence of an harmonic function H such that

$$(i\varphi_*, d\varphi_*) = dH \quad . \tag{6. 6}$$

Since $|\varphi_*| \equiv 1$, we have

$$d\varphi_* = (i\varphi_*, d\varphi_*)i\varphi_* \quad . \tag{6. 7}$$

Combining (6. 6) and (6. 7) we deduce that

$$d(\varphi_*e^{-iH}) = 0 \quad \text{in } \mathcal{D}'(B_r(x_0)) \tag{6. 8}$$

Thus φ_* is a harmonic map from $B_r(x_0)$ into S^1 , in particular $\varphi_* \in C^\infty(B_r(x_0))$.

Our aim, now, is to prove that there exists a subsequence, still denoted φ_n such that for n sufficiently large we have

$$|\varphi_n| \geq \frac{1}{2} \quad \text{in } B_{r/4}(x_0) \quad . \tag{6. 9}$$

Indeed, as we have already seen below, (6. 9) and the weak convergence of h_n imply the strong convergence in $B_{r/4}(x_0)$.

Let $0 < s < r$, s will be chosen sufficiently small later. Consider $s\mathbb{Z}^3 \cap B_r(x_0)$ as in the end of the proof of theorem V.1 $B_r(x_0) \cap \text{supp}\mathbb{L}^{\varepsilon_n}$ is included in C/s cubes $C(2s, x)$ where $x \in s\mathbb{Z}^3$ where C does not depend neither on s nor on ε_n . Extract a subsequence such that the bad cubes in $B_r(x_0)$ are always the same and denote by T_s this union of cubes.

Let $\Omega_{r,s} = B_r(x_0) \setminus \bar{T}_s$ we claim that $h_n \rightarrow h_*$ strongly in $H_{loc}^1(\Omega_{r,s})$. Let ω be an open set such that $\omega \subset\subset \Omega_{r,s}$ and χ be a non negative cut-off function such that

$$\begin{cases} \chi \equiv 1 & \text{in } \omega \\ \chi \equiv 0 & \text{in } \mathbb{R}^3 \setminus \Omega_{r,s} \quad . \end{cases} \tag{6. 10}$$

From (2. 15) and the fact that $|\varphi_n| \geq \frac{1}{2}$ in $\Omega_{r,s}$ we have

$$-d\left(\frac{1}{|\varphi_n|^2}d^*\tilde{h}_n\right) + \tilde{h}_n = 0 \quad \text{in } \Omega_{r,s} \quad . \tag{6. 11}$$

Let K_n be the solution of

$$\begin{cases} -\Delta K_n + K_n = f_n & \text{in } \Omega_{r,s} \\ K_n = 0 & \text{on } \partial\Omega_{r,s} \end{cases}, \tag{6. 12}$$

where $f_n = d^*\tilde{h}_n/|d^*\tilde{h}_n|^\beta$ in ω and $f_n \equiv 0$ in $\Omega_{r,s} \setminus \omega$ for $0 < \beta < 1$ (fixed later). Let $k_n = dK_n$. For q and β such that $q(1 - \beta) < 2$ and $2 < q < 3$ we have by classical elliptic estimates

$$\left(\int_{\Omega_{r,s}} |k_n|^{\frac{3q}{3-q}} \right)^{\frac{1}{q} - \frac{1}{3}} + \left(\int_{\Omega_{r,s}} |d^*k_n|^q \right)^{1/q} \leq C \left(\int_{\omega} |d^*\tilde{h}_n|^{q(1-\beta)} \right)^{1/q} \tag{6. 13}$$

Multiplying (6. 11) by χk_n and integrating on $\Omega_{r,s}$ we have

$$\begin{aligned} & \int_{\Omega_{r,s}} \chi d^*k_n d^*\tilde{h}_n + \int_{\Omega_{r,s}} k_n \cdot d^*\tilde{h}_n \wedge d\chi + \int_{\Omega_{r,s}} \chi k_n \tilde{h}_n \\ &= \int_{\Omega_{r,s}} \left(1 - \frac{1}{|\varphi_n|^2} \right) \left[d^*\tilde{h}_n d^*k_n \chi + k_n \cdot d^*\tilde{h}_n \wedge \chi \right]. \end{aligned} \tag{6. 14}$$

In the other hand, taking the exterior derivative of (6. 12), multiplying by $\chi \tilde{h}_n$ and integrating on $\Omega_{r,s}$ we have

$$\begin{aligned} & \int_{\Omega_{r,s}} \chi d^*\tilde{h}_n \cdot d^*k_n + \chi \tilde{h}_n \cdot k_n + \tilde{h}_n \cdot d\chi \wedge d^*k_n \\ &= - \int_{\omega} |d^*\tilde{h}_n|^{2-\beta} \end{aligned} \tag{6. 15}$$

Combining (6. 14) and (6. 15) we get

$$\begin{aligned} \int_{\omega} |d^*\tilde{h}_n|^{2-\beta} &\leq C \left[\int_{\Omega_{r,s}} (1 - |\varphi_n|^2) |d^*\tilde{h}_n| |d^*k_n| + \right. \\ &\quad \left. + \int_{\Omega_{r,s}} |k_n| |d^*\tilde{h}_n| + |\tilde{h}_n| |d^*k_n| \right] \end{aligned} \tag{6. 16}$$

First of all we have

$$\begin{aligned} & \left| \int_{\Omega_{r,s}} (1 - |\varphi_n|^2) |d^*\tilde{h}_n| (|d^*k_n| + |k_n|) \right| \leq \\ & \leq C \varepsilon_n^{2(1-\frac{1}{2}-\frac{1}{q})} \left(\log \frac{1}{\varepsilon_n} \right)^{1/2} \left[\int_{\omega} |d^*\tilde{h}_n|^{q(1-\beta)} \right]^{1/q}. \end{aligned} \tag{6. 17}$$

we have also

$$\int_{\Omega_{r,s}} |k_n| |d^*\tilde{h}_n| \leq \left(\int_{\Omega_{r,s}} |k_n|^{\frac{3q}{3-q}} \right)^{\frac{1}{q} - \frac{1}{3}} \times \left(\int_{\Omega_{r,s}} |d^*\tilde{h}_n|^{\frac{3q}{4q-3}} \right)^{1 + \frac{1}{3} - \frac{1}{q}} \tag{6. 18}$$

and for $2 < q < 3$, $\frac{2}{3} < \frac{5}{6} < \frac{4}{3} - \frac{1}{q} < 1$, thus (6. 18) implies

$$\int_{\Omega_{r,s}} |k_n| |d^* \tilde{h}_n| \leq C \left(\int_{\omega} |d^* \tilde{h}_n|^{q(1-\beta)} \right)^{1/q} \tag{6. 19}$$

Finally

$$\int_{\Omega_{r,s}} |\tilde{h}_n| |d^* k_n| \leq \left(\int_{\Omega_{r,s}} |\tilde{h}_n|^{1-\frac{1}{q}} \right) \left(\int_{\Omega_{r,s}} |d^* k_n|^q \right)^{1/q} \tag{6. 20}$$

and since \tilde{h}_n is bounded in $W^{1,p}(B_r(x_0))$ for any $1 < p < \frac{3}{2}$, \tilde{h}_n is bounded in $L^s(B_r(x_0))$ for any $1 < s < 3$ and since $2 < q, \frac{q}{q-1} < 2$ and (6. 20) combined with (6. 13) gives

$$\int_{\Omega_{r,s}} |\tilde{h}_n| |d^* k_n| \leq C \left(\int_{\omega} |d^* \tilde{h}_n|^{q(1-\beta)} \right)^{1/q} \tag{6. 21}$$

Combining (6. 16), (6. 17), (6. 19) and (6. 20) we have

$$\int_{\omega} |d^* \tilde{h}_n|^{2-\beta} \leq C \left(\int_{\omega} |d^* \tilde{h}_n|^{q(1-\beta)} \right)^{1/q}$$

Taking $2 - \beta = q(1 - \beta)$, $2 < q < 3$ implies $\frac{3}{2} < 2 - \beta < 2$ thus we have

$$\forall 1 \leq p < 2 \quad \int_{\omega} |d^* \tilde{h}_n|^p < C \quad \text{indep. of } n \tag{6. 22}$$

Let $\omega' \subset\subset \omega$ since we have

$$\Delta \tilde{h}_n = d(d^* \tilde{h}_n) \quad \text{in } \omega \tag{6. 23}$$

and since \tilde{h}_n is uniformly bounded in $W^{1,p}(\omega)$ $1 \leq p < \frac{3}{2}$ in view of (6. 22) we have, by classical elliptic estimates that \tilde{h}_n is uniformly bounded in $W^{1,p}(\omega')$ $1 \leq p < 2$.

Thus we have proved that \tilde{h}_n is uniformly bounded in any compact set included in $\Omega_{r,s}$ for the norm $W^{1,p}$ $1 \leq p < 2$.

We establish now the strong H^1_{loc} convergence of \tilde{h}_n in $\Omega_{r,s}$.
 Multiplying (6. 11) by $\chi \tilde{h}_n$ and integrating on $\Omega_{r,s}$ we have

$$\int_{\Omega_{r,s}} \chi \frac{1}{|\varphi_n|^2} |d^* \tilde{h}_n|^2 + \int_{\Omega_{r,s}} \frac{1}{|\varphi_n|^2} \tilde{h}_n \cdot d\chi \wedge d^* \tilde{h}_n + \chi |\tilde{h}_n|^2 = 0. \tag{6. 24}$$

Thus we have

$$\int_{\omega} |d^* \tilde{h}_n|^2 \leq C \left(\int_{supp\chi} |\tilde{h}_n|^4 \right)^{1/4} \left(\int_{\Omega_{r,s}} |d^* \tilde{h}_n|^{\frac{4}{3}} \right)^{3/4} \tag{6. 25}$$

Since \tilde{h}_n is uniformly bounded for the norm $W^{1,p}$ for any $p < 2$ on $supp\chi \subset\subset \Omega_{r,s}$, \tilde{h}_n converges to h_* in $L^4(supp\chi)$ and we obtain that $d^* \tilde{h}_n$ is uniformly bounded for the L^2 norm on any compact set included in $\Omega_{r,s}$.

Multiplying (6. 11) by $\chi(\tilde{h}_n - \tilde{h}_*)$, we have

$$\int_{\Omega_{r,s}} \chi \frac{1}{|\varphi_n|^2} |d^* \tilde{h}_n|^2 - \int_{\Omega_{r,s}} \chi \frac{1}{|\varphi_n|^2} d^* \tilde{h}_* . d^* \tilde{h}_n + \int_{\Omega_{r,s}} \frac{1}{|\varphi_n|^2} (\tilde{h}_n - \tilde{h}_*) . d\chi \wedge d^* \tilde{h}_n \int_{\Omega_{r,s}} \chi (\tilde{h}_n - \tilde{h}_*) . \tilde{h}_n = 0 \quad .$$

Hence we deduce that

$$\int_{\Omega_{r,s}} \chi \frac{1}{|\varphi_n|^2} |d^* \tilde{h}_n|^2 \longrightarrow \int_{\Omega_{r,s}} \chi |d^* \tilde{h}_*|^2 \quad \text{as } n \rightarrow \infty \quad .$$

By lower semicontinuity this implies the strong convergence of $d^* \tilde{h}_n$ to $d^* \tilde{h}_*$ in $L^2(K)$ for any compact K included in $\Omega_{r,s}$. Because of (6. 23), since we have already strong convergence of \tilde{h}_n to h_* in $L^2(\Omega_{r,s})$ we deduce that \tilde{h}_n strongly converges to h_* in $H^1_{loc}(\Omega_{r,s})$. This convergence, combined with (2. 15), yields

$$\left(i \frac{\varphi_n}{|\varphi_n|}, d \frac{\varphi_n}{|\varphi_n|} \right) \longrightarrow (i\varphi_*, d\varphi_*) \quad \text{in } L^2_{loc}(\Omega_{r,s}) \quad . \quad (6. 26)$$

Moreover, since $\varphi_* \in C^\infty(B_r(x_0), S^1)$, the degree of φ_* on all the generators of $\pi_1(\Omega_{r,s})$ is 0 and the convergence (6. 26) implies that this is also the case for $\varphi_n/|\varphi_n|$ for n chosen sufficiently large. Let Γ be one of the regular curves in $\Omega_{r,s}$ which generates $\pi_1(\Omega_{r,s})$ and let U be a tubular neighborhood of Γ in $\Omega_{r,s}$. Consider a closed 2-form ω_Γ with support in U which is the Poincaré dual of Γ (see [5] I.6 for the construction of ω_Γ) that is, for any closed 1-form ψ in U , ω_Γ verifies

$$\int_U \psi \wedge \omega_\Gamma = \int_\Gamma \psi \quad . \quad (6. 27)$$

Since $\varphi_n/|\varphi_n|$ is regular in U

$$d \left(i \frac{\varphi_n}{|\varphi_n|}, d \frac{\varphi_n}{|\varphi_n|} \right) = 0 \quad \text{in } U \quad . \quad (6. 28)$$

Thus

$$\int_U \left(\frac{i\varphi_n}{|\varphi_n|}, d \frac{\varphi_n}{|\varphi_n|} \right) \wedge \omega_\Gamma = \int_\Gamma \left(\frac{i\varphi_n}{|\varphi_n|}, d \frac{\varphi_n}{|\varphi_n|} \right) = 2\pi \text{ deg}_\Gamma \frac{\varphi_n}{|\varphi_n|} \quad , \quad (6. 29)$$

and in the same way

$$\int_U (i\varphi_*, d\varphi_*) \wedge \omega_\Gamma = \int_\Gamma (i\varphi_*, d\varphi_*) = 2\pi \text{ deg}_\Gamma \varphi_* = 0 \quad . \quad (6. 30)$$

Since

$$\int_U \left(\frac{i\varphi_n}{|\varphi_n|}, d \frac{\varphi_n}{|\varphi_n|} \right) \wedge \omega_\Gamma = \int_\Gamma (i\varphi_*, d\varphi_*) \wedge \omega_\Gamma \quad , \quad (6. 31)$$

we have, for n chosen sufficiently large, $\text{deg}_\Gamma \frac{\varphi_n}{|\varphi_n|} = 0$.

We prove now that $|\varphi_n| \geq \frac{1}{2}$ in $B_{r/4}(x_0)$ that is $T_{r,s} \cap B_{r/4}(x_0) = \emptyset$. Since we have already established the strong convergence in $H^1_{loc}(B_r(x_0) \setminus T_{r,s})$ of

\tilde{h}_n , in view of Lemma 3.1, to get the desired result, it suffices to prove that, for n sufficiently large, we have

$$\frac{2}{r} \int_{B_{r/2}(x_0)} |\nabla_{\tilde{A}} \varphi_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 + |h_n|^2 < \eta \log \left(\frac{r}{2\varepsilon_n} \right) \quad (6. 32)$$

Applying the mean value formula we deduce the existence of $\rho \in (\frac{r}{2}; r)$ such that

$$\left\{ \begin{array}{l} \partial B_\rho(x_0) \cap T_{r,s} \subset N \text{ balls of radius } s \\ \text{denoted } \cup_{i=1}^N B(a_i, s) \text{ where } N \text{ is indep. of } s \text{ and } \varepsilon_n. \\ \int_{\partial B_\rho(x_0)} |\nabla_{\tilde{A}_n} \varphi_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 + |h_n|^2 \leq C \log \left(\frac{1}{\varepsilon_n} \right) \\ \int_{\partial B_\rho(x_0) \setminus \cup_{i=1}^N B(a_i, 2s)} |\nabla \varphi_n|^2 \leq C(r, s) \end{array} \right. \quad (6. 33)$$

This is possible because $T_{r,s}$ is included in C/s balls of radius s (where C is independent of s and ε_n), because Lemma 2.5 holds and φ_n converges strongly in $H_{loc}^1(T_{r,s})$. First of all, since N is independent on s and ε_n , there exists $\lambda \geq 2$ (independent on s and ε_n) $N' \leq N$ and $(b_j)_{j \in N'}$ such that

$$\left\{ \begin{array}{l} b_j \in \partial B_\rho(x_0) \\ \cup_{i=1}^N B(a_i, s) \subset \cup_{j=1}^{N'} B(b_j, \lambda s) \\ \text{for } j \neq j' \quad B(b_j, 2\lambda s) \cap B(b_{j'}, 2\lambda s) = \emptyset \end{array} \right. \quad (6. 34)$$

To prove (6. 32), since (φ_n, \tilde{A}_n) is a minimizer of the energy, it suffices to find $\varphi'_n : B_\rho(x_0) \rightarrow \mathbb{C}$ such that $\varphi'_n = \varphi_n$ on $\partial B_\rho(x_0)$ and such that $(\varphi'_n, \tilde{A}_n)$ verifies (6. 32).

Construction of φ'_n .

First, for $\rho \geq t \geq \rho - s$, we contract the bad parts contained in the $B(b_j, \lambda s)$ in the following way. By the mean value formula we can find $\mu \in (\lambda, 2\lambda)$ such that

$$\forall j \leq N' \quad \int_{\partial(B(b_j, \mu s) \cap \partial B_\rho(x_0))} |\nabla \varphi_n|^2 \leq C(r, s) \quad (6. 35)$$

and

$$\forall j \leq N' \quad \int_{\partial(B(b_j, \mu s) \cap \partial B_\rho(x_0))} \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 \leq \frac{C}{s} \log \frac{1}{\varepsilon_n} \quad (6. 36)$$

where C is independent of n and s .

For $\rho - s \leq t \leq \rho$ denote by T_t the following set

$$T_t = \frac{t}{\rho} \left[\partial B_\rho(x_0) \setminus \cup_{j=1}^{N'} B(b_j, \mu s) \right]$$

and

$$T = \bigcup_{\rho-s \leq t \leq \rho} T_t$$

Let $B^2(x, l)$ be the geodesic ball in $\partial B(x_0, |x - x_0|)$ of centre x and radius $2 \arcsin(l/2|x|)$ corresponding to the chord of length l .

In $B^2(\frac{t}{\rho}b_j, \frac{t}{\rho}\mu s)$ a point x is represented by the coordinates (θ, l, t) where θ is the angle in geodesic polar coordinates in $B^2(\frac{t}{\rho}b_j; \frac{t}{\rho}\mu s)$, l is the distance to $\frac{t}{\rho}b_j$ and $t = |x|$.

We have seen above that, for n sufficiently large, the degree of $\varphi_n/|\varphi_n|$ on $\partial(B^2(b_j, \mu s) \cap B_\rho(x_0)) = \partial B^2(b_j, \mu s)$ is equal to 0. Thus there exists a regular function $\psi_n^j : \partial B^2(b_j, \mu s) \rightarrow \mathbb{R}$ such that

$$\left\{ \begin{array}{l} \frac{\varphi_n}{|\varphi_n|}(\theta, \mu s, \rho) = e^{i\psi_n^j(\theta)} e^{i\theta_n^j} \quad \text{in } \partial B^2(b_j, \mu s) \\ \text{and} \\ \int_{\partial B^2(b_j, \mu s)} \psi_n^j = 0 \quad , \end{array} \right.$$

where θ_n^j is the representant, in $[0; 2\pi[$ of the mean value of $\varphi_n/|\varphi_n|$ along $\partial B^2(b_j, \mu s)$. From (6. 35) and Poincaré inequality, we deduce that

$$\int_{\partial B^2(b_j, \mu s)} |\psi_n^j|^2 \leq C \int_{\partial B^2(b_j, \mu s)} |\nabla \psi_n^j|^2 \leq C(s, r) \quad . \quad (6. 37)$$

Let $C^j = \cup_{\rho-s < t < \rho} \frac{t}{\rho} B^2(b_j, \mu s) = \cup_{\rho-s < t < \rho} B^2(\frac{t}{\rho}b_j, \frac{t}{\rho}\mu s)$. C^j is starshaped from $b_j^s = (1 - \frac{s}{\rho})b_j$. Denote by π^j the map from \bar{C}^j into ∂C^j which associates to $x \in \bar{C}^j$ the point of ∂C^j which is on the line starting from b_j^s to x . We have

$$|\nabla \pi^j|(x) \leq \frac{C s}{|x - b_j^s|} \quad (6. 38)$$

where C is independent of s .

Define $\varphi_n'(x)$ for $\rho - s \leq |x| \leq \rho$ in the following way

$$\left\{ \begin{array}{l} \varphi_n'(x) = \varphi_n\left(\frac{\rho x}{t}\right) \quad \text{in } \bar{T}_t \\ \varphi_n'(\theta, l, \rho - s) = \left[1 + (|\varphi_n|(\theta, \mu s; \rho) - 1) \frac{l}{\mu s} \frac{\rho}{\rho - s} \right] \times \\ \quad \times \exp \left[i \frac{l}{\mu s} \frac{\rho}{\rho - s} \psi_n^j(\theta, \mu s, \rho) \right] \times e^{i\theta_n^j} \quad (6. 39) \\ \text{in } B^2\left(\frac{\rho - s}{\rho}b_j, \frac{\rho - s}{\rho}\mu s\right) \\ \text{and } \varphi_n'(x) = \varphi_n'(\pi^j(x)) \quad \text{in } C^j \end{array} \right.$$

First of all in view of (6. 33) we have

$$\int_T |\nabla \varphi_n'|^2 \leq C \int_{\rho-s}^\rho \int_{T_\rho} |\nabla \varphi_n|^2 \leq C(s, r) \quad (6. 40)$$

and

$$\int_T \frac{1}{\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 \leq C \int_{\rho-s}^\rho \int_{T_\rho} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n|^2)^2 \leq Cs \log \frac{1}{\varepsilon_n} \quad (6.41)$$

On $\partial T \cap \partial C_j$ we have also, in view of (6.35)

$$\int_{\partial T \cap \partial C_j} |\nabla \varphi'_n|^2 \leq C \int_{\rho-s}^\rho \int_{\partial B^2(b_j, \mu s)} |\nabla \varphi_n|^2 \leq C(s, r) \quad (6.42)$$

and in view of (6.36)

$$\begin{aligned} \int_{\partial T \cap \partial C_j} \frac{1}{\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 &\leq C \int_{\rho-s}^\rho \int_{\partial B^2(b_j, \mu s)} \frac{1}{\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 \\ &\leq C \log \frac{1}{\varepsilon_n} \quad (6.43) \end{aligned}$$

In $\frac{\rho-s}{\rho} \times B^2(b_j, \mu s)$ we have

$$\begin{aligned} |\nabla_{\mathbb{T}} \varphi'_n| &\leq C \left[\left| \frac{1}{l} \frac{\partial \varphi'_n}{\partial \theta} \right| + \left| \frac{\partial \varphi'_n}{\partial l} \right| \right] \\ &\leq C \left[\left| \frac{1}{\mu s} \frac{\partial \psi_n^j}{\partial \theta} \right| + \frac{1}{\mu s} |\psi_n^j| + \left| \frac{1}{\mu s} \frac{\partial |\varphi_n|}{\partial \theta} \right| (\theta, \mu s, \rho) \right. \\ &\quad \left. + \frac{1}{\mu s} (|\varphi_n|(\theta, \mu s, \rho) - 1) \right] \quad (6.44) \end{aligned}$$

The integration of (6.44) in $\frac{\rho-s}{\rho} B^2(b_j, \mu s)$ gives, in view of (6.35)

$$\begin{aligned} \int_{\frac{\rho-s}{\rho} B^2(b_j, \mu s)} |\nabla_{\mathbb{T}} \varphi'_n|^2 &\leq \\ &\leq C \int_0^{\frac{\rho-s}{\rho} \mu s} \int_0^{2\pi} \left(|\nabla \varphi_n|^2(\theta, \mu s, \rho) + \frac{C}{s} \right) l \, dl \, d\theta \quad (6.45) \end{aligned}$$

and note that $|1 - |\varphi'_n|(\theta, l, \rho - s)| \leq |1 - |\varphi_n|(\theta, \mu s, \rho)|$ thus we have, using (6.36),

$$\begin{aligned} \int_{\frac{\rho-s}{\rho} B^2(b_j, \mu s)} \frac{1}{\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 &\leq \\ &\leq \int_0^{\frac{\rho-s}{\rho} \mu s} \int_0^{2\pi} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n|^2)^2(\theta, \mu s, \rho) \, l \, dl \, d\theta \\ &\leq \int_0^{\frac{\rho-s}{\mu} s} \int_0^{2\pi} \frac{l}{\mu s} \, dl \int_{\partial(B(b_j, \mu s) \cap \partial B_\rho(x_0))} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n|^2)^2 \leq C \log \frac{1}{\varepsilon_n} \quad (6.46) \end{aligned}$$

Finally, let C_k^j be the homotetic of C^j by the homotetic of centre b_j^s and rate k , we have

$$\begin{aligned} \int_{C^j} |\nabla \varphi_n'|^2 &\leq C \int_{v=0}^s dv \int_{\partial C_{v/s}^j} |\nabla_{\top} \varphi_n(\pi^j(x))|^2 |\nabla \pi^j|^2(x) \\ &\leq C \int_{v=0}^s \frac{s^2}{v^2} dv \int_{\partial C_{v/s}^j} |\nabla_{\top} \varphi_n(\pi^j(x))|^2 \\ &\leq C s \int_{\partial C^j} |\nabla_{\top} \varphi_n|^2 \leq C s \log \frac{1}{\varepsilon_n} \end{aligned} \tag{6. 47}$$

Moreover we have

$$\begin{aligned} \int_{C^j} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n'|^2)^2 &\leq C \int_0^s \frac{1}{\varepsilon_n^2} \int_{\partial C_{v/s}^j} (1 - |\varphi_n|^2)^2 (\pi^j(x)) \\ &\leq C \int_0^s \frac{v^2}{s^2} dv \int_{\partial C^j} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n|^2)^2 \\ &\leq C s \int_{\partial C^j} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n|^2)^2 \leq C s \log \frac{1}{\varepsilon_n} \end{aligned} \tag{6. 48}$$

Combining (6. 40), (6. 41), (6. 47) and (6. 48) we obtain that φ_n' verifies

$$\int_{B_{\rho}(x_0) \setminus B_{\rho-s}(x_0)} |\nabla \varphi_n'|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n'|^2)^2 \leq C s \log \frac{1}{\varepsilon_n} \tag{6. 49}$$

Moreover $|\varphi_n'| \geq \frac{1}{2}$ on $\partial B_{\rho-s}(x_0)$ and from (6. 45) we deduce

$$\int_{\partial B_{\rho-s}(x_0)} |\nabla_{\top} \varphi_n'|^2 \leq C(r, s) \tag{6. 50}$$

and from (6. 46)

$$\int_{\partial B_{\rho-s}(x_0)} \frac{1}{\varepsilon_n^2} (1 - |\varphi_n'|^2)^2 \leq C \log \frac{1}{\varepsilon_n} \tag{6. 51}$$

Now we fill-up the ball $B_{\rho-s}(x_0)$ as we did at the end of the proof of Lemma 3.1.

Take the harmonic extension of ψ_n , where ψ_n verifies $e^{i\psi_n} = \varphi_n'/|\varphi_n'|$ on $\partial B_{\rho-s}(x_0)$ and if ξ_n is the solution of

$$\begin{cases} -\varepsilon_n^2 \Delta \xi_n + \xi_n = 1 & \text{in } B_{\rho-s}(x_0) \\ \xi_n = |\varphi_n'| & \text{on } \partial B_{\rho-s}(x_0) \end{cases} .$$

One verifies that $\varphi_n' = \xi_n e^{i\psi_n}$ in $B_{\rho-s}(x_0)$ is such that

$$\int_{B_{\rho-s}(x_0)} |\nabla \varphi_n'|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n'|^2)^2 \leq C(r, s) \tag{6. 52}$$

Combining (6. 49) and (6. 52) we have extended φ_n by φ'_n in such a way that

$$\int_{B_\rho(x_0)} |\nabla \varphi'_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 \leq C s \log \frac{1}{\varepsilon_n} + C(r, s) \quad (6. 53)$$

Since h_n weakly converges in $W^{1,p}(B_\rho(x_0))$ for any $p < \frac{3}{2}$, h_n strongly converges in $L^2(B_\rho(x_0))$ thus we have

$$\frac{1}{\rho} \int_{B_\rho(x_0)} |\nabla_{\tilde{A}_n} \varphi'_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi'_n|^2)^2 + |h_n|^2 \leq C \frac{s}{\rho} \log \frac{1}{\varepsilon_n} + C(r, s) \quad .$$

Choose now s small enough such that $\frac{Cs}{\rho} < \frac{\eta}{2}$ then we obtain that, for ε_n sufficiently small,

$$\frac{1}{\rho} \int_{B_\rho(x_0)} |\nabla_{\tilde{A}_n} \varphi_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 + |h_n|^2 < \eta \log \frac{\rho}{\varepsilon_n}$$

This implies by Lemma 3.1 that $|\varphi_n| \geq \frac{1}{2}$ in $B_{r/4}(x_0) \subset B_{\rho/2}(x_0)$, and h_n strongly converges in $H^1(B_{r/4}(x_0))$. ■

7. THE LIMITING SINGULAR SET IDENTIFIED AS A MINIMAL CONNECTION

In this part we conclude the proof of the theorem 1.2 stated in the introduction by proving the following theorem

THEOREM 7.1. The 1-dimensional rectifiable current \mathbb{L} given by the theorem 5.1 is a minimal connection between the $(x_i, d_i)_{i=1\dots n}$. ■

Proof of theorem 7.1.

We denote by M the rectifiable set associated to \mathbb{L} , by $\xi(y)$ the approximate tangent vector (with the orientation given by \mathbb{L}) at $y \in M$ and by $\theta(y)$ the positive integer multiplicity of \mathbb{L} at the same point. The rectifiability of M implies that it can be included in a countable union of 1-dimensional disjoint submanifolds of \mathbb{R}^3 (i.e $M \subset \cup_{i=1}^\infty \mathcal{C}_i$) modulo an \mathcal{H}^1 - measure zero set that we do not consider in working with \mathbb{L} . We have, using the notation $M_i = M \cap \mathcal{C}_i$

$$L = M(\mathbb{L}) = \int_M \theta \, d\mathcal{H}^1 = \sum_{i=1}^\infty \int_{M \cap \mathcal{C}_i} \theta \, d\mathcal{H}^1 \quad .$$

Let $\alpha > 0$, we claim that for n sufficiently large we have

$$\int_{M_{\tilde{x}}} |\nabla_{\tilde{A}_n} \varphi_n|^2 + \frac{1}{2\varepsilon_n^2} (1 - |\varphi_n|^2)^2 + |h_n|^2 \geq 2\pi(L - \alpha) \log \frac{1}{\varepsilon_n} - C(\alpha) \quad (7. 1)$$

where $C(\alpha)$ does not depend on n .

Let $N \in \mathbb{N}$ be such that

$$\sum_{i \geq N} \int_{M_i} \theta \, d\mathcal{H}^1 < \alpha \quad . \quad (7. 2)$$

For $i \leq N$ and $z \in \mathcal{C}_i$ denote by $N_i(z)$ the plane passing by z perpendicular to \mathcal{C}_i and by $D_i(z, s)$ the two dimensional disk in this plane of centre z and

radius s . z also denotes sometimes the normal coordinate along \mathcal{C}_i .

Let $i \leq N$ be given. For $s > 0$ sufficiently small the $D_i(z, s)$ for $z \in \mathcal{C}_i$ are disjoint and $\cup_{z \in \mathcal{C}_i} D_i(z, s)$ is a regular open tubular neighborhood of \mathcal{C}_i denoted by $U_i(s)$. One verifies, using the coarea formula, that

$$\forall f \in L^1(U_i(s)) \tag{7.3}$$

$$\int_{U_i(s)} f = \int_{\mathcal{C}_i} dz \int_{D_i(z,s)} f(x) |1 - C_i(z)\nu(z) \cdot (x - z)| d\mathcal{H}^2$$

where $C_i(z)$ is the curvature of \mathcal{C}_i at z and $\nu(z)$ is first normal of \mathcal{C}_i at z when $C_i(z) \neq 0$.

Choose s sufficiently small such that

$$\sup_{i \leq N} \left[\sup_{z \in \mathcal{C}_i} C_i(z) \right] \times s \leq \alpha \tag{7.4}$$

Let $\mathcal{A}_i^N(s)$ for $i \leq N$ be the following subset of \mathcal{C}_i

$$\mathcal{A}_i^N(s) = \{z \in \mathcal{C}_i / \exists j \neq i, j \leq N \text{ s.t. } D_i(z, s) \cap \cup_j(s) \neq \emptyset\}$$

Since the $\mathcal{C}_i, i \leq N$ are disjoint submanifolds of \mathbb{R}^3 , for s chosen sufficiently small we can ensure

$$\forall i \leq N \int_{\mathcal{C}_i \cap \mathcal{A}_i^N(s)} \theta d\mathcal{H}^1 \leq \frac{\alpha}{N} \tag{7.5}$$

Denote by $\mathcal{B}_i^N(s)$, for $i \leq N$, the following subset of \mathcal{C}_i

$$\mathcal{B}_i^N(s) = \{z \in \mathcal{C}_i \setminus \mathcal{A}_i^N(s) / \exists j > N, \exists x \in M_j \cap D_i(z, s) \text{ s.t. } \theta(x) \neq 0\}$$

For $i \leq N$ and $j > N$ we consider the set

$$R_i^j(s) = M_j \cap U_i(s) \setminus \left[U_i(s) \cap \bigcup_{k \neq i, k \leq N} U_k(s) \right] .$$

Clearly the $R_i^j(s)$ are all disjoint. Moreover we have

$$\mathcal{H}^1(\mathcal{B}_i^N(s)) \leq \sum_{j > N} \int_{R_i^j(s)} \theta d\mathcal{H}^1 \tag{7.6}$$

Since $\sum_{i \leq N} \int_{\mathcal{C}_i} \theta d\mathcal{H}^1 < +\infty$, there exists $\beta > 0$ independent of $i \leq N$ such that

$$\forall i \leq N \quad \mathcal{H}^1(\mathcal{B}_i^N(s)) \leq \beta \Rightarrow \int_{\mathcal{B}_i^N(s)} \theta d\mathcal{H}^1 \leq \frac{\alpha}{N} . \tag{7.7}$$

We claim that, for s sufficiently small, $\mathcal{H}^1(\mathcal{B}_i^N(s)) \leq \beta$ for $i \leq N$. Indeed, let N' such that

$$\sum_{j > N'} \int_{M_j} \theta d\mathcal{H}^1 \leq \frac{\beta}{2N} , \tag{7.8}$$

this implies clearly that

$$\forall i \leq N \quad \sum_{j > N'} \int_{R_i^j(s)} \theta d\mathcal{H}^1 \leq \frac{\beta}{2N} . \tag{7.9}$$

Fix N' such that (7. 8) holds. Since the $\mathcal{C}_j, N < j \leq N'$ are regular curves disjoint from the $\mathcal{C}_i \ 1 \leq i \leq N$ for s sufficiently small $\sum_{j=N+1}^{N'} \mathcal{H}^1(R_i^j(s))$ is as small as we want for any $i \leq N$ thus we can ensure

$$\forall i \leq N \quad \sum_{j=N+1}^{N'} \int_{R_i^j(s)} \theta d\mathcal{H}^1 \leq \frac{\beta}{2N} \quad . \quad (7. 10)$$

Combining (7. 9), (7. 10), (7. 6) and (7. 7) for s sufficiently small we have

$$\forall i \leq N \quad \int_{B_i^N(s)} \theta d\mathcal{H}^1 \leq \frac{\alpha}{N} \quad . \quad (7. 11)$$

Fix $s > 0$ such that (7. 5) and (7. 11) hold. For $t > 0$ to be chosen sufficiently small bellow, consider the lattice $t\mathbb{Z}^3$ in \mathbb{R}^3 except the points included in $\cup_{i=1}^n B(x_i, \frac{t}{\eta})$, where η is the constant given by Lemma 3.1. Because of Lemma 3.1 we can put the bad set out of $\cup_{i=1}^n B(x_i, \frac{t}{\eta})$ (i.e., $y \notin \cup_{i=1}^n B(x_i, \frac{t}{\eta})$) such that $|\varphi_n|(y) \leq \frac{1}{2}$ in C/t cubes $C(2t, y_l)$ where C does not depend neither on n nor on t and where the y_l are points of the lattice.

Extract a subsequence such that the bad cubes $C(2t, y_l)$ are always the same in $\cup_{i=1}^N U_i(s)$.

First of all for t sufficiently small we can ensure

$$\int_{M \cap \cup_{i=1}^n B(x_i, t/\eta)} \theta d\mathcal{H}^1 \leq \alpha \quad . \quad (7. 12)$$

Denote by $N_t^i(z)$, for $i \leq N$ and $z \in \mathcal{C}_i$, the cardinal of the following set

$$N_t^i(z) = \# \{y_l / C(2t, y_l) \cap D_i(z, s) \neq \emptyset\}$$

Since a cube $C(2t, y_l)$ can intersect $D_i(z, s)$ for z in an interval of size at most $4t$ we obtain that

$$\int_{\mathcal{C}_i} N_t^i(z) d\mathcal{H}^1 \leq C^i \quad , \quad (7. 13)$$

where C^i does not depend on t . Consider β such that

$$\forall E \subset \mathcal{C}_i \quad \mathcal{H}^1(E) \leq \beta \Rightarrow \int_{\mathcal{C}_i \cap E} \theta d\mathcal{H}^1 \leq \frac{\alpha}{N} \quad . \quad (7. 14)$$

Since C^i does not depend on t , there exists $N^i \in \mathbb{N}$ independent on t such that

$$\forall i \leq N \quad \mathcal{H}^1(\{z \in \mathcal{C}_i \ s.t. \ N_t^i(z) \geq N^i\}) \leq \beta \quad (7. 15)$$

where β is given by (7. 14) independently of $i \leq N$.

Let $\bar{N} = \max_{i \leq N} N^i$. Denote by $\mathcal{C}_i^N(s, t)$ the set

$$\mathcal{C}_i^N(s, t) = \{z \in \mathcal{C}_i \ s.t. \ N_t^i(z) \geq \bar{N}\} \quad .$$

For t verifying

$$0 < t < \frac{s}{4\bar{N} + 4} \quad (7. 16)$$

we can ensure that there exists $k_z \in [1; 2\bar{N} + 1]$ for any $z \in \mathcal{C}_i \setminus \mathcal{C}_i^N(s, t)$ such that the annulus

$$L_i(z, s) = D_i \left(z; \frac{k_z + 1}{2\bar{N} + 2} s \right) \setminus D_i \left(z; \frac{k_z}{2\bar{N} + 2} s \right)$$

does not intersect the bad set (i.e. $\cup_{l \leq \frac{\varepsilon}{t}} C(2t, y_l)$). Fix t verifying (7. 12) and (7. 16). From part 6 we know that φ_n tends strongly in H^1 to φ_* in the complement of the bad set $\cup_{l \leq \frac{\varepsilon}{t}} C(2t, y_l)$ in $\cup_{i \leq N} U_i(s)$. Thus for n sufficiently large we have

$$\forall i \leq N \quad \forall z \in \mathcal{C}_i \setminus \mathcal{C}_i^N(s, t)$$

$$\text{deg} \left(\frac{\varphi_n}{|\varphi_n|}; \partial D_i \left(\frac{k_z}{2\bar{N} + 2} s, z \right) \right) = \text{deg} \left(\varphi_*; \partial D_i \left(\frac{k_z}{2\bar{N} + 2} s, z \right) \right) \quad (7. 17)$$

Moreover using Lemma A.5 we have

$$\left\{ \begin{array}{l} \forall i \leq N, \quad \mathcal{H}^1 - a.e. z \in \mathcal{C}_i \setminus \mathcal{A}_i^N \cup \mathcal{B}_i^N \cup \mathcal{C}_i^N \\ \text{deg} \left(\varphi_*; \partial D_i \left(\frac{k_z s}{2\bar{N} + 2}, z \right) \right) = -\theta(z) \end{array} \right. \quad (7. 18)$$

Let $\alpha' = 1 - \alpha$, using Lemma 6.1 the bad set in $\cup_{l \leq \frac{\varepsilon}{t}} C(2t, y_l)$ can be put in $C(\alpha')/\varepsilon_n^{\alpha'}$ disjoint open sets of size $\approx \varepsilon_n^{\alpha'}$ verifying conditions 1, 2, ...,6. Denote by $P_{\alpha'}^i(z)$ for $i \leq N$ and $z \in \mathcal{C}_i$ the cardinal of the following set

$$P_{\alpha'}^i(z) = \# \{R_l / R_l \cap D_i(z, s) \neq \emptyset\} \quad .$$

Since each R_l given by Lemma 6.1 for any α' can intersect $D_i(z, s)$ for z in an interval of size at most $4\varepsilon^{\alpha'}$ we obtain that

$$\int_{\mathcal{C}_i} P_{\alpha'}^i(z) d\mathcal{H}^1 \leq C^i(\alpha') \quad (7. 19)$$

where $C^i(\alpha')$ does not depend on ε_n . Consider once again β given by (7. 14), there exists $P_{\alpha'}^i$, independent of ε_n such that

$$\forall i \leq N \quad \mathcal{H}^1 \left(\{z \in \mathcal{C}_i s.t. P_{\alpha'}^i(z) \geq P_{\alpha'}^i\} \right) \leq \beta \quad (7. 20)$$

Let $\bar{P}_{\alpha'} = \max_{i \leq N} P_{\alpha'}^i$ and denote by $\mathcal{D}_i^N(s)$ the set

$$\mathcal{D}_i^N(s) = \{z \in \mathcal{C}_i s.t. P_{\alpha'}^i(z) \geq \bar{P}_{\alpha'}\} \quad .$$

Using Lemma A.6 for $\mathcal{H}^1 a.e. z \in \mathcal{C}_i \setminus \mathcal{A}_i^N \cup \mathcal{B}_i^N \cup \mathcal{C}_i^N \cup \mathcal{D}_i^N$ we have in any trivialisation of E over $D_i(z, s)$

$$\int_{D_i \left(\frac{k_z + 1}{2\bar{N} + 2} s, z \right) \setminus \cup R_l} \left| \nabla \frac{\varphi_n}{|\varphi_n|} \right|^2 \geq 2\pi \theta(z) \alpha' \log \frac{1}{\varepsilon_n} - 2\pi \log \left(\frac{s}{2\bar{N} + 2} \right) - 2\pi (\bar{P}_{\alpha'})^5 \quad (7. 21)$$

where d is the maximal degree of $\varphi_n/|\varphi_n|$ around each R_l and is uniformly bounded independently of ε_n because of conditions 4 and 5 of Lemma 4.1.

By lowering η in Lemma 3.1 we can impose that $|\varphi_n| \geq \alpha'$ in the complement of the R_i . Thus we have

$$\mathcal{H}^1 - a.e.z \in \mathcal{C}_i \setminus \mathcal{A}_i^N \cup \mathcal{B}_i^N \cup \mathcal{C}_i^N \cup \mathcal{D}_i^N$$

$$\int_{D_i(z,s)} |\nabla \varphi_n|^2 \geq 2\pi\theta(z)\alpha'^3 \log \frac{1}{\varepsilon_n} - 2\pi\alpha'^2 \log \left(\frac{s}{2N+2} \right) - 2\pi(\bar{P}_{\alpha'})^5 \tag{7.22}$$

Since we are working in $\cup_{i=1}^N U_i(s) \setminus \cup_{i=1}^n B(x_i, \frac{t}{\eta})$ we have the strong convergence $\tilde{A}_n \rightarrow \tilde{A}$ in L^2 thus

$$\int_{\cup_{i=1}^N U_i(s) \setminus \cup_{i=1}^n B(x_i, t/\eta)} |\tilde{A}_n|^2 \leq C \tag{7.23}$$

where C does not depend on ε_n . Integrating (7.21) along the $\mathcal{C}_i \setminus \mathcal{A}_i^N \cup \mathcal{B}_i^N \cup \mathcal{C}_i^N \cup \mathcal{D}_i^N$ using (7.3) and (7.22), since we have the bound using (7.5), (7.11), (7.15) and (7.20)

$$\int_{\mathcal{A}_i^N \cup \mathcal{B}_i^N \cup \mathcal{C}_i^N \cup \mathcal{D}_i^N} \theta d\mathcal{H}^1 \leq \frac{4\alpha}{N}$$

we get, using also (7.12),

$$\int_{\cup_{i=1}^N U_i(s) \setminus \cup_{i=1}^n B(x_i, t/\eta)} |\nabla_{\tilde{A}} \varphi|^2 \geq$$

$$2\pi\alpha'^3 \left(\int_{\cup_{i=1}^N M_i} \theta d\mathcal{H}^1 - 5\alpha \right) \log \frac{1}{\varepsilon_n} - C(\alpha) \tag{7.24}$$

where $C(\alpha)$ does not depend on N . Since $\alpha' = 1 - \alpha$ considering (7.2), (7.23) implies (7.1).

Combining the upper bound of the total energy given by Lemma 2.1 and (5.47) which says that \mathbb{L} is a connection between the $(x_i, d_i)_{i=1 \dots n}$, the lower bound (7.1) implies that \mathbb{L} is a minimal connection. ■

APPENDIX A. APPENDIX

Lemma A.1. *Let φ, η, ξ be 0, 1 and 2-forms of $B_1 \subset \mathbb{R}^3$ we have*

$$\frac{\partial}{\partial r} d\varphi = d \frac{\partial \varphi}{\partial r} - \frac{d\varphi}{r} + \frac{\partial \varphi}{\partial r} d(\log r)$$

$$\frac{\partial}{\partial r} d\eta = d \frac{\partial \eta}{\partial r} - \frac{d\eta}{r} + d(\log r) \wedge \frac{\partial \eta}{\partial r}$$

and

$$\frac{\partial}{\partial r} d\xi = d \frac{\partial \xi}{\partial r} - \frac{d\xi}{r} + d(\log r) \wedge \frac{\partial \xi}{\partial r}$$

■

Proof of Lemma A.1.

$$\begin{aligned} \frac{\partial}{\partial r}d\varphi &= \sum_i \frac{\partial}{\partial r} \frac{\partial \varphi}{\partial x_i} dx_i = \sum_{ik} \frac{\partial}{\partial x_k} \frac{\partial \varphi}{\partial x_i} \frac{x_k}{|x|} dx_i \\ &= \sum_{ik} \frac{\partial}{\partial x_i} \frac{\partial \varphi}{\partial x_k} \frac{x_k}{|x|} dx_i \\ &= \sum_i \frac{\partial}{\partial x_i} \frac{\partial \varphi}{\partial r} dx_i - \sum_{ik} \frac{\partial \varphi}{\partial x_k} \frac{\partial}{\partial x_i} \left(\frac{x_k}{|x|} \right) dx_i \\ &= d \left(\frac{\partial \varphi}{\partial r} \right) - \frac{1}{r} d\varphi + \frac{\partial \varphi}{\partial r} d(\log r). \end{aligned}$$

The two other equalities are proved in the same way by direct computations. ■

Lemma A.2. *Let φ and η be 0 and 1 forms in $B_1 \subset \mathbb{R}^3$. for $r > 0$ let d_\top and $(*)$ the exterior derivative and the Hodge operator on $\partial B_r \subset \mathbb{R}^3$ and let Δ_r be the Laplace-Beltrani operator on ∂B_r . We have*

$$\begin{aligned} d_\top \frac{\partial \varphi}{\partial r} &= \frac{\partial}{\partial r} d_\top \varphi + \frac{1}{r} d_\top \varphi \\ d_\top \frac{\partial \varphi}{\partial r} &= \frac{\partial}{\partial r} d_\top \eta + \frac{1}{r} d_\top \eta \end{aligned}$$

Let i_r the isometric embedding in $i_r : \partial B_r \rightarrow \mathbb{R}^3$

$$\frac{\partial}{\partial r} (*) i_r^\# \eta = (*) i_r^\# \frac{\partial \eta}{\partial r}$$

and

$$\frac{\partial}{\partial r} \Delta_r \varphi - \Delta_r \frac{\partial \varphi}{\partial r} = -\frac{2}{r} \Delta_r \varphi. \quad \blacksquare$$

Proof of Lemma A.2.

By definition

$$d_\top \frac{\partial \varphi}{\partial r} = d \frac{\partial \varphi}{\partial r} - \frac{\partial}{\partial r} \frac{\partial \varphi}{\partial r} dr \quad ,$$

thus by Lemma A.1

$$\begin{aligned} d_\top \frac{\partial \varphi}{\partial r} &= \frac{\partial}{\partial r} d\varphi + \frac{d\varphi}{r} - \frac{\partial \varphi}{\partial r} \frac{dr}{r} - \frac{\partial^2 \varphi}{\partial r^2} dr \\ &= \frac{\partial}{\partial r} \left(d_\top \varphi + \frac{\partial \varphi}{\partial r} dr \right) + \frac{d\varphi}{r} - \frac{\partial \varphi}{\partial r} \frac{dr}{r} - \frac{\partial^2 \varphi}{\partial r^2} dr \\ &= \frac{\partial}{\partial r} (d_\top \varphi) + \frac{d\varphi}{r} - \frac{\partial \varphi}{\partial r} \frac{dr}{r} \\ &= \frac{\partial}{\partial r} d_\top \varphi + \frac{1}{r} d_\top \varphi \end{aligned}$$

Let

$$\omega = *dr = \frac{xy \wedge dz + ydz \wedge dx + zdx \wedge dy}{r}$$

By definition

$$d_{\top} \frac{\partial \eta}{\partial r} = \left\langle d \frac{\partial \eta}{\partial r}, \omega \right\rangle \omega,$$

thus

$$\begin{aligned} d_{\top} \frac{\partial \eta}{\partial r} &= \left\langle \frac{\partial}{\partial r} d\eta + \frac{d\eta}{r} - \frac{dr}{r} \wedge \frac{\partial \eta}{\partial r}, \omega \right\rangle \omega \\ &= \left\langle \frac{\partial}{\partial r} d\eta + \frac{d\eta}{r}, \omega \right\rangle \omega \end{aligned}$$

Since $\frac{\partial \omega}{\partial r} = 0$ we have

$$\begin{aligned} d_{\top} \frac{\partial \eta}{\partial r} &= \frac{\partial}{\partial r} [\langle d\eta, \omega \rangle \omega] + \frac{1}{r} \langle d\eta, \omega \rangle \omega \\ &= \frac{\partial}{\partial r} (d_{\top} \eta) + \frac{1}{r} d_{\top} \eta. \end{aligned}$$

We have $(*)i_r^{\#} \eta = (*)(\eta \wedge dr)$, thus

$$\begin{aligned} \frac{\partial}{\partial r} (*)i_r^{\#} \eta &= * \frac{\partial}{\partial r} (\eta \wedge dr) = * \left(\frac{\partial \eta}{\partial r} \wedge dr \right) \\ &= (*)i_r^{\#} \frac{\partial \eta}{\partial r}. \end{aligned}$$

Finally

$$\begin{aligned} \Delta_r \frac{\partial \varphi}{\partial r} &= (*)d_{\top} (*)d_{\top} \frac{\partial \varphi}{\partial r} = (*)d_{\top} (*) \left[\frac{\partial}{\partial r} d_{\top} \varphi + \frac{1}{r} d_{\top} \varphi \right] \\ &= (*)d_{\top} \frac{\partial}{\partial r} [(*)d_{\top} \varphi] + \frac{1}{r} \Delta_r \varphi \\ &= (*) \frac{\partial}{\partial r} [d_{\top} (*)d_{\top} \varphi] + \frac{1}{r} (*)d_{\top} (*)d_{\top} \varphi + \frac{1}{r} \Delta_r \varphi \\ &= \frac{\partial}{\partial r} \Delta_r \varphi + \frac{2}{r} \Delta_r \varphi \end{aligned}$$

■

Lemma A.3. *Let φ be a 0-form on $B_1 \subset \mathbb{R}^3$ taking its values into \mathbb{C} , we have*

$$\Delta_r \left(i\varphi, \frac{\partial \varphi}{\partial r} \right) - \frac{\partial}{\partial r} (i\varphi, \Delta_r \varphi) = 2(*)d_{\top} \left[\left(\frac{\partial \varphi}{\partial r}, i(*)d_{\top} \varphi \right) \right] + \frac{2}{r} (i\varphi, \Delta_r \varphi)$$

■

Proof of Lemma A.3.

$$\begin{aligned}
 \Delta_r \left(i\varphi, \frac{\partial\varphi}{\partial r} \right) &= (*)d_{\top}(*)d_{\top}(i\varphi, \frac{\partial\varphi}{\partial r}) \\
 &= (*)d_{\top}(*) \left[\left(\frac{\partial\varphi}{\partial r}, i d_{\top}\varphi \right) + \left(i\varphi, d_{\top} \frac{\partial\varphi}{\partial r} \right) \right] \\
 &= (*) \left(id_{\top}(*)d_{\top}\varphi, \frac{\partial\varphi}{\partial r} \right) + (*) \left(d_{\top} \frac{\partial\varphi}{\partial r}, i(*)d_{\top}\varphi \right) \\
 &+ (*) \left(id_{\top}\varphi, (*)d_{\top} \frac{\partial\varphi}{\partial r} \right) + (*) \left(i\varphi, d_{\top}(*)d_{\top} \frac{\partial\varphi}{\partial r} \right)
 \end{aligned}$$

where we have used the following notation, if a and b are two \mathbb{C} 1-form,

$$(a, b) = \frac{1}{2}[a \wedge \bar{b} + \bar{a} \wedge b],$$

and since $((*)a) \wedge b = ((*)b) \wedge a$ we have

$$((*)a, b) = ((*)b, a).$$

This implies

$$\begin{aligned}
 \Delta_r \left(i\varphi, \frac{\partial\varphi}{\partial r} \right) &= \left(i\Delta_r\varphi, \frac{\partial\varphi}{\partial r} \right) + 2(*) \left(d_{\top} \frac{\partial\varphi}{\partial r}, i(*)d_{\top}\varphi \right) \\
 &+ \left(i\varphi, \Delta_r \frac{\partial\varphi}{\partial r} \right) \quad ,
 \end{aligned}$$

and finally

$$\begin{aligned}
 \Delta_r \left(i\varphi, \frac{\partial\varphi}{\partial r} \right) &= \left(i\Delta_r\varphi, \frac{\partial\varphi}{\partial r} \right) + 2(*)d_{\top} \left[\left(\frac{\partial\varphi}{\partial r}, i(*)d_{\top}\varphi \right) \right] \\
 &- 2 \left(\frac{\partial\varphi}{\partial r} i\Delta_r\varphi \right) + \left(i\varphi, \frac{\partial}{\partial r} \Delta_r\varphi \right) + \left(i\varphi, \frac{2}{r} \Delta_r\varphi \right) \\
 &= \frac{\partial}{\partial r} \left(i\varphi, \Delta_r\varphi \right) - \left(i \frac{\partial\varphi}{\partial r}, \Delta_r\varphi \right) + \left(i\Delta_r\varphi, \frac{\partial\varphi}{\partial r} \right) \\
 &- 2 \left(\frac{\partial\varphi}{\partial r}, i\Delta_r\varphi \right) + 2(*)d_{\top} \left[\left(\frac{\partial\varphi}{\partial r}, i(*)d_{\top}\varphi \right) \right] \\
 &+ \frac{2}{r} \left(i\varphi, \Delta_r\varphi \right).
 \end{aligned}$$

This yields the desired result. ■

Lemma A.4. *Let f and η be 0 and 1-forms in $B_1 \subset \mathbb{R}^3$ and $\omega = *dr$, we have*

$$\langle \Delta(f\omega); \omega \rangle = \Delta f - \frac{2}{r^2} f$$

and

$$d_{\top}\Delta\eta = \Delta d_{\top}\eta + \frac{4}{r^2} \frac{\partial}{\partial r} [r \langle d\eta; \omega \rangle] \omega - \frac{2}{r} * d \langle d\eta; \omega \rangle$$



Proof of Lemma A.4.

We have

$$\Delta(f\omega) = \Delta f \omega + 2 \sum_k \frac{\partial f}{\partial k} \frac{\partial \omega}{\partial k} + f \Delta \omega.$$

Moreover one verifies that $\Delta \omega = -\frac{2}{r^2} \omega$ and in the other hand $\forall k < \frac{\partial \omega}{\partial k}; \omega \rangle = 0$ thus

$$\langle \Delta(f\omega); \omega \rangle = \Delta f - \frac{2}{r^2} f \quad .$$

By definition

$$d_{\top} \Delta \eta = \langle d\Delta \eta; \omega \rangle \omega = \langle \Delta d\eta, \omega \rangle \omega \quad .$$

In the other hand

$$\begin{aligned} \Delta d_{\top} \eta &= \Delta [\langle d\eta; \omega \rangle \omega] = \sum_i \Delta \left[\langle d\eta; \omega \rangle \frac{x_i}{r} \right] dx_{i+1} \wedge dx_{i-1} \\ &= 2 \sum_i \nabla \langle d\eta; \omega \rangle \nabla \left(\frac{x_i}{r} \right) dx_{i+1} \wedge dx_{i-1} \\ &+ \Delta \langle d\eta; \omega \rangle \omega - \frac{2}{r^2} \langle d\eta; \omega \rangle \omega \quad . \end{aligned} \tag{A. 1}$$

We have also

$$\begin{aligned} \Delta \langle d\eta; \omega \rangle &= \langle d\Delta \eta; \omega \rangle + 2 \sum_i \nabla (d\eta)_i \nabla \left(\frac{x_i}{r} \right) + \langle d\eta; \Delta \omega \rangle \\ &= \langle d\Delta \eta; \omega \rangle + 2 \sum_{il} \frac{\partial}{\partial l} (d\eta)_i \frac{\partial}{\partial l} \left(\frac{x_i}{r} \right) - \frac{2}{r^2} \langle d\eta; \omega \rangle \\ &= \langle d\Delta \eta; \omega \rangle + 2 \sum_i \frac{\partial}{\partial i} (d\eta)_i \frac{1}{r} + 2 \sum_{il} \frac{\partial}{\partial l} (d\eta)_i x_i \frac{\partial}{\partial l} \left(\frac{1}{r} \right) \\ &- \frac{2}{r^2} \langle d\eta; \omega \rangle \end{aligned} \tag{A. 2}$$

Clearly $\sum_i \frac{\partial (d\eta)_i}{\partial i} = *d(d\eta) = 0$, and since $\frac{\partial}{\partial r} \left(\frac{x_i}{r} \right) = 0$

$$\begin{aligned} -2 \sum_{il} \frac{\partial}{\partial l} (d\eta)_i x_i \frac{x_l}{r^3} &= -\frac{2}{r} \sum_l \frac{x_l}{r} \frac{\partial}{\partial l} \left(\sum_i (d\eta)_i \frac{x_i}{r} \right) \\ &= -\frac{2}{r} \frac{\partial}{\partial r} \langle d\eta; \omega \rangle. \end{aligned}$$

This implies

$$\begin{aligned} \Delta \langle d\eta; \omega \rangle &= \langle d\Delta \eta; \omega \rangle - \frac{2}{r} \frac{\partial}{\partial r} \langle d\eta; \omega \rangle - \frac{2}{r^2} \langle d\eta; \omega \rangle \\ &= \langle d\Delta \eta; \omega \rangle - \frac{2}{r^2} \frac{\partial}{\partial r} [r \langle d\eta; \omega \rangle]. \end{aligned} \tag{A. 3}$$

Finally

$$\begin{aligned} \nabla \langle d\eta; \omega \rangle \cdot \nabla \left(\frac{x_i}{r} \right) &= \frac{1}{r} \frac{\partial}{\partial x_i} \langle d\eta; \omega \rangle \\ &- \frac{x_i}{r^2} \frac{\partial}{\partial r} \langle d\eta; \omega \rangle \quad . \end{aligned} \tag{A. 4}$$

Combining (A. 1), (A. 3) and (A. 4) we have

$$\begin{aligned} \Delta d_{\top} \eta &= \langle d\Delta\eta; \omega \rangle \omega - \frac{2}{r^2} \frac{\partial}{\partial r} [r \langle d\eta; \omega \rangle] \omega \\ &+ \frac{2}{r} * d \langle d\eta; \omega \rangle - \frac{2}{r} \frac{\partial}{\partial r} [\langle d\eta; \omega \rangle] \omega \\ &- \frac{2}{r^2} \langle d\eta; \omega \rangle \omega. \end{aligned}$$

This yields the desired result. ■

Lemma A.5. *Let Ω and ω be two bounded regular subsets of \mathbb{R}^2 such that Ω is simply connected and $\omega \subset\subset \Omega$. Let $a > 0$ and denote by C and \mathcal{C} the following cylinders of \mathbb{R}^3 .*

$$C = \omega \times (-a, a) \quad \text{and} \quad \mathcal{C} = \Omega \times (-a, a).$$

Let $T(\theta, M)$ be a 1-dimensionnal, finite mass, integer multiplicity rectifiable current with support included in C and let ψ be $W^{1,1}(\mathcal{C}, S^1)$ map such that ψ is regular in $\mathcal{C} \setminus C$ and

$$*d(i\psi, d\psi) = 2\pi \mathbb{T} \quad \text{in} \quad \mathcal{D}_1(\mathcal{C}) \tag{A. 5}$$

Let Γ be a closed regular curve in $\Omega \setminus \omega$ and Σ in Ω such that $\partial\Gamma = \Sigma$. Denote by $\Gamma_t, \Sigma_t, \omega_t$ and Ω_t the sets $\Gamma \times \{t\}, \Sigma \times \{t\}, \omega \times \{t\}$ and $\Omega \times \{t\}$ for $t \in (-a, a)$ then we have

$$\begin{aligned} \text{for a.e. } t \in (-a, a) \quad \text{deg}_{\Gamma_0} \psi &= \text{deg}_{\Gamma_t} \psi \\ &= - \sum_{y \in M \cap \Sigma_t} \text{sgn} \langle \xi; e_3 \rangle (y) \theta(y) \end{aligned} \tag{A. 6}$$

where $e_3 = (0, 0, 1)$ and Γ_t is oriented by e_3 . ■

Proof of Lemma A.5.

First of all we claim that

$$\begin{aligned} \text{for a.e } t \in (-a, a) \quad \sum_{y \in M \cap \omega_t} \theta(y) &< +\infty \\ \text{and} \quad \forall y \in M \cap \omega_t \quad \langle \xi(y); e_3 \rangle &\neq 0 \end{aligned} \tag{A. 7}$$

This can be proved using the coarea formula of Federer (see for instance [9] pages 15 and 34). Let t be the function on M which associates the third coordinate. This is a Lipschitz function in M (which is rectifiable subset of

\mathbb{R}^3) moreover for \mathcal{H}^1 a.e. $y \in M$ $apJ_1 t(y) = | \langle \xi(y); e_z \rangle |$ thus for any \mathcal{H}^1 LM integrable function g on M we have

$$\begin{aligned} \int_M | \langle \xi(y), e_3 \rangle | g d\mathcal{H}^1 &= \int_{-a}^a \int_{M \cap \omega_t} g d\mathcal{H}^0 dt \\ &= \int_{-a}^a \sum_{y \in M \cap \omega_t} g(y) dt \end{aligned} \tag{A. 8}$$

Take first $g = \theta$, since $M(\mathbb{T}) < +\infty$ we have

$$\int_{-a}^a \sum_{y \in M \cap \omega_t} \theta(y) dt \leq M(\mathbb{T}) \quad . \tag{A. 9}$$

For any $\alpha > 0$ take $g_\alpha = \frac{\theta(y)}{\alpha + | \langle \xi(y), e_3 \rangle |}$, we have

$$\int_{-a}^{+a} \sum_{y \in M \cap \omega_t} \frac{\theta(y)}{\alpha + | \langle \xi(y), e_3 \rangle |} \leq M(\mathbb{T}) \quad . \tag{A. 10}$$

Thus, let

$$E = \left\{ \begin{array}{l} t \in (-a, a) \text{ s.t. } \sum_{y \in M \cap \omega_t} \theta(y) < +\infty \text{ and} \\ \exists y \in M \cap \omega_t \text{ s.t. } \theta(y) \neq 0 \text{ and } \langle \xi(y), e_3 \rangle = 0 \end{array} \right\}$$

from (A. 10) we deduce that

$$\forall \alpha > 0 \quad |E| \times \frac{1}{\alpha} \leq M(\mathbb{T}) \tag{A. 11}$$

This implies (A. 7).

Since Γ_t is homotopic to Γ_0 in $\mathcal{C} \setminus C$ we have clearly

$$deg_{\Gamma_t} \psi = deg_{\Gamma_0} \psi \quad . \tag{A. 12}$$

Let $t_0 \in (-a, a)$ such that (A. 7) holds for $t = t_0$ and let U be a tubular neighborhood of Γ chosen sufficiently small such that

$$U \subset \mathcal{C} \setminus C$$

and U is diffeomorphic to a neighborhood of the zero section of the normal bundle $N\Gamma$ of $\Gamma \subset \mathcal{C}$.

We can suppose that this diffeomorphism sends linearly each intersection of U and any normal plane in \mathcal{C} to Γ (at γ) to a neighborhood of 0 in the corresponding fiber ($\pi^{-1}\gamma = N_\gamma\Gamma$) of $N\Gamma$.

Let χ be the following function on $\Omega \times \{t_0\}$

$$\begin{cases} 0 \leq \chi \leq 1 \\ \chi \equiv 0 \text{ in } \Omega \times \{t_0\} \setminus \Sigma \times \{t_0\} \\ \chi \equiv 1 \text{ in } \omega \times \{t_0\} \\ \text{supp}(d\chi) \subset U \cap \Omega \times \{t_0\}, \end{cases} \tag{A. 13}$$

and define on all of \mathcal{C} $\chi(x, y, t) \equiv \chi(x, y, t_0)$.

Let H be a non negative function of $C_c^\infty(\mathbb{R}; \mathbb{R}_+)$ such that

$$\begin{cases} H \equiv 1 \text{ in a neighborhood of } 0 \\ \int_{\mathbb{R}} H = 1 \end{cases} \tag{A. 14}$$

Let $H_\lambda^{t_0} = \frac{1}{\lambda} H\left(\frac{t-t_0}{\lambda}\right)$ for $\lambda > 0$ chosen sufficiently small

$$supp(H_\lambda^{t_0} \times d\chi \wedge dt) \subset U$$

Let

$$\eta_{\Gamma_{t_0}} = H_\lambda^{t_0} d\chi \wedge dt \tag{A. 15}$$

We clearly have $d\eta_{\Gamma_{t_0}} = 0$, moreover let $\gamma \in \Gamma_{t_0}$ and let P be the vertical plane (containing e_3) passing by γ , perpendicular to Γ_{t_0} and oriented by the unit tangent vector $\dot{\gamma}$ of Γ_{t_0} (oriented by e_3 i.e. such that $\dot{\gamma} \wedge e_3$ is the unit outward normal of $\partial\Sigma_{t_0}$) we have

$$\begin{aligned} \int_{P_\gamma} \eta_{\Gamma_{t_0}} &= \int_{P_\gamma} d\chi \wedge H_\lambda^{t_0} dt = \int_{P_\gamma \cap \Gamma_{t_0}} d\chi \times \int_{-a}^a H_\lambda^{t_0} dt \\ &= 1 \end{aligned}$$

where P_γ is the component of $P \cap U$ which contains γ . Thus $\eta_{\Gamma_{t_0}}$ is the Thom Class of $N\Gamma_{t_0}$ (see [5] chapter I.6) and from proposition 6.24 of [5] this is the Poincaré dual of Γ_{t_0} in \mathcal{C} . In the other hand, since $supp(\eta_{\Gamma_{t_0}}) \subset U$, we have

$$\int_{\mathcal{C}} (i\psi, d\psi) \wedge \eta_{\Gamma_{t_0}} = \int_U (i\psi, d\psi) \wedge \eta_{\Gamma_{t_0}}$$

and since $(i\psi, d\psi)$ is regular in U and $d(i\psi, d\psi) = 0$ in U , the Poincaré duality implies

$$\int_{\mathcal{C}} (i\psi, d\psi) \wedge \eta_{\Gamma_{t_0}} = \int_{\Gamma_{t_0}} (i\psi, d\psi) = 2\pi deg_{\Gamma_{t_0}} \psi. \tag{A. 16}$$

Remark that $\eta_{\Gamma_{t_0}} = d(\chi H_\lambda^{t_0} dt)$ thus, since $supp(\chi H_\lambda^{t_0} dt) \subset \mathcal{C}$, applying (A. 5) we have

$$\begin{aligned} \int_{\mathcal{C}} (i\psi, d\psi) \wedge \eta_{\Gamma_{t_0}} &= - \langle *d(i\psi, d\psi); \chi H_\lambda^{t_0} dt \rangle \\ &= -2\pi \langle T; \chi H_\lambda^{t_0} dt \rangle \end{aligned} \tag{A. 17}$$

By hypothesis $supp\mathbb{T} \subset C$ and $\chi \equiv 1$ in C thus

$$\langle \mathbb{T}; \chi H_\lambda^{t_0} dt \rangle = \int_M H_\lambda^{t_0} \langle \xi; e_z \rangle \theta d\mathcal{H}^1 \tag{A. 18}$$

Applying the coarea formula (A.8) for $g = H_\lambda^{t_0} \frac{\langle \xi, e_3 \rangle}{|\langle \xi, e_3 \rangle|} \theta$ we have

$$\begin{aligned} \langle \mathbb{T}; \chi H_\lambda^{t_0} dt \rangle &= \int_{-a}^{+a} \sum_{y \in M \cap \omega_t} H_\lambda^{t_0}(t) \theta(y) \frac{\langle \xi, e_3 \rangle}{|\langle \xi, e_3 \rangle|} dt \\ &= \int_{-a}^{+a} \frac{1}{\lambda} H \left(\frac{t - t_0}{\lambda} \right) f(t) dt \end{aligned} \tag{A.19}$$

where f is the L^1 function on $(-a, a)$ such that

$$f(t) = \sum_{y \in M \cap \omega_t} \theta(y) \frac{\langle \xi; e_3 \rangle}{|\langle \xi; e_3 \rangle|} \tag{A.20}$$

For almost every $t \in (-a, a)$

$$\frac{1}{\lambda} \int_{t_0 - \lambda}^{t_0 + \lambda} |f(t) - f(t_0)| dt \longrightarrow 0 \text{ as } \lambda \rightarrow +\infty \tag{A.21}$$

Suppose the t_0 we have chosen verifying (A.7) is also a Lebesgue point for f verifying (A.21), thus we have

$$\begin{aligned} |\langle T, \chi H_\lambda^{t_0} dt \rangle - f(t_0)| &= \left| \int_{-a}^{+a} \frac{1}{\lambda} H \left(\frac{t - t_0}{\lambda} \right) (f(t) - f(t_0)) dt \right| \\ &\leq \frac{1}{\lambda} \|H\|_{L^\infty} \int_{t_0 - \lambda}^{t_0 + \lambda} |f(t) - f(t_0)| dt \end{aligned} \tag{A.22}$$

where we have used the fact that $\frac{1}{\lambda} \int_{-a}^{+a} H \left(\frac{t - t_0}{\lambda} \right) = 1$. Combining (A.16), (A.17), (A.19), (A.20) (A.21) and (A.22) we get the desired result. ■

Lemma A.6. *Let $0 < R_1 < R_2$, $0 < \eta < R_1$, N points a_1, \dots, a_N in the ball $B_{R_1}(0)$ of \mathbb{R}^2 such that*

$$\forall 1 \leq i \leq N \quad B_\eta(a_i) \subset B_{R_1}(0) \quad ,$$

$$\forall i \neq j \quad B_\eta(a_i) \cap B_\eta(a_j) = \emptyset \quad ,$$

and let d_1, \dots, d_N be N integers in $\{-D, -D + 1, \dots, D - 1, D\}$. For any $u : \bar{B}_{R_2}(0) \setminus \cup_{i=1}^N B_\eta(a_i) = \Omega_\eta \rightarrow S^1$ such that

$$\text{deg}_{\partial B_\eta(a_i)} u = d_i \quad ,$$

we have

$$\int_{\Omega_\eta} |\nabla u|^2 \geq 2\pi \left| \sum_{i=1}^N d_i \right| \log \frac{R_2 - R_1}{\eta} - 2\pi \log 2 D N^5 \tag{A.23}$$

■

Remark A.1. *The quantity $-2\pi \log 2 D N^5$ in the lower bound (A. 23) is not optimal at all. We were just looking for a lower bound $2\pi \left| \sum_{i=1}^N d_i \right| \log \frac{R_2 - R_1}{\eta}$ modulo a term which does not depend on η .*

Proof of Lemma A.6.

Let $d = \sum_{i=1}^N d_i$. suppose $d > 0$. We are going to prove by induction on n the following proposition:

Under the previous hypothesis for $n \leq d$ we have

$$\int_{\Omega_\eta} |\nabla u|^2 \geq 2\pi n \log \frac{R_2 - R_1}{\eta} - 2\pi \log 2 D N^5 \tag{A. 24}$$

the case $n = 1$.

We construct by induction on i for $1 \leq i \leq i_0 \leq N$, a sequence $\eta \leq r_i \leq R_2 - R_1$ such that

$$\left\{ \begin{array}{l} 1 - \quad r_i = \eta \text{ if } B_\eta(a_i) \subset B_{r_k}(a_k) \text{ for some } k < i \\ 2 - \quad r_i \text{ is the largest real number such that} \\ \quad r_i \leq R_2 - R_1 \quad , \\ \quad \partial B_{r_i}(a_i) \cap (\cup_{j>i} B_\eta(a_j) \cup_{k<i} B_{r_k}(a_k)) = \emptyset \quad , \\ \quad \forall \eta \leq r < r_i \text{ for which } \partial B_r(a_i) \cap (\cup_{j>i} B_\eta(a_j) \cup_{k<i} B_{r_k}(a_k)) = \emptyset \\ \quad \text{we have } \text{deg}(u, \partial B_r(a_i)) \neq 0 \\ \quad \text{and either } \text{deg}(u, \partial B_{r_i}(a_i)) = 0 \text{ or } r_i = R_2 - R_1 \end{array} \right.$$

we stop the construction of the sequence when we obtain an $i = i_0$ for which $r_{i_0} = R_2 - R_1$.

It is clear that there exists some $i_0 \in [1, N]$ such that $r_{i_0} = R_2 - R_1$. Indeed $(B_{r_i}(a_i))_{1 \leq i \leq i_0}$ is a union of disks which are either disjoint, or one nested in the other. Moreover if $B_{r_i}(a_i)$ is not contained in an other $B_{r_k}(a_k)$ we have for $i < i_0$

$$\text{deg}(u, \partial B_{r_i}(a_i)) = 0 \quad ,$$

thus $\cup_{i < i_0} B_{r_i}(a_i)$ is a union of disjoint disks such that the degree of u on each component of the total boundary is equal to zero. Since $\sum_{i=1}^N d_i = d > 0$, for topological reasons, it would not be possible to enclose all our balls $B_\eta(a_i)$ in a neutral cluster (i.e. a domain ω such that $\text{deg}(u, \partial\omega) = 0$), thus such an $i_0 \leq N$ for which $r_{i_0} = R_2 - R_1$ exists.

We introduce some notations for $1 \leq i \leq i_0$. Let I_i be the following set

$$I_i = \{ \eta \leq r < r_i \text{ s.t. } \partial B_r(a_i) \cap (\cup_{j>i} B_\eta(a_j) \cup_{k<i} B_{r_k}(a_k)) = \emptyset \}$$

and we have

$$\text{if } r_i \neq \eta \quad \forall r \in I_i \quad \text{deg}(u, \partial B_r(a_i)) \neq 0. \tag{A. 25}$$

Let E_i and F_i be the following sets

$$\text{for } i \leq i_0 \quad E_i = \{l \neq i \text{ s.t. } B_{r_l}(a_l) \subset B_{r_i}(a_i)\} \quad ,$$

and

$$\text{for } i \leq i_0 \quad F_i = \left\{ \begin{array}{l} l \neq i \text{ s.t. } B_{r_l}(a_l) \subset B_{r_i}(a_i) \text{ and} \\ \forall l' \neq l, l' < i \quad B_{r_l}(a_l) \not\subset B_{r_{l'}}(a_{l'}) \end{array} \right\} \quad .$$

We clearly have

$$E_i = \bigcup_{l \in F_i} E_l \cup \{i\} \quad , \tag{A. 26}$$

and moreover

$$\forall l \neq l' \in F_i \quad (E_l \cup \{l\}) \cap (E_{l'} \cup \{l'\}) = \emptyset \quad . \tag{A. 27}$$

Thus we have for any $i \leq i_0$

$$\begin{aligned} \int_{B_{r_i}(a_i) \cap \Omega_\eta} |\nabla u|^2 &\geq \int_{I_i} dr \int_{\partial B_{r_i}} |\nabla u|^2 + \sum_{j \in E_i} \int_{I_j} dr \int_{\partial B_r} |\nabla u|^2 \\ &= \int_{I_i} dr \int_{\partial B_r} |\nabla u|^2 + \sum_{j \in F_i} \left[\int_{I_j} dr \int_{\partial B_r} |\nabla u|^2 \right. \\ &\quad \left. + \sum_{k \in E_j} \int_{I_k} dr \int_{\partial B_r} |\nabla u|^2 \right] \end{aligned} \tag{A. 28}$$

Let us prove by induction for $i \leq i_0$ that

$$\begin{aligned} \int_{I_i} dr \int_{\partial B_r} |\nabla u|^2 + \sum_{j \in E_i} \int_{I_j} dr \int_{\partial B_r} |\nabla u|^2 &\geq 2\pi \log \frac{r_i}{\eta} - \\ &\quad - 2\pi \log 4 \#E_i. \end{aligned} \tag{A. 29}$$

Suppose (A. 29) is true for $i \leq l - 1$. Thus we have

$$|I_l| \geq |r_l - \eta| - \sum_{k \in F_l} 2r_k \quad . \tag{A. 30}$$

For $r \in I_l$ we have $\text{deg}(u, \partial B_r(a_l)) \neq 0$, thus, this implies

$$\forall r \in I_l \quad \int_{\partial B_r(a_l)} |\nabla u| \geq 2\pi \quad , \tag{A. 31}$$

and by Cauchy-Schwartz inequality we get

$$\forall r \in I_l \quad \int_{\partial B_r(a_l)} |\nabla u|^2 \geq \frac{2\pi}{r} \quad . \tag{A. 32}$$

Combining (A. 32) and (A. 30) we have

$$\begin{aligned} \int_{I_l} dr \int_{\partial B_r} |\nabla u|^2 &\geq \int_{\eta + \sum_{k \in F_l} 2r_k}^{r_l} \frac{2\pi}{r} dr \\ &= 2\pi \log \frac{r_l}{\eta} - 2\pi \log \left(1 + \sum_{k \in F_l} 2 \frac{r_k}{\eta} \right) \end{aligned} \tag{A. 33}$$

Combining (A. 28) for $i = l$, (A. 33) and the fact that (A. 29) is true for $i \leq l - 1$ we have

$$\begin{aligned} & \int_{I_l} \int_{\partial B_r} |\nabla u|^2 + \sum_{j \in E_l} \int_{I_j} \int_{\partial B_r} |\nabla u|^2 \geq \\ & \int_{I_l} \int_{\partial B_r} |\nabla u|^2 + \sum_{k \in F_l} \left[2\pi \log \frac{r_k}{\eta} - 2\pi \log 4 \#E_k \right] \geq 2\pi \log \frac{r_l}{\eta} \quad (A. 34) \\ & - 2\pi \log \left[1 + \sum_{k \in F_l} \frac{2r_k}{\eta} \right] + 2\pi \log \left[\prod_{k \in F_l} \right] \left(\frac{r_k}{\eta} \right) - 2\pi \log 4 \sum_{k \in F_l} \#E_k \quad . \end{aligned}$$

Since $r_k \geq \eta$ we clearly have that

$$1 + \sum_{k \in F_l} \frac{2r_k}{\eta} \leq 4 \#F_l \prod_{k \in F_l} \frac{r_k}{\eta} \quad .$$

Thus (A. 34) yields

$$\begin{aligned} & \int_{I_l} \int_{\partial B_r} |\nabla u|^2 + \sum_{j \in E_l} \int_{I_j} \int_{\partial B_r} |\nabla u|^2 \geq 2\pi \log \frac{r_l}{\eta} - 2\pi \log(4 \#F_l) \\ & - 2\pi \log 4 \sum_{k \in F_l} \#E_k \quad . \quad (A. 35) \end{aligned}$$

From (A. 26) and (A. 27) we deduce that

$$\#E_l = \#F_l + \sum_{k \in F_l} \#E_k \quad . \quad (A. 36)$$

Combining (A. 35) and (A. 36) we obtain that the proposition is true for $i = l$, thus (A. 29) holds for $i \leq i_0$.

Applying inequality (A. 29) for $i = i_0$, since $r_{i_0} = R_2 - R_1$, we obtain (A. 24) for $d \leq 1 = n$.

The case $n > 1$

Suppose (A. 24) is true for any positive integer $n \leq n'$. Take a configuration of $(a_i, d_i)_{1 \leq i \leq N}$ such that $\sum_{i=1}^N d_i \leq n' + 1$.

If $N = 1$ the proof of (A. 24) for $n' + 1$ is straightforward. Suppose now $N > 1$ and let

$$L = \max_{i \neq i'} |a_i - a_{i'}|.$$

we can always assume that $|a_1 - a_2| = L$. We clearly have

$$\begin{aligned} \int_{\Omega_\eta \setminus L+\eta(a_i)} |\nabla u|^2 &\geq \int_{L+\eta}^{R_2-R_1} \int_{\partial B_r} |\nabla u|^2 \\ &\geq 2\pi(n'+1)^2 \log \frac{R_2 - R_1}{L + \eta}. \tag{A.37} \\ &\geq 2\pi(n'+1) \log \frac{R_2 - R_1}{L + \eta} \end{aligned}$$

If $L \leq 2^{2(N^2+1)}\eta$, $\log \frac{R_2-R_1}{L+\eta} \geq \log \frac{R_2-R_1}{\eta} - 2(N^2 + 2) \log 2$ and since $N \geq 2$ (A.37) directly implies (A.24) for $d \geq n' + 1$. If $L \geq 2^{2(N^2+1)}\eta$, there exists necessarily $k_1 \in [1, 2N]$ such that

$$\forall i \quad B_\eta(a_i) \cap [B_{2^{-k_1}L}(a_1) \setminus B_{2^{-k_1-1}L}(a_1)] = \emptyset$$

In fact, since $\eta \leq 2^{-2N^2}L$ each $B_\eta(a_i)$ for $i = 2, \dots, N$ can intersect at most two $B_{2^{-k}L}$ for $k \in [1, 2N]$.

In the same way, by induction for any $j \leq N$ such that $a_j \notin \cup_{i \leq j-1} B_{2^{-k_i}L}(a_i)$ we construct $k_j \in [2jN; 2(j+1)N]$ verifying

$$\forall i \quad B_\eta(a_i) \cap [B_{2^{-k_j}L}(a_j) \setminus B_{2^{-k_j-1}L}(a_j)] = \emptyset$$

and we have

$$\forall i < j \quad B_{\frac{3}{2}2^{-k_i}L}(a_i) \cap B_{2^{-k_j}L}(a_j) = \emptyset$$

We are lead to introduce the following notation, $C(a, s, r)$ denotes the annulus for $s \geq 1, r \leq \eta$ and $a \in B_{R_1}(0)$

$$C(a, s, r) = B_{sr}(a) \setminus B_r(a)$$

We have constructed disjoint annuli $C(a_j, \frac{3}{2}, 2^{-k_j-1}L)$ for some j in $[1, N]$ such that $k_j \leq 2N^2$ and the balls $B_\eta(a_k)$ are included in the $B_{2^{-k_j-1}L}(a_j)$. The sum of the degrees of u on $\partial B_{2^{-k_j-1}L}(a_j)$ is equal to $d \geq n' + 1$.

Suppose there is exactly one of those $j = j_1$ such that $\text{deg}(u, \partial B_{2^{-k_{j_1}-1}L}(a_{j_1})) \neq 0$, then we must have

$$\text{deg}(u, \partial B_{2^{-k_{j_1}-1}L}(a_{j_1})) = d$$

In that case we repeat the procedure in $B_{2^{k_{j_1}-1}L}(a_{j_1})$. That is, let $R_{2,2} = 3/2 \cdot 2^{-k_{j_1}-1}L$, $R_{1,2} = 2^{-k_{j_1}-1}L$ and N_2 be the number of a_k in $B_{R_{2,2}}(a_{j_1})$.

If $N_2 = 1$.

We clearly have

$$\begin{aligned} \int_{\Omega_\eta \setminus B_\eta(a_{j_1})} |\nabla u|^2 &\geq 2\pi(n' + 1)^2 \log \frac{R_{2,2} - R_{1,2}}{\eta} \\ &\geq 2\pi(n' + 1) \log \frac{L}{\eta} - (k_{j_1} + 2) \log 2 (n' + 1) \quad (\text{A. 38}) \\ &\geq 2\pi(n' + 1) \log \frac{L}{\eta} - 2(N^2 + 1) \log 2 (n' + 1) \end{aligned}$$

and combining (A. 38) with (A. 37) we get the result.

If $N_2 > 1$.

Let L_2 be the following length

$$L_2 = \max_{i \neq i', a_i, a_{i'} \in B_{R_{2,2}}(a_{j_1})} \{|a_i - a_{i'}|\} = |a_{i_2} - a_{i'_2}|$$

and we have

$$\begin{aligned} \int_{B_{R_{2,2}}(a_{j_1}) \setminus B_{L_2+\eta}(a_{i_2})} |\nabla u|^2 &\geq \int_{L_2+\eta}^{R_{2,2}-R_{1,2}} \int_{\partial B_r} |\nabla u|^2 \\ &\geq 2\pi(n' + 1)^2 \log \frac{R_{2,2} - R_{1,2}}{L_2 + \eta} \quad (\text{A. 39}) \\ &\geq 2\pi(n' + 1) \log \frac{R_{2,2} - R_{1,2}}{L_2 + \eta} \end{aligned}$$

If $L_2 \leq 2^{2N_2+2}\eta$.

In this case we have

$$\log \frac{R_{2,2} - R_{1,2}}{L_2 + \eta} \geq \log \frac{L}{\eta} - (k_{j_1} + 4 + 2N_2^2) \log 2$$

but $(k_{j_1} + 4 + 2N_2^2) \geq 2(N^2 + N_2^2 + 2)$ and since $N \leq (N_2 + 1)$ and $N_2 \geq 2$ we easily get that $2(N^2 + N_2^2 + 2) < N^4$. Finally combining (A. 39) and (A. 37) we get the desired result.

If $L_2 > 2^{2N_2^2}\eta$

As previously we construct disjoint annuli $C(a_j, \frac{3}{2}, 2^{-k_j-1}L_2)$ for some j in $[1, N]$ such that $a_j \in B_{R_{2,2}}(a_{j_1})$ and $k_j \leq 2N_2^2$ in such a way that all the $B_\eta(a_k)$ in $B_{R_{2,2}}(a_{j_1})$ are included in the $B_{2^{-k_j-1}L_2}(a_j)$.

As before the sum of the degrees of v on the $\partial B_{2^{-k_j-1}L_2}(a_j)$ is equal to $d \geq n' + 1 \dots$

In this way we construct by induction for $l \leq p$ the following sequences together

$$\begin{aligned} N_l \in \mathbb{N}^* \quad s.t. \quad N_{l-1} \leq N_l + 1 \\ a_{j_l} \in \{a_1, \dots, a_N\} \end{aligned}$$

$$k_{j_l} \in \mathbb{N}^* \text{ also denoted } k_l \text{ s.t. } k_l \leq 2N_l^2$$

$$L_l > 0 \text{ s.t. } L_l \leq 2^{-k_{l-1}-1} L_{l-1}$$

$$\text{and s.t. for any } l' > l \quad B_{L_l}(a_{j_{l'}}) \subset B_{R_{1,l}}(a_{j_l})$$

$$\text{where } R_{1,l} = 2^{-k_{l-1}-1} L_{l-1} \quad \text{and} \quad R_{2,l} = \frac{3}{2} R_{1,l}$$

and finally such that

$$\int_{B_{R_{2,l}}(a_{j_{l-1}}) \setminus B_{L_l}(a_{j_l})} |\nabla u|^2 \geq 2\pi (n' + 1) \log \frac{R_{2,l} - R_{1,l}}{L_l + \eta} \tag{A. 40}$$

We stop the construction at $l = p$ when one of the following happens.

Let N_p be the number of a_k in $B_{R_{1,p}}(a_{j_{p-1}})$

$$N_p = 1$$

In this case we have

$$\int_{B_{R_{2,p}}(a_{j_{p-1}}) \setminus B_{\eta}(a_{j_p})} |\nabla u|^2 \geq 2\pi (n + 1)^2 \log \frac{R_{2,p} - R_{1,p}}{\eta} \tag{A. 41}$$

Summing (A. 40) for $l \leq p - 1$, since the $B_{R_{2,l}}(a_{j_{l-1}}) \setminus B_{L_l}(a_{j_l})$ are disjoint, we get

$$\begin{aligned} \int_{\Omega_\eta \setminus B_{R_{2,p}}(a_{j_{p-1}})} |\nabla u|^2 &\geq 2\pi (n' + 1) \log \prod_{l \geq p-1} \frac{R_{2,l} - R_{1,l}}{L_l + \eta} \\ &\geq 2\pi (n' + 1) \log \frac{R_2 - R_1}{L_{p-1} + \eta} - 2\pi (n' + 1) \log 2 \sum_{l=1}^{p-2} (2 + k_l) \\ &\geq 2\pi (n' + 1) \log \frac{R_2 - R_1}{L_{p-1} + \eta} - 2\pi (n' + 1) \log 2 \times 2 \sum_{l=1}^{p-2} (1 + N_l^2). \end{aligned} \tag{A. 42}$$

Adding, now (A. 41) and (A. 42) we have

$$\begin{aligned} \int_{\Omega_\eta} |\nabla u|^2 &\geq 2\pi (n' + 1) \log \frac{R_2 - R_1}{\eta} - \\ &\quad - 2\pi (n' + 1) \log 2 \times 2 \sum_{l=1}^{p-1} (1 + N_l^2) \quad . \end{aligned} \tag{A. 43}$$

Moreover $2 \sum_{l=1}^{p-1} (1 + N_l^2) \leq 2(p - 1)(1 + N^2) \leq 2(N - 1)(1 + N^2)$

Since we are dealing with $N \geq 3$ ($N = 1$ or 2 have already been considered), we have $2N^3 - 2(N^2 - N + 1) < N^4$ and we get the desired result because $n' + 1 \leq ND$.

$N_p > 1$. Let L_p be the following length

$$L_p = \max_{i \neq i' \quad a_i, a_{i'} \in B_{R_{2,p}}(a_{j_{p-1}})} \{|a_i - a_{i'}|\} = |a_{i_p} - a_{i'_p}| \quad .$$

We have (A. 40) for $l = p$

Case 1: $L_p \leq 2^{2N^2+2}\eta$

In this case we have

$$\begin{aligned} \int_{B_{R_{2,p}}(a_{j_{p-1}}) \setminus B_{L_p}(a_{j_{p-1}})} |\nabla u|^2 &\geq 2\pi(n'+1) \log 2^{-(k_{p-1}+2)} \frac{L_{p-1}}{L_{p+\eta}} \\ &\geq 2\pi(n'+1) \log \frac{L_{p-1}}{\eta} - 2\pi(n'+1) \log 2 \times 2(N_p^2 + N_{p-1}^2 + 2) \end{aligned} \tag{A.44}$$

Combining (A.44) and (A.42) we obtain

$$\begin{aligned} \int_{\Omega_\eta} |\nabla u|^2 &\geq 2\pi(n'+1) \log \frac{R_2 - R_1}{\eta} - \\ &\quad - 2\pi(n'+1) \log 2 \times 2 \sum_{l=1}^p (1 + N_l^2) \end{aligned} \tag{A.45}$$

Since $N_p > 1, p < N, 2 \sum_{l=1}^p (1 + N_l^2) \geq 2(N-1)(1 + N^2) < N^4$, we get the result because $n'+1 \geq ND$.

Case 2: $L_p > 2^{2N_p^2+2}\eta$

We construct, as before, disjoint annuli $(C(a_{p_i}, \frac{3}{2}, 2^{-k_i-1}L_p))_{1 \leq i \leq n_p}$ (where $n_p \geq N_p$ and $k_i \in [1, 2N_p^2]$) such that all the $B_\eta(a_k)$ in $B_{R_{2,p}}(a_{j_{p-1}})$ are included in the $B_{2^{-k_i-1}L_p}(a_{p_i})$. The sum of the degrees d_i of u on the circles $\partial B_{2^{-k_i-1}L_p}(a_{p_i})$ is equal to $d \geq n'+1$ and in this case there exists two disjoint $\partial B_{2^{-k_i-1}L_p}(a_{p_i})$ such that $d_i \neq 0$ (otherwise we continue the previous construction).

We can apply the hypothesis (A.24) for $n \leq n'$ on the $B_{\frac{3}{2}2^{-k_i-1}L_p}(a_{p_i})$ and we have $\forall 1 \leq i \leq n_p$

$$\begin{aligned} \int_{B_{\frac{3}{2}2^{-k_i-1}L_p}(a_{p_i}) \cap \Omega_\eta} |\nabla u|^2 &\geq 2\pi \min(|d_i|, n') \log \frac{2^{-k_i-2}L_p}{\eta} \\ &\quad - 2\pi \log 2D(N_p^i)^5 \end{aligned} \tag{A.46}$$

where N_p^i is the number of a_k in $B_{2^{-k_i-2}L_p}(a_{p_i})$. In particular we have

$$\sum_{i=1}^{n_p} N_p^i = N_p \tag{A.47}$$

Since $\sum_{i=1}^{n_p} |d_i| \geq n'+1$ and since there exists at least two i such that $|d_i| \neq 0$ we have $\sum_{i=1}^{n_p} \min(|d_i|, n') \geq n'+1$ and summing (A.46) for $i = 1, \dots, n_p$ we

get

$$\begin{aligned} \int_{B_{L_p}(a_{j_{p-1}})} |\nabla u|^2 &\geq 2\pi(n'+1) \log \frac{L_p}{\eta} - \\ &\quad - 2\pi D \log 2 \left[\sum_{i=1}^{n_p} [N_p^i(k_i+2) + (N_p^i)^5] \right] \quad (\text{A. 48}) \\ &\geq 2\pi(n'+1) \log \frac{L_p}{\eta} - 2\pi D \log 2 \left[2N_p(N_p^2+1) + \sum_{i=1}^{n_p} (N_p^i)^5 \right]. \end{aligned}$$

Combining (A. 42), the first line of (A. 44) and (A. 48) we obtain

$$\begin{aligned} \int_{\Omega_\eta} |\nabla u|^2 &\geq 2\pi(n'+1) \log \frac{R_2 - R_1}{\eta} - \\ &\quad - 2\pi D \log 2 \left[2N \sum_{l=1}^{p-1} (1 + N_l^2) + 2N_p(N_p^2 + 1) + \sum_{i=1}^{n_p} (N_p^i)^5 \right] \quad (\text{A. 49}) \end{aligned}$$

In the other hand we have the following bound

$$\begin{aligned} \text{let } n \geq 2, \quad \text{and} \quad (\alpha_i)_{i=1\dots n} \text{ s.t. } \alpha_i \leq 1 \\ \sum_{i=1}^n \alpha_i^5 \leq \left(\sum_{i=1}^n \alpha_i\right)^5 - 2\left(\sum_{i=1}^n \alpha_i\right)^3 - 2\sum_{i=1}^n \alpha_i \quad (\text{A. 50}) \end{aligned}$$

Indeed let l be such that $\alpha_l = \min_{i=1\dots n} \alpha_i$. Thus we have

$$\begin{aligned} \left(\sum_{i=1}^n \alpha_i\right)^5 &\geq \alpha_l^5 + \left(\sum_{i \neq l} \alpha_i\right)^5 + C_5^2 \alpha_l^2 \left(\sum_{i \neq l} \alpha_i\right)^3 + C_5^1 \alpha_l^4 \left(\sum_{i \neq l} \alpha_i\right) \quad (\text{A. 51}) \\ &\geq \sum_{i=1}^n \alpha_i^5 + \frac{5 \times 4 \times 3}{2} \left(\sum_{i \neq l} \alpha_i\right)^3 + 5 \left(\sum_{i \neq l} \alpha_i\right) \end{aligned}$$

and since $\alpha_l = \min_{i=1\dots n} \alpha_i$ we have $\sum_{i \neq l} \alpha_i \geq \frac{n-1}{n} \sum_{i=1}^n \alpha_i$ and since $n \geq 2$ we have $\frac{n-1}{n} \leq \frac{1}{2}$, thus (A. 51) yields

$$\left(\sum_{i=1}^n \alpha_i\right)^5 \geq \sum_{i=1}^n \alpha_i^5 + \frac{5 \times 4 \times 3}{2^4} \left(\sum_{i=1}^n \alpha_i\right)^3 + \frac{5}{2} \sum_{i=1}^n \alpha_i$$

which implies (A. 50). Combining (A. 50) and (A. 49) we obtain

$$\begin{aligned} \int_{\Omega_\eta} |\nabla u|^2 &\geq 2\pi(n'+1) \log \frac{R_2 - R_1}{\eta} - \\ &\quad - 2\pi D \log 2 \left[2N(p-1)(1 + N^2) + N_p^5 \right], \quad (\text{A. 52}) \end{aligned}$$

but $p-1 \leq N - N_p$ thus $2N(p-1)(1 + N^2) + N_p^5 \leq 2N(N - N_p)(1 + N^2) + N_p^5$ and since $N_p \geq 2$ one verifies that

$$N^5 \geq 2N(N - N_p)(1 + N^2) + N_p^5$$

and (A. 52) implies

$$\int_{\Omega_\eta} |\nabla u|^2 \geq 2\pi(n'+1) \log \frac{R_2 - R_1}{\eta} - 2\pi D \log 2 N^5. \quad (\text{A. 53})$$

This is the desired result. ■

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