

## DO SHAPE MEMORY ALLOY CABLES RESTRAIN THE VIBRATIONS OF GIRDER BRIDGES? A MATHEMATICAL POINT OF VIEW

V. RÉGNIER\* 

**Abstract.** We study the energy decay of a damped Euler-Bernoulli beam which is subject to a pointwise feedback force representing a Shape Memory Alloy (SMA) cable. The problem we consider is that of the paper by [Liu *et al. Int. J. Struct. Stab. Dyn.* **17** (2017) 1750076] but, for simplicity, our modelization does not take into account the additional stiffness term they considered. An explicit expression is given for the resolvent of the underlying operator as well as its eigenvalues and eigenfunctions. We show the exponential decay of the energy. The fastest decay rate is given by the supremum of the real part of the spectrum of the infinitesimal generator of the underlying semigroup since we prove the existence of a Riesz basis. To the question "Do Shape Memory Alloy cables restrain the vibrations of girder bridges?", the experiments in [Liu *et al. Int. J. Struct. Stab. Dyn.* **17** (2017) 1750076] answer positively. Our study does not allow to give a definite answer yet. The only presence of these cables may not be enough. Some physical parameters have to be chosen carefully.

**Mathematics Subject Classification.** 74K10, 35B40, 35Q74, 34L10, 34L15, 34L20, 35B35, 35E15, 93D23.

Received June 24, 2021. Accepted October 3, 2022.

### 1. INTRODUCTION

In this paper, we consider the same problem as in the paper by Liu *et al.* [12] which is rewritten for simplicity:

$$\frac{\partial^2 u}{\partial t^2}(x, t) + a \frac{\partial^4 u}{\partial x^4}(x, t) + b \frac{\partial u}{\partial t}(x, t) + \alpha \frac{\partial u}{\partial t}(\xi, t) \delta_\xi = 0 \quad \text{for } (x, t) \text{ in } (0; 1) \times (0, \infty) \quad (1.1)$$

$$u(0, t) = u(1, t) = \frac{\partial^2 u}{\partial x^2}(0, t) = \frac{\partial^2 u}{\partial x^2}(1, t) = 0 \quad \text{for } t \text{ in } (0, \infty) \quad (1.2)$$

$$u(x, 0) = u_0(x), \quad \frac{\partial u}{\partial t}(x, 0) = u_1(x) \quad \text{for } x \text{ in } (0; 1) \quad (1.3)$$

where  $a$ ,  $b$  and  $\alpha$  are strictly positive constants.

---

*Keywords and phrases:* Euler-Bernoulli beam, pointwise stabilization, resolvent operator, eigenvalues and their localization, eigenfunctions, Riesz basis, exponential stability, rate of decay.

Univ. Polytechnique Hauts-de-France, INSA Hauts-de-France, CERAMATHS-Laboratoire de Matériaux Céramiques et de Mathématiques, 59313 Valenciennes, France.

\* Corresponding author: [virginie.regnier@uphf.fr](mailto:virginie.regnier@uphf.fr)

In this problem the function  $u$  denotes the transverse displacement of the bridge deck represented by a beam with a uniform section,  $\delta_\xi$  is the Dirac mass concentrated in the point  $\xi \in (0; 1)$ .

Denote by  $\rho$  the mass density of the beam,  $A$  the area of the cross section of the beam,  $EI$  the second moment of area of the cross-section and  $c$  the damping of the beam. Then (1.1)–(1.3) coincides with the system in [12] with  $a = \frac{EI}{\rho A}$ ,  $b = \frac{c}{\rho A}$ ,  $\alpha = \frac{c_{SMA}}{\rho A}$  and  $k_{SMA} = 0$ .

The values  $k_{SMA}$  and  $c_{SMA}$  are respectively the equivalent stiffness and equivalent damping of a damper which represents the Shape Memory Alloy (SMA) cable at the position  $x$ . They are calculated by (26) and (27) of [12].

In fact, the problem considered in [12] is more precisely the following one with  $\beta > 0$ :

$$\frac{\partial^2 u}{\partial t^2}(x, t) + a \frac{\partial^4 u}{\partial x^4}(x, t) + b \frac{\partial u}{\partial t}(x, t) + \left( \alpha \frac{\partial u}{\partial t}(\xi, t) + \beta u(\xi, t) \right) \delta_\xi = 0 \quad \text{for } (x, t) \text{ in } (0; 1) \times (0, \infty) \quad (1.4)$$

$$u(0, t) = u(1, t) = \frac{\partial^2 u}{\partial x^2}(0, t) = \frac{\partial^2 u}{\partial x^2}(1, t) = 0 \quad \text{for } t \text{ in } (0, \infty) \quad (1.5)$$

$$u(x, 0) = u_0(x), \quad \frac{\partial u}{\partial t}(x, 0) = u_1(x) \quad \text{for } x \text{ in } (0; 1). \quad (1.6)$$

Indeed Liu *et al.* have realized experiments in which the value  $k_{SMA}$  (equivalent stiffness of the damper) does not vanish. We start here with a simpler problem but we conjecture that the result will not change since the term we keep is the most important one for the stabilization. In this paper,  $\beta = 0$ . The most important results are generalized to the case  $\beta > 0$  in remarks (see Rems. 4.5 and 4.8).

Note that the case  $a = 1$ ,  $b > 0$  and  $\alpha = 0$  is treated in the paper by Ammari *et al.* [4]. They even consider  $L^\infty(0; 1)$  functions for  $b$ , which are positive and non-negative on an open subset of  $(0; 1)$ . For a constant  $b$ , the energy is proved to decrease exponentially and the fastest decay rate is given by the supremum of the real part of the spectrum of the infinitesimal generator of the underlying semigroup.

The case  $a = 1$ ,  $b = 0$  and  $\alpha > 0$  is treated in the paper by Ammari and Tucsnak [5] with two types of boundary conditions. They study the energy decay of a Euler-Bernoulli beam which is subject to a pointwise feedback force (given by the Dirac term). They show that both uniform and non uniform energy decay may occur depending on the boundary conditions. In the case of non uniform decay in the energy space (which is the case we are interested in), they give explicit polynomial decay estimates valid for regular initial data. Their method consists of deducing the decay estimates from observability inequalities for the associated undamped problem *via* sharp trace regularity results.

The same problem is studied by the same authors in [6] one year later. It is the second example of the applications given in Section 5. The estimates are unchanged for this example but the paper gives more general results: under a regularity assumption, the authors show that observability properties for the undamped problem imply decay estimates for the damped problem.

Let  $u$  be a regular solution of system (1.1)–(1.3). Its associated total energy is defined by

$$E(t) = \frac{1}{2} \int_0^1 (|u_t(x, t)|^2 + a|u_{xx}(x, t)|^2) dx. \quad (1.7)$$

Then a classical computation using parts integration gives:

$$\frac{d}{dt} E(t) = -b \left( \int_0^1 |u_t(x, t)|^2 dx \right) - \alpha |u_t(\xi, t)|^2 \leq 0. \quad (1.8)$$

Hence system (1.1)–(1.3) is dissipative in the sense that its associated energy is non increasing with respect to time.

A lot of studies have been performed by many authors in the stabilization of Euler-Bernoulli beams. Some of them can be found in the bibliography of the already cited papers: [4–6].

The control of networks of Euler-Bernoulli or Timoshenko beams was studied in the papers by Dekoninck and Nicaise [11], Castro and Zuazua [9], Mercier and Régnier [14–16] and Akil *et al.* [2] for example. Spectral methods are used like in this paper.

More recently, Euler-Bernoulli beams are coupled with wave equations with a Kelvin-Voigt damping in the paper by Akil *et al.* [3] for example. See also the references of this paper.

The main goal of this work is to obtain the energy decay rate of the damped Euler-Bernoulli beam subject to a pointwise feedback force representing a Shape Memory Alloy (SMA) cable described by (1.1) to (1.3).

First we establish the well-posedness and strong stability of the problem. Then an explicit expression for the resolvent is given as well as the eigenvalues and eigenfunctions of the associated dissipative operator. We study the localization of the eigenvalues of the operator for small values of  $\alpha > 0$ . We prove that the eigenfunctions are quadratically close to those of the case  $\alpha = 0$ . Thus they form a Riesz basis using a result of the book by Pöschel and Trubowitz [19].

At last, we give the explicit exponential decay estimate of the energy for regular initial data. The presence of a SMA cable ( $\alpha > 0$ ) with a weak damping (small value for  $\alpha$ ) may not restrain the vibrations of girder bridges. Our modelization does not take into account the stiffness term considered in [12] ( $\beta = 0$  here). The term which contributes most to the damping is  $\alpha u_t$ . That is why we start with this situation. See Remark 5.2 which confirms this intuition. The stiffness term ( $\beta \delta_\xi u$ ) even seems to diminish the decay rate of the energy.

Note that the results were not obvious. One may think that adding a damping term ( $\alpha > 0$  versus  $\alpha = 0$ ) always leads to a better decay rate of the energy. This is not so clear. What we already know (*cf.* the conclusion of this paper) is that, if  $\alpha = 0$ , increasing the value of  $b$  does not always lead to a higher decay rate of the energy, which is rather counterintuitive.

This paper does not confirm the results of [12]. It does not refute them either. The situation is more complicated. First they have an additional stiffness term but we conjecture that the results will be analogous. Indeed the eigenfunctions are unchanged and the characteristic equation is similar (*cf.* Rems. 4.5 and 4.8). Moreover the large eigenvalues of the case  $\alpha > 0$  and  $\beta = 0$  are proved to be close to those of the case  $\alpha = 0$  and  $\beta = 0$  (see Prop. 4.6) and Rem. 5.2 proves that adding the stiffness term ( $\beta > 0$ ) leads to some eigenvalues with a real part larger than the supremum of the real part of the eigenvalues with  $\beta = 0$ .

All that is not necessarily in contradiction with the experiments of [12], since we have excluded some values for  $\xi$  (*cf.* Thm. 4.3). As it is said in the conclusion of this paper, some more work is required to give a definite answer to the question in its title.

The main result of the paper is the Theorem proved in its last section.

**Theorem 1.1.** (*Exponential stability and decay rate for small values of  $\alpha$* )

*System (1.1)–(1.3) is exponentially stable and, for any  $\epsilon > 0$ , there exists  $\alpha_0$  such that for any  $0 < \alpha < \alpha_0$ , for any solution  $U$  of (2.5) with initial data  $U_0 \in D(\mathcal{A})$  (*cf.* (2.3) and (2.4)), there exist constants  $C > 0$  and  $\omega_0(\epsilon) < 0$  depending on  $a$ ,  $b$  and  $\alpha$  such that the energy defined by (1.7) satisfies:*

$$E(t) \leq C e^{2\omega_0(\epsilon)t} \|U_0\|_{D(\mathcal{A})}^2, \forall t > 0,$$

where

$$\begin{cases} \omega_0(\epsilon) := -\frac{b}{2} + \epsilon, & \text{if } b \leq 2\sqrt{a}\pi^2 \\ \omega_0(\epsilon) := \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}) + \epsilon, & \text{if } b > 2\sqrt{a}\pi^2. \end{cases}$$

## 2. WELL-POSEDNESS AND STRONG STABILITY

In this section we study the existence, uniqueness and strong stability of the solution of system (1.1)–(1.3). The energy space  $\mathcal{H}$  is defined as follows

$$\mathcal{H} = [H^2(0; 1) \cap H_0^1(0; 1)] \times L^2(0; 1) \quad (2.1)$$

with the inner product defined by

$$(U_1, U_2)_{\mathcal{H}} = \int_0^1 (au_{1,xx}\overline{u_{2,xx}} + v_1\overline{v_2})dx, \quad (2.2)$$

for all  $U_1 = (u_1, v_1)$ ,  $U_2 = (u_2, v_2) \in \mathcal{H}$ .

Here again  $a$  is a strictly positive constant (as in the introduction).

**Remark 2.1.** The norm  $(U, U)_{\mathcal{H}}^{\frac{1}{2}}$  induced by (2.2) is equivalent to the usual norm of  $\mathcal{H}$ .

For shortness we denote by  $\|\cdot\|$  the  $L^2(\Omega)$ -norm.

Now, we define a linear unbounded operator  $\mathcal{A} : D(\mathcal{A}) \rightarrow \mathcal{H}$  by:

$$D(\mathcal{A}) = \left\{ U \in \mathcal{H} : (u, v) \in [H^4(0, \xi) \cap H^4(\xi, 1) \cap H^2(0; 1)] \times H^2(0; 1), u(0) = v(0) = 0 \right. \\ \left. u(1) = v(1) = 0, u_{xx}(0) = u_{xx}(1) = 0, u_{xx}(\xi^-) = u_{xx}(\xi^+), u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-) = -\frac{\alpha}{a}v(\xi) \right\} \quad (2.3)$$

$$\mathcal{A}(u, v) = (v, -au_x^{(4)} - bv - \alpha v(\xi)\delta_\xi), \quad \forall U = (u, v) \in D(\mathcal{A}). \quad (2.4)$$

Then we rewrite formally System (1.1)–(1.3) into the evolution equation

$$\begin{cases} U_t = \mathcal{A}U, \\ U(0) = U_0, \quad U_0 \in \mathcal{H} \end{cases} \quad (2.5)$$

with  $U = (u, u_t)$ .

**Proposition 2.2.** *The operator  $\mathcal{A}$  is  $m$ -dissipative in the energy space  $\mathcal{H}$ .*

*Proof.* On the first hand, the dissipativeness holds since we can check using integrations by parts:

$$\Re(\mathcal{A}U, U)_{\mathcal{H}} = -b \left( \int_0^1 |v|^2 dx \right) - \alpha |v(\xi)|^2 \leq 0, \forall U = (u, v) \in \mathcal{D}(\mathcal{A}). \quad (2.6)$$

On the other end, the maximality is proved in the following way.

Let  $f := (f_1; f_2) \in \mathcal{H}$ . We look for  $U := (u; v) \in \mathcal{D}(\mathcal{A})$  solution of

$$- \mathcal{A}U = f \quad (2.7)$$

or equivalently

$$\begin{cases} f_1 = -v \\ f_2 = -au_x^{(4)} - bv - \alpha v(\xi)\delta_\xi \end{cases} \Leftrightarrow \begin{cases} v = -f_1 \\ au_x^{(4)} = f_2 + bf_1 - \alpha f_1(\xi)\delta_\xi \end{cases} \quad (2.8)$$

Assume that such a solution  $u$  exists, then multiplying the second identity by a function  $\phi \in V := H^2(0; 1) \cap H_0^1(0; 1)$ , integrating in space and using integrations by parts, it follows, since  $U \in \mathcal{D}(\mathcal{A})$  and  $v = -f_1$ :

$$\begin{aligned} & \int_0^1 au_x^{(2)} \phi_x^{(2)} dx + [au_x^{(3)}(\xi^-) - au_x^{(3)}(\xi^+)]\phi(\xi) = \int_0^1 (f_2 + bf_1)\phi dx - \alpha f_1(\xi)\phi(\xi) \quad (2.9) \\ \Leftrightarrow & \int_0^1 au_x^{(2)} \phi_x^{(2)} dx + [\alpha v(\xi)]\phi(\xi) = \int_0^1 (f_2 + bf_1)\phi dx + \alpha v(\xi)\phi(\xi) \Leftrightarrow \int_0^1 au_x^{(2)} \phi_x^{(2)} dx = \int_0^1 (f_2 + bf_1)\phi dx. \end{aligned}$$

This problem has a unique solution  $u \in V = H^2(0; 1) \cap H_0^1(0; 1)$  by Lax-Milgram's lemma, because the left-hand side of the last equation of (2.9) is coercive on  $V$ .

If we consider  $\phi \in (\mathcal{D}(0, \xi) \cap \mathcal{D}(\xi, 1)) \subset V$ , then  $u$  satisfies

$$au_x^{(4)} = f_2 + bf_1 - \alpha f_1(\xi)\delta_\xi \text{ in the distributional sense.}$$

This directly implies that  $u \in H^4(0; \xi) \cap H^4(\xi; 1)$  since  $f_1$  and  $f_2$  belong to  $L^2(0; 1)$ .

Now, integrating by parts leads to:

$$\begin{aligned} & \int_0^1 (f_2 + bf_1 - \alpha f_1(\xi)\delta_\xi)\phi dx = \int_0^\xi au_x^{(4)}\phi dx + \int_\xi^1 au_x^{(4)}\phi dx = au_x^{(3)}(\xi^-)\phi(\xi^-) - au_x^{(3)}(\xi^+)\phi(\xi^+) \\ & + au_x^{(2)}(\xi^+)\phi_x(\xi^+) - au_x^{(2)}(\xi^-)\phi_x(\xi^-) + \int_0^\xi au_x^{(2)}\phi_x^{(2)} dx + \int_\xi^1 au_x^{(2)}\phi_x^{(2)} dx. \end{aligned}$$

Consequently, by taking particular test functions  $\phi$  and using  $v = -f_1$ , we obtain the last three conditions in  $\mathcal{D}(\mathcal{A})$  for  $u$ . Thus we get  $U \in \mathcal{D}(\mathcal{A})$  satisfying (2.7), which finishes the proof of maximality.  $\square$

Using Lumer-Phillips Theorem (see Pazy [18], Thm. 1.4.3), the operator  $\mathcal{A}$  generates a  $C_0$ -semigroup of contractions  $e^{t\mathcal{A}}$  on  $\mathcal{H}$ . Then, we have the following result.

**Theorem 2.3.** *(Existence and uniqueness)*

(1) If  $U_0 \in \mathcal{D}(\mathcal{A})$ , then system (2.5) has a unique strong solution

$$U \in C^0(\mathbb{R}_+, \mathcal{D}(\mathcal{A})) \cap C^1(\mathbb{R}_+, \mathcal{H}).$$

(2) If  $U_0 \in \mathcal{H}$ , then system (2.5) has a unique weak solution

$$U \in C^0(\mathbb{R}_+, \mathcal{H}).$$

Now, we have the following general strong stability result.

**Theorem 2.4.** *(Strong stability)*

System (1.1)–(1.3) is strongly stable, i.e. for any solution  $U$  of (2.5) with initial data  $U_0 \in \mathcal{H}$ , it holds

$$\lim_{t \rightarrow \infty} E(t) = 0,$$

where  $E(t)$  is defined by (1.7).

*Proof.* Since  $\mathcal{A}$  generates a contraction semigroup and its resolvent is compact in  $\mathcal{H}$ , using Arendt-Batty Theorem (see the book by W. Arendt and C.J.K. Batty, [7], p. 837), system (1.1)–(1.3) is strongly stable if and only if  $\mathcal{A}$  does not have purely imaginary eigenvalues.

Assume that  $\mathcal{A}$  has a purely imaginary eigenvalue denoted by  $i\mu$  with  $\mu \in \mathbb{R}$ . Then there exists  $U = (u, v) \neq (0; 0) \in \mathcal{D}(\mathcal{A})$  such that  $v = i\mu u$  and  $\mathcal{A}U = i\mu U$ . Using (2.6) leads to

$$\Re [i\mu(U, U)_{\mathcal{H}}] = 0 = -b \left( \int_0^1 |v(x)|^2 dx \right) - \alpha |v(\xi)|^2. \quad (2.10)$$

This implies  $|v(\xi)| = 0$  and  $\int_0^1 |v(x)|^2 dx = 0$  (since  $\alpha > 0$  and  $b > 0$ ). Thus  $v \equiv 0$  almost everywhere on  $(0; 1)$  and since  $v = i\mu u$ ,  $u \equiv 0$  almost everywhere on  $(0; 1)$ . Now  $u$  and  $v$  are continuous on  $(0; 1)$  and they vanish at 0 and 1 due to their belonging to  $D(\mathcal{A})$ . Thus  $u \equiv v \equiv 0$ . This contradicts the fact that  $U \neq (0; 0)$ .  $\square$

### 3. EXPLICIT EXPRESSION FOR THE RESOLVENT

In this section we give an explicit expression of the resolvent  $(\mu I - \mathcal{A})^{-1}$  and prove some useful estimates. In fact such estimates are useful since later on, we will use a result of Borichev and Tomilov ([8], Thm. 2.4) which involves the norm operator of  $(\mu I - \mathcal{A})^{-1}$  with  $\mu \in \mathbb{C}$ .

Let  $F = (u_1, v_1) \in \mathcal{H}$ , we look for a solution  $U = (u, v) \in \mathcal{D}(\mathcal{A})$  of

$$(\mu I - \mathcal{A})U = F, \quad \mu \in \mathbb{C}. \quad (3.1)$$

The explicit expression for the resolvent we give in next Proposition 3.2 involves the restriction on  $[0; 1]$  of the classical convolution product of two functions on  $\mathbb{R}$ . Let us recall the definition and two useful properties we established with D. Mercier in [17].

**Lemma 3.1.** *(A technical lemma)*

Let  $\psi \in C^\infty([0, \infty[)$  and  $f$  in  $L^2(0; 1)$  be two functions and define their convolution product  $\psi \star f$  on  $[0; 1]$  by :

$$(\psi \star f)(x) = \int_0^x \psi(x-s)f(s)ds, \quad \forall x \in [0; 1]. \quad (3.2)$$

Then the following two properties hold:

1.  $(\psi \star f) \in H^1(0; 1)$  and its derivative is:

$$(\psi \star f)'(x) = \int_0^x \psi'(x-s)f(s)ds + \psi(0)f(x), \quad \forall x \in [0; 1]. \quad (3.3)$$

2. If  $\psi(0) = 0$  is also assumed, then  $(\psi \star f) \in H^2(0; 1)$  and its second derivative is:

$$(\psi \star f)''(x) = \int_0^x \psi''(x-s)f(s)ds + \psi'(0)f(x), \quad \forall x \in [0; 1]. \quad (3.4)$$

*Proof.* 1. The functions  $\psi$  extended by 0 on  $(-\infty, 0)$  and  $f$  extended by 0 on  $\mathbb{R}$  outside  $[0; 1]$  are still called  $\psi$  and  $f$  respectively. Then the convolution product defined by (3.2) is extended by the classical convolution product on  $\mathbb{R}$  i.e. by

$$(\psi \star f)(x) = \int_{\mathbb{R}} \psi(x-s)f(s)ds, \quad \forall x \in \mathbb{R}. \quad (3.5)$$

It is well known that  $(\psi \star f)' = (\psi')_{\text{dist}} \star f$  where  $(\psi')_{\text{dist}}$  is the derivative of  $\psi$  in the distributional sense. Due to the property of  $\psi$  and its extension on  $\mathbb{R}$  we have

$$(\psi')_{\text{dist}} = \psi' + \psi(0)\delta_0,$$

where  $\delta_0$  is the Dirac distribution at  $x = 0$ . The property (3.3) follows from this remark.

2. (3.4) is a consequence of (3.3). □

Note that weaker assumptions could be made on  $\psi$  for this lemma ( $\psi \in C^2([0; 1])$  is sufficient).

**Notation for the following Proposition.** The Heaviside step function is denoted by  $H(x, 0)$  and defined by

$$H(x, 0) := \begin{cases} 0, \forall x < 0 \\ 1, \forall x \geq 0. \end{cases} \quad (3.6)$$

Let  $a, b$  and  $\alpha$  be strictly positive real numbers,  $\xi$  a real number in  $(0; 1)$ ,  $\mu$  and  $\lambda$  complex numbers. Let the spaces  $\mathcal{H}$  and  $D(\mathcal{A})$  be defined by (2.1) and (2.3). Let  $F = (u_1, v_1) \in \mathcal{H}$ .

Introduce the expressions:

$$f_1^0 := (\mu + b)u_1 + v_1 \quad (3.7)$$

$$Det^\alpha(\lambda, \xi) := 4\lambda^2 \left\{ \begin{aligned} &[-2\lambda^3 \sinh(\lambda) - \frac{\alpha\mu}{a} \sinh(\lambda\xi) \sinh(\lambda(\xi - 1))] \sin(\lambda) \\ &+ \frac{\alpha\mu}{a} \sin(\lambda\xi) \sinh(\lambda) \sin(\lambda(\xi - 1)) \end{aligned} \right\} \quad (3.8)$$

$$Det^0(\lambda) := -8\lambda^5 \sinh(\lambda) \sin(\lambda) \quad (3.9)$$

$$u_0(\lambda, x) := \frac{1}{2a\lambda^3} [\sin(\lambda x) - \sinh(\lambda x)] H(x, 0), \forall x \in (0; 1) \quad (3.10)$$

$$\begin{cases} A(\lambda, \alpha) := \lambda^2(u_0(\lambda, \cdot) \star f_1^0)(1) + 2\lambda^2\alpha u_1(\xi)u_0(\lambda, 1 - \xi) \\ B(\lambda, \alpha) := (u_0(\lambda, \cdot) \star f_1^0)_{xx}(1) + 2\alpha u_1(\xi)u_{0,xx}(\lambda, 1 - \xi) \\ C(\lambda, \xi) := \lambda^2(u_0(\lambda, \cdot) \star f_1^0)(\xi). \end{cases} \quad (3.11)$$

For simplicity, the dependency of  $A, B$  and  $C$  on  $\lambda, \alpha$  and  $\xi$  is omitted in the following.

$$\Delta(\lambda, \xi) := \frac{4\lambda^2\mu}{a} (\sin(\lambda\xi) \sin(\lambda(\xi - 1)) \sinh(\lambda) - \sinh(\lambda\xi) \sinh(\lambda(\xi - 1)) \sin(\lambda)) \quad (3.12)$$

$$\Delta_3(\lambda, \xi, \alpha) := \frac{4\lambda\mu}{a} \left\{ \begin{aligned} &\sin(\lambda(\xi - 1)) \left( \frac{A+B}{2} [\sinh(\lambda\xi) - \sin(\lambda\xi)] - C \sinh(\lambda) \right) \\ &+ \sinh(\lambda(\xi - 1)) \left( \frac{A-B}{2} [\sinh(\lambda\xi) - \sin(\lambda\xi)] + C \sin(\lambda) \right) \end{aligned} \right\} \quad (3.13)$$

$$\Delta_1(\lambda, \xi, \alpha) := \frac{4\lambda^3\mu}{a} \left\{ \begin{aligned} &\sin(\lambda(\xi - 1)) \left( -\frac{A+B}{2} [\sinh(\lambda\xi) + \sin(\lambda\xi)] + C \sinh(\lambda) \right) \\ &+ \sinh(\lambda(\xi - 1)) \left( -\frac{A-B}{2} [\sinh(\lambda\xi) + \sin(\lambda\xi)] + C \sin(\lambda) \right) \end{aligned} \right\} \quad (3.14)$$

$$\Delta_0(\lambda, \xi, \alpha) := -4\lambda^3 \left\{ 2C \sin(\lambda) \sinh(\lambda) - (A+B) \sin(\lambda) \sinh(\lambda\xi) - (A-B) \sin(\lambda\xi) \sinh(\lambda) \right\}. \quad (3.15)$$

**Proposition 3.2.** (*Explicit expression for the resolvent of the operator  $\mathcal{A}$* )

Let  $a, b$  and  $\alpha$  be strictly positive real numbers,  $\xi$  a real number in  $(0; 1)$  and  $\mu$  a complex number.

Let the spaces  $\mathcal{H}$  and  $D(\mathcal{A})$  be defined by (2.1) and (2.3). Let  $F = (u_1, v_1) \in \mathcal{H}$ .

Denote by  $\lambda$  the complex number, such that  $\exists k \in \mathbb{Z}, \arg(\lambda) \in [-\pi/4 + 2k\pi; \pi/4 + 2k\pi)$ , satisfying:

$$\lambda^4 = -\frac{b\mu + \mu^2}{a}. \quad (3.16)$$

The notation introduced just before this Proposition is adopted. For simplicity, the dependency of  $A, B$  and  $C$  on  $\lambda, \alpha$  and  $\xi$  is omitted in the following.

The operator  $\mathcal{A}_0$  is the operator  $\mathcal{A}$  defined by (2.3) and (2.4) with  $\alpha = 0$ .

$$R(\mu, \mathcal{A}_0)F := (u_0 \star f_1^0)(x) + \frac{a}{\text{Det}^0(\lambda)} \left[ \begin{aligned} &4\lambda^4 \{ (A-B) \sinh(\lambda) + (A+B) \sin(\lambda) \} u_0(x, \lambda) \\ &+ 4\lambda^6 \{ (A+B) \sin(\lambda) - (A-B) \sinh(\lambda) \} u_{0,xx}(x, \lambda) \end{aligned} \right]. \quad (3.17)$$

Then the solution  $U = (u, v) \in \mathcal{D}(\mathcal{A})$  of  $(\mu I - \mathcal{A})U = F$ , with  $F = (u_1, v_1) \in \mathcal{H}$ , is given by:

$$u(x) = \frac{1}{\text{Det}^\alpha(\lambda, \xi)} \left( \text{Det}^0(\lambda) \cdot R(\mu, \mathcal{A}_0)F + \alpha \left[ \begin{aligned} &\Delta(\lambda, \xi)(u_0 \star f_1^0)(x) + a\Delta_3(\lambda, \xi, \alpha)u_0(\lambda, x) \\ &+ a\Delta_1(\lambda, \xi, \alpha)u_{0,xx}(\lambda, x) - \mu\Delta_0(\lambda, \xi, \alpha)u_0(\lambda, x - \xi) \end{aligned} \right] \right). \quad (3.18)$$

*Proof.* The computation of an explicit expression for the resolvent of  $\mathcal{A}$  requires the search for a solution  $u$  in  $H^4(0; \xi) \cap H^2(0; 1)$  (resp. in  $H^4(\xi; 1) \cap H^2(0; 1)$ ) of the following problem on  $(0; \xi)$  (resp. on  $(\xi; 1)$ ):

$$\left\{ \begin{aligned} &\alpha u_x^{(4)} + (b\mu + \mu^2)u + \mu\alpha\delta_\xi u = f_1^\alpha \\ &v = \mu u - u_1 \\ &u(0) = u(1) = 0 \\ &u_{xx}(0) = u_{xx}(1) = 0 \\ &u_{xx}(\xi^-) = u_{xx}(\xi^+) \\ &u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-) = g_1 \end{aligned} \right. \quad (3.19)$$

with  $f_1^\alpha := (\mu + b)u_1 + v_1 + \alpha\delta_\xi u_1$  and  $g_1 := \frac{\alpha}{a}(-\mu u(\xi) + u_1(\xi))$ . Note that  $f_1^0$  defined by (3.7) is  $f_1^\alpha$  with  $\alpha = 0$ . First we still denote by  $u$  the function defined by

$$u(x) := \begin{cases} 0, & \forall x < 0 \text{ and } x > \xi \\ u(x), & \forall x \in (0; \xi) \end{cases} \quad (3.20)$$



where  $u$  is a solution of (3.19) on  $(0; \xi)$ .

Denote by  $\mathcal{L}$  the classical Laplace transform. Using four successive integrations by parts and the properties of  $u$  at  $0, 1, \xi^-$  and  $\xi^+$  leads to:

$$[\mathcal{L}u_x^{(4)}](p) = p^4[\mathcal{L}u](p) + [u_x^{(3)}(\xi^-) + pu_x^{(2)}(\xi^-) + p^2u_x(\xi^-) + p^3u(\xi^-)]e^{-p\xi} - [u_x^{(3)}(0) + p^2u_x(0)]. \quad (3.21)$$

Applying the Laplace transform to Problem (3.19) on  $(0; \xi)$  gives:

$$[ap^4 + b\mu + \mu^2][\mathcal{L}u](p) + \alpha\mu u(\xi)e^{-p\xi} + a[u_x^{(3)}(\xi^-) + pu_x^{(2)}(\xi^-) + p^2u_x(\xi^-) + p^3u(\xi^-)]e^{-p\xi} - a[u_x^{(3)}(0) + p^2u_x(0)] = \mathcal{L}(f_1^\alpha \cdot \chi_{(0;\xi)}). \quad (3.22)$$

We proceed similarly on  $(\xi; 1)$  and sum both functions to get:

$$\begin{aligned} [\mathcal{L}u](p) &= \frac{1}{ap^4 + b\mu + \mu^2} \left[ [\mathcal{L}f_1^\alpha](p) + au_x^{(3)}(0) + (u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-))ae^{-p\xi} - au_x^{(3)}(1)e^{-p} \right. \\ &\quad \left. + ap^2[u_x(0) - u_x(1)e^{-p}] - \alpha\mu u_1(\xi)e^{-p\xi} \right] \\ &= \frac{1}{ap^4 + b\mu + \mu^2} \left[ [\mathcal{L}f_1^\alpha](p) + au_x^{(3)}(0) + \alpha u(\xi)e^{-p\xi} - au_x^{(3)}(1)e^{-p} + ap^2[u_x(0) - u_x(1)e^{-p}] \right]. \end{aligned} \quad (3.23)$$

The function  $u_0$  defined by (3.10) satisfies:  $\left[ \mathcal{L} \left[ \frac{1}{2a\lambda^3} [\sin(\lambda \cdot) - \sinh(\lambda \cdot)] \right] \right] (p) = \frac{1}{ap^4 + b\mu + \mu^2}$ . Thus

$$u(x) = [u_0(\lambda, \cdot) \star f_1^\alpha](x) + au_x^{(3)}(0) \cdot u_0(\lambda, x) + au_x(0) \cdot u_{0,xx}(\lambda, x) + [\alpha u_1(\xi) - 2\alpha\mu u(\xi)] \cdot u_0(\lambda, x - \xi). \quad (3.24)$$

An analogous problem is treated in Section 2.1 of the paper by Majkut for example [13]. Taking into account  $u(1) = u_{xx}(1) = 0$  gives the following two equations:

$$\begin{cases} \lambda^2[\sinh(\lambda) + \sin(\lambda)]u_x^{(3)}(0) + [\sinh(\lambda) - \sin(\lambda)]u_x(0) - \frac{\alpha\mu}{a}[\sinh(\lambda(1 - \xi)) - \sin(\lambda(1 - \xi))]u(\xi) \\ = -2\lambda^3(u_0 \star f_1^0)(1) - 4\lambda^3\alpha u_1(\xi)u_0(1 - \xi) \\ \lambda^2[\sinh(\lambda) - \sin(\lambda)]u_x^{(3)}(0) + [\sinh(\lambda) + \sin(\lambda)]u_x(0) - \frac{\alpha\mu}{a}[\sinh(\lambda(1 - \xi)) + \sin(\lambda(1 - \xi))]u(\xi) \\ = -2\lambda(u_0 \star f_1^0)^{(2)}(1) - 4\lambda\alpha u_1(\xi)u_{0,xx}(1 - \xi). \end{cases} \quad (3.25)$$

(3.4) has been used here since  $u_0(0) = u_{0,x}(0) = 0$ . The last step is to evaluate (3.24) at  $x = \xi$ :

$$[\sinh(\lambda\xi) - \sin(\lambda\xi)]u_x^{(3)}(0) + \lambda^2[\sinh(\lambda\xi) + \sin(\lambda\xi)]u_x(0) - 2\lambda^3u(\xi) = -2\lambda^3(u_0 \star f_1^0)(\xi). \quad (3.26)$$

Now (3.25) and (3.26) form a  $3 \times 3$  system where the unknowns are  $u_x^{(3)}(0)$ ,  $u_x(0)$  and  $u(\xi)$ . Its determinant is  $Det^\alpha(\lambda)$ . Solving this system gives (3.18) after long calculations.  $\square$

**Remark 3.3.** The expression  $R(\lambda, \mathcal{A}_0)$  represents the resolvent of the operator  $\mathcal{A}_0$  which is  $\mathcal{A}$  with  $\alpha = 0$ . Thus the expression for the solution  $u$  given by (3.18) is  $R(\lambda, \mathcal{A}_0)F$  if  $\alpha = 0$ .

#### 4. EIGENVALUES AND EIGENFUNCTIONS OF THE OPERATOR $\mathcal{A}$

The eigenvalues and eigenfunctions of the operator  $\mathcal{A}$  defined by (2.3) and (2.4) in Section 2 are explicitly calculated. For the sake of completeness, the case  $\alpha = 0$  which corresponds to the operator  $\mathcal{A}_{a_0}$  of the paper by Ammari *et al.* [4]) is recalled here. Note the dependency of  $\mathcal{A}$  on  $a, b, \alpha$  and  $\xi$ .

First a technical Lemma will be useful in the proof of Theorem 4.3 to handle with particular values for  $\xi$ .

**Lemma 4.1.** *(Is  $\mu = -b$  an eigenvalue of  $\mathcal{A}$ ?)*

Assume that  $a, b, \alpha$  are strictly positive constants,  $\xi \in (0; 1)$ . The following problem on  $H^4(0; \xi) \cap H^2(0; 1)$  (resp. on  $H^4(\xi; 1) \cap H^2(0; 1)$ ):

$$\begin{cases} au_x^{(4)} - b\alpha\delta_\xi u = 0 \\ u(0) = u(1) = 0 \\ u_{xx}(0) = u_{xx}(1) = 0 \\ u_{xx}(\xi^-) = u_{xx}(\xi^+) \\ u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-) = \frac{\alpha}{a}bu(\xi) \end{cases} \quad (4.1)$$

has no other solution than  $u \equiv 0$  if the polynomial  $P(x) = 1 - \frac{2\alpha b}{3}x^2(1-x)^2$  has no solution in  $(0; 1)$ . There are three cases:

- Case  $\alpha b > 6$ :  $P$  has exactly two roots in  $(0; 1)$  called  $\xi_1$  and  $\xi_2$  in the following.
- Case  $\alpha b = 6$ :  $P$  has exactly one root in  $(0; 1)$  called  $\xi_0$  in the following.
- Case  $\alpha b < 6$ : the polynomial  $P$  does not vanish on  $(0; 1)$ .

*Proof.* The proof starts like that of Proposition 3.2.

Denote by  $u_0$  the function defined by  $u_0(x) = \frac{x^3}{6a} \cdot H(x, 0), \forall x \in (0; 1)$ .

It holds  $[\mathcal{L} \left( \frac{x^3}{6a} \right)](x \mapsto p) = \frac{1}{ap^4}, \forall p \in (0; +\infty)$ . Thus

$$u(x) = au_x^{(3)}(0) \cdot u_0(\lambda, x) + au_x(0) \cdot u_{0,xx}(\lambda, x) + \frac{\alpha}{a}bu(\xi) \cdot u_0(\lambda, x - \xi). \quad (4.2)$$

Taking into account  $u(1) = u_{xx}(1) = 0$  gives the following two equations:

$$\begin{cases} \frac{1}{6}u_x^{(3)}(0) + u_x(0) + \frac{\alpha}{a}bu(\xi)(1-\xi)^3 = 0 \\ u_x^{(3)}(0) + 2\alpha bu(\xi)(1-\xi) = 0. \end{cases} \quad (4.3)$$

The last step is to evaluate (4.2) at  $x = \xi$ :

$$\frac{\xi^3}{6}u_x^{(3)}(0) + \xi u_x(0) - u(\xi) = 0. \quad (4.4)$$

Now (4.3) and (4.4) form a  $3 \times 3$  system where the unknowns are  $u_x^{(3)}(0)$ ,  $u_x(0)$  and  $u(\xi)$ . An obvious solution is  $(0; 0; 0)$ . Its determinant is  $P(\xi)$ .

The system has no other solution than  $(0; 0; 0)$  if and only if  $P(\xi)$  does not vanish.

The roots of  $P$  follow from a classical study of the variations of the function  $P$ . □

As announced before, we recall a result already proved by Ammari *et al.* [4] written with our notation.

**Theorem 4.2.** (*Eigenvalues and eigenfunctions of the operator  $\mathcal{A}$  - Case  $\alpha = 0$* )

Let  $\mathcal{A}$  be the operator defined by (2.3) and (2.4) in Section 2. Assume that  $\alpha = 0$ . Denote by  $\mu_n$ ,  $n \in \mathbb{Z}$  the eigenvalues of  $\mathcal{A}$ .

1. If  $b \in \mathbb{R}^{+*} - \{2\sqrt{an^2\pi^2}, n \in \mathbb{Z}^*\}$ , then  $\exists n_0$  such that  $2\sqrt{an_0^2\pi^2} < b < 2\sqrt{a}(n_0 + 1)^2\pi^2$  and

$$\mu_n^\pm = \begin{cases} \frac{1}{2} \left( -b \pm \sqrt{b^2 - 4an^4\pi^4} \right), n = 1, 2, \dots, n_0 \\ \frac{1}{2} \left( -b \pm i\sqrt{4an^4\pi^4 - b^2} \right), n > n_0. \end{cases} \quad (4.5)$$

The algebraic multiplicity of these eigenvalues is 1.

2. If  $\exists n_0$  such that  $b = 2\sqrt{an_0^2\pi^2}$  then  $\mu_{n_0} = -\frac{b}{2}$  and the algebraic multiplicity of this eigenvalue may not exceed 2.

In both cases, the associated eigenfunction is defined on  $(0; 1)$  by

$$\Phi_n^\pm(x) = \sin(n\pi x)(1, \mu_n^\pm), n \geq 1. \quad (4.6)$$

*Proof.* The proof is that of Lemma 2.2 of [4], mutatis mutandis.  $\square$

The general case where  $\alpha$  does not vanish is new. The result is stated in the following.

**Theorem 4.3.** (*Eigenvalues and eigenfunctions of the operator  $\mathcal{A}$  - Case  $\alpha \neq 0$* )

Let  $\mathcal{A}$  be the operator defined by (2.3) and (2.4) in Section 2. Denote by  $\mu_n$ ,  $n \in \mathbb{Z}$  the eigenvalues of  $\mathcal{A}$  and by  $\lambda$  the complex numbers satisfying  $\lambda^4 = -\frac{b\mu + \mu^2}{a}$ . One of these values of  $\lambda$  is such that  $\exists k \in \mathbb{Z}$ ,  $\arg(\lambda) \in [-\pi/4 + 2k\pi; \pi/4 + 2k\pi)$ . The others are  $-\lambda$  and  $\pm i\lambda$  and they are still denoted by  $\lambda$ .

1. Case  $\alpha b < 6$ :  $\mu = -b$  is not an eigenvalue of  $\mathcal{A}$ .

If  $\xi \notin \mathbb{Q}$ ,  $\mu$  is an eigenvalue of  $\mathcal{A}$  if and only if the corresponding values of  $\lambda$  satisfy  $\lambda \notin \{k\pi, k \in \mathbb{Z}\}$ ,  $\lambda \notin \{ik\pi, k \in \mathbb{Z}\}$  and the characteristic equation:

$$2(\mu + b) \sinh(\lambda) \sin(\lambda) + \alpha \lambda [\sin(\lambda) \sinh(\lambda\xi) \sinh(\lambda(1 - \xi)) - \sinh(\lambda) \sin(\lambda\xi) \sin(\lambda(1 - \xi))] = 0 \quad (4.7)$$

and the associated eigenfunction is  $\Phi(\mu, \cdot) = (\phi(\lambda, \cdot))(1, \mu)$  where  $\phi(\lambda, \cdot)$  is defined on  $(0; 1)$  by

$$\begin{aligned} \phi(\lambda, x) := & \frac{1}{|\lambda|^2} \cdot e^{(\xi-2)|\lambda|} \left\{ \sin(\lambda) \sinh(\lambda) [\sin(\lambda(x - \xi)) - \sinh(\lambda(x - \xi))] \cdot H(x, \xi) \right. \\ & \left. + \sinh(\lambda) \sin(\lambda(1 - \xi)) \sin(\lambda x) - \sin(\lambda) \sinh(\lambda(1 - \xi)) \sinh(\lambda x) \right\} \end{aligned} \quad (4.8)$$

with  $H(x, \xi)$  the Heaviside step function at  $\xi$  (cf. (4.12)).

Denote by  $r = r(a, b)$  the real positive number such that  $r^4 = \frac{b^2}{4a}$ . Then, if

$$\alpha = -\frac{4ar^3 \sinh(r) \sin(r)}{b[\sinh(r) \sin(r\xi) \sin(r(\xi - 1)) - \sin(r) \sinh(r\xi) \sinh(r(\xi - 1))]} \quad (4.9)$$

then  $\mu = -b/2$  is an eigenvalue of  $\mathcal{A}$  with algebraic multiplicity at least 2 and its geometric multiplicity is 1.

If  $\alpha$  takes any other value, the algebraic multiplicity of  $\mu$  as an eigenvalue of  $\mathcal{A}$  is at least 1 and its geometric multiplicity is 1.

2. Case  $\alpha b = 6$ : there exists  $\xi_0$  in  $(0; 1)$  such that, if  $\xi = \xi_0$ ,  $\mu = -b$  is an eigenvalue of  $\mathcal{A}$  ( $= \mathcal{A}(\xi)$ ). If  $\xi \neq \xi_0$ ,  $\mu = -b$  is not an eigenvalue of  $\mathcal{A}$  ( $= \mathcal{A}(\xi)$ ). In both cases, if  $\xi \notin \mathbb{Q}$ , the other eigenvalues and the associated eigenfunctions are given by the same expressions as in the preceding case.
3. Case  $\alpha b > 6$ : there exist  $\xi_1$  and  $\xi_2$  in  $(0; 1)$  such that, if  $\xi = \xi_1$  or  $\xi = \xi_2$ ,  $\mu = -b$  is an eigenvalue of  $\mathcal{A}$  ( $= \mathcal{A}(\xi)$ ). If  $\xi \notin \{\xi_1, \xi_2\}$ ,  $\mu = -b$  is not an eigenvalue of  $\mathcal{A}$  ( $= \mathcal{A}(\xi)$ ). In both cases, if  $\xi \notin \mathbb{Q}$ , the other eigenvalues and the associated eigenfunctions are given by the same expressions as in the first case.

*Proof.* First of all, the operator  $\mathcal{A}$  has already been proved to have no imaginary eigenvalues (cf. the proof of Thm. 2.4). In particular, 0 is not an eigenvalue.

The computation of the eigenelements of  $\mathcal{A}$  requires the search for a solution  $u$  in  $H^4(0; \xi) \cap H^2(0; 1)$  (resp. in  $H^4(\xi; 1) \cap H^2(0; 1)$ ) of the following problem on  $(0; \xi)$  (resp. on  $(\xi; 1)$ ):

$$\begin{cases} au_x^{(4)} + (b\mu + \mu^2)u + \mu\alpha\delta_\xi u = 0 \\ u(0) = u(1) = 0 \\ u_{xx}(0) = u_{xx}(1) = 0 \\ u_{xx}(\xi^-) = u_{xx}(\xi^+) \\ u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-) = -\frac{\alpha}{a}\mu u(\xi) \end{cases} \quad (4.10)$$

**Case  $\alpha b < 6$ :** since neither  $\mu = 0$  nor  $\mu = -b$  is an eigenvalue of  $\mathcal{A}$  (cf. Lem. 4.1),  $\lambda$  cannot vanish and a particular solution of Problem (4.10) is:

$$u_p(\lambda, x) := -\frac{\alpha\mu}{2a\lambda^3}u(\xi) [\sinh(\lambda(x - \xi)) - \sin(\lambda(x - \xi))] H(x, \xi), \forall x \in (0; 1) \quad (4.11)$$

where  $H(x, \xi)$  is for the Heaviside step function defined by

$$H(x, \xi) := \begin{cases} 0, \forall x < \xi \\ 1, \forall x \geq \xi. \end{cases} \quad (4.12)$$

This result is not new: it can be found in [13] for example.

The general solution of this problem (i.e. with  $u_x^{(3)}(\xi^+) - u_x^{(3)}(\xi^-) = 0$ ) can be written as:

$$u_g(\lambda, x) := \begin{cases} P_1 \cosh(\lambda x) + Q_1 \sinh(\lambda x) + R_1 \cos(\lambda x) + S_1 \sin(\lambda x), \forall x \in (0; \xi] \\ P_2 \cosh(\lambda(x - \xi)) + Q_2 \sinh(\lambda(x - \xi)) + R_2 \cos(\lambda(x - \xi)) + S_2 \sin(\lambda(x - \xi)), \forall x \in [\xi; 1). \end{cases} \quad (4.13)$$

Using the boundary conditions for  $u$  and  $u_{xx}$  as well as the continuity of  $u_{xx}$  and  $u_x^{(3)}$  at  $\xi$ , it holds:

$$\begin{cases} P_1 = R_1 = 0 \\ P_2 = Q_1 \sinh(\lambda\xi), Q_2 = Q_1 \cosh(\lambda\xi), R_2 = S_1 \sin(\lambda\xi) \text{ and } S_2 = S_1 \cos(\lambda\xi) \\ \sinh(\lambda)Q_1 = -\frac{\alpha\mu}{2a\lambda^3}u(\xi) \sinh(\lambda(\xi - 1)) \text{ and } \sin(\lambda)S_1 = \frac{\alpha\mu}{2a\lambda^3}u(\xi) \sin(\lambda(\xi - 1)). \end{cases} \quad (4.14)$$

If  $\xi \notin \mathbb{Q}$ ,  $\lambda \notin \{k\pi, k \in \mathbb{Z}^*\}$  and  $\lambda \notin \{ik\pi, k \in \mathbb{Z}^*\}$ . Indeed, if  $\exists k \in \mathbb{Z}^*$ , such that  $\lambda = k\pi$ , then  $\sin(\lambda) = 0$  but  $\sinh(\lambda) \neq 0$  and the last identity of (4.14) is:

$$\frac{\alpha\mu}{2a\lambda^3}u(\xi) \sin(k\pi(\xi - 1)) = 0. \quad (4.15)$$

Since  $\alpha \neq 0$ ,  $\mu \neq 0$  and  $\sin(k\pi(\xi - 1)) \neq 0$  (the expression only vanishes if  $\xi \in \mathbb{Q}$ ), it holds  $u(\xi) = 0$ . Now, since  $\sinh(\lambda) \neq 0$ , it implies  $Q_1 = 0$  and consequently,  $P_2 = Q_2 = 0$ . The vanishing of  $u(\xi)$  also implies that of  $u_p$ . Thus

$$u(\lambda, x) = u_g(\lambda, x) := \begin{cases} S_1 \sin(n\pi x), \forall x \in (0; \xi] \\ S_1 \sin(n\pi\xi) \cos(n\pi(x - \xi)) + S_1 \cos(n\pi\xi) \sin(n\pi(x - \xi)) = S_1 \sin(n\pi x), \forall x \in [\xi; 1). \end{cases} \quad (4.16)$$

Now,  $u(\xi) = 0 = S_1 \sin(n\pi\xi)$  and  $S_1 \neq 0$  (otherwise  $u \equiv 0$ ), then  $\sin(n\pi\xi) = 0$ , which contradicts the fact that  $\xi \notin \mathbb{Q}$ .

Analogously,  $\lambda \notin \{ik\pi, k \in \mathbb{Z}^*\}$  if  $\xi \notin \mathbb{Q}$ .

Now, since  $\lambda \notin \{k\pi, k \in \mathbb{Z}^*\}$  and  $\lambda \notin \{ik\pi, k \in \mathbb{Z}^*\}$ , (4.14) becomes:

$$\begin{cases} P_1 = R_1 = 0 \\ Q_1 = -\frac{\alpha\mu}{2a\lambda^3}u(\xi) \frac{\sinh(\lambda(\xi - 1))}{\sinh(\lambda)} \text{ and } S_1 = \frac{\alpha\mu}{2a\lambda^3}u(\xi) \frac{\sin(\lambda(\xi - 1))}{\sin(\lambda)} \\ P_2 = Q_1 \sinh(\lambda\xi), Q_2 = Q_1 \cosh(\lambda\xi), R_2 = S_1 \sin(\lambda\xi) \text{ and } S_2 = S_1 \cos(\lambda\xi). \end{cases} \quad (4.17)$$

Combining that with classical trigonometric formulae leads to:

$$u_g(\lambda, x) := Q_1 \sinh(\lambda x) + S_1 \sin(\lambda x), \forall x \in (0; 1). \quad (4.18)$$

In particular  $u(\xi) = u_p(\lambda, \xi) + u_g(\lambda, \xi) = u_g(\lambda, \xi)$  *i.e.*

$$u(\xi) = -\frac{\alpha\mu}{2a\lambda^3}u(\xi) \left( \frac{\sinh(\lambda(\xi - 1))}{\sinh(\lambda)} \sinh(\lambda\xi) - \frac{\sin(\lambda(\xi - 1))}{\sin(\lambda)} \sin(\lambda\xi) \right). \quad (4.19)$$

This implies the following equation, since  $u(\xi) \neq 0$ :

$$2a\lambda^3 \sinh(\lambda) \sin(\lambda) - \alpha\mu [\sin(\lambda) \sinh(\lambda\xi) \sinh(\lambda(1 - \xi)) - \sinh(\lambda) \sin(\lambda\xi) \sin(\lambda(1 - \xi))] = 0. \quad (4.20)$$

Multiplying both sides by  $\lambda \neq 0$ , replacing  $(a\lambda^4)$  by  $(-b\mu - \mu^2)$  and dividing by  $\mu \neq 0$  leads to the characteristic equation (4.7).

At last,  $u(\xi)$  can be rewritten as:

$$\begin{aligned} u(\xi) = u_g(\xi) &= -\frac{\sinh(\lambda(\xi - 1))}{\sinh(\lambda)} \sinh(\lambda\xi) + \frac{\sin(\lambda(\xi - 1))}{\sin(\lambda)} \sin(\lambda\xi) \\ &= \frac{2a\lambda^3}{\alpha\mu} \text{ due to the characteristic equation.} \end{aligned} \quad (4.21)$$

Hence the expression for the eigenfunction (4.8). Note that the factor  $\frac{1}{|\lambda|^2} e^{(\xi-2)|\lambda|}$  is aimed at making the function bounded with  $\lambda$  in  $\mathcal{H}$ .

The last part of the proof concerns the multiplicity of the eigenvalues. Since the characteristic equation (4.7) is invariant under the transform  $\lambda \mapsto i\lambda$ , to each root  $\lambda$  such that  $\lambda \in \{z \in \mathbb{C} - \{k\pi, k \in \mathbb{Z}\}, \exists l \in \mathbb{Z}, \arg(z) \in [-\pi/4 + 2l\pi; \pi/4 + 2l\pi)\}$  correspond three other roots which are  $\pm i\lambda$  and  $-\lambda$ . Up to a multiplicative constant, the expression for the eigenfunction is invariant under the transform  $\lambda \mapsto i\lambda$ , thus the geometric multiplicity of  $\mu$  as an eigenvalue is always 1.

Now (3.16) is equivalent to

$$\mu(\lambda) = \frac{-b + \delta(\lambda)}{2} \quad (4.22)$$

where  $\pm\delta(\lambda)$  are the (complex) square roots of  $(b^2 - 4a\lambda^4)$ .

The value  $\mu = -b/2$  is an eigenvalue of the operator  $\mathcal{A}$  if and only if  $\delta(\lambda) = 0$  which is the case for the value  $\alpha$  defined by (4.9). In that case, the algebraic multiplicity is at least 2.

If  $\alpha$  takes any other value, it is at least 1. □

**Remark 4.4.** The characteristic equation (4.7) (found for  $\alpha \neq 0$  and  $\xi \notin \mathbb{Q}$ ) becomes  $\sinh(\lambda)\sin(\lambda) = 0$ , if  $\alpha$  is replaced by 0 in (4.7). The solutions of this equation are  $\lambda = n\pi$ ,  $n \in \mathbb{Z}$  and  $\lambda = in\pi$ ,  $n \in \mathbb{Z}$ . Now, if  $\alpha$  tends to 0,  $\alpha b < 6$  so  $\mu = -b$  is not an eigenvalue of  $\mathcal{A}$  and since 0 is not an eigenvalue of  $\mathcal{A}$ ,  $\lambda \neq 0$ . Thus the corresponding eigenvalues are those given for Case  $\alpha = 0$  in Theorem 4.2, which is coherent.

**Remark 4.5.** (Case  $\beta > 0$ ) The eigenfunctions of Problem (1.4)–(1.6) with the additional term  $\beta\delta_\xi$  are unchanged (*cf.* (4.8)) and the characteristic equation (4.7) becomes:

$$2a\lambda^3 \sinh(\lambda)\sin(\lambda) - (\alpha\mu + \beta) [\sin(\lambda)\sinh(\lambda\xi)\sinh(\lambda(1-\xi)) - \sinh(\lambda)\sin(\lambda\xi)\sin(\lambda(1-\xi))] = 0. \quad (4.23)$$

Multiplying both sides by  $\lambda \neq 0$ , replacing  $(a\lambda^4)$  by  $(-b\mu - \mu^2)$  and dividing by  $\mu \neq 0$  leads to the following characteristic equation:

$$2(\mu + b)\sinh(\lambda)\sin(\lambda) + \left(\alpha + \frac{\beta}{\mu}\right)\lambda [\sin(\lambda)\sinh(\lambda\xi)\sinh(\lambda(1-\xi)) - \sinh(\lambda)\sin(\lambda\xi)\sin(\lambda(1-\xi))] = 0. \quad (4.24)$$

The asymptotic behaviour of the eigenvalues is often useful for the study of stabilization. The following Proposition gives what we call the asymptotic characteristic equation.

**Proposition 4.6.** (*Large eigenvalues of the operator  $\mathcal{A}$* )

Let  $\mathcal{A}$  be the operator defined by (2.3) and (2.4) in Section 2. Denote by  $\mu_n$ ,  $n \in \mathbb{Z}$  the eigenvalues of  $\mathcal{A}$  and by  $\lambda$  the complex numbers satisfying  $\lambda^4 = -\frac{b\mu + \mu^2}{a}$ . One of these values of  $\lambda$  is such that  $\exists k \in \mathbb{Z}, \arg(\lambda) \in [-\pi/4 + 2k\pi; \pi/4 + 2k\pi)$ . The others are  $-\lambda$  and  $\pm i\lambda$  and they are still denoted by  $\lambda$ .

Assume that  $\alpha > 0$ ,  $\alpha b < 6$  and  $\xi \notin \mathbb{Q}$ .

Then  $\mu$  is a large eigenvalue of  $\mathcal{A}$  if and only if the corresponding values of  $\lambda$  satisfy the asymptotic characteristic equation:

$$\sinh(\lambda)\sin(\lambda) = 0. \quad (4.25)$$

The algebraic multiplicity of  $\mu$  as a large eigenvalue of  $\mathcal{A}$  is exactly 1 and its geometric multiplicity is 1.

*Proof.* We start from the characteristic equation (4.20) in which the trigonometric functions are replaced by:

$$\begin{cases} \sin(\lambda) = \frac{e^{\Im(\lambda)}}{2i} \left[ -e^{-i\Re(\lambda)} + e^{-2\Im(\lambda)} e^{i\Re(\lambda)} \right] \\ \sinh(\lambda) = \frac{e^{\Re(\lambda)}}{2} \left[ e^{i\Im(\lambda)} - e^{-2\Re(\lambda)} e^{-i\Im(\lambda)} \right] \end{cases} \quad (4.26)$$

where  $\Re$  and  $\Im$  are the real and imaginary parts respectively. Denoting by

$$\begin{cases} S(\lambda) := -e^{-i\Re(\lambda)} + e^{-2\Im(\lambda)} e^{i\Re(\lambda)} \\ Sh(\lambda) := e^{i\Im(\lambda)} - e^{-2\Re(\lambda)} e^{-i\Im(\lambda)} \end{cases} \quad (4.27)$$

the characteristic equation (4.20) is equivalent to:

$$\frac{e^{\Im(\lambda)+\Re(\lambda)}}{4i} [2a\lambda^3 S(\lambda)Sh(\lambda) - \alpha\mu[S(\lambda)Sh(\lambda\xi)Sh(\lambda(\xi-1)) - Sh(\lambda)S(\lambda\xi)S(\lambda(\xi-1))]] = 0 \quad (4.28)$$

and, since  $\lambda = 0$  is excluded here, it is also:

$$S(\lambda)Sh(\lambda) - \alpha \frac{\mu}{2a\lambda^3} [S(\lambda)Sh(\lambda\xi)Sh(\lambda(\xi-1)) - Sh(\lambda)S(\lambda\xi)S(\lambda(\xi-1))] = 0. \quad (4.29)$$

The expression  $[S(\lambda)Sh(\lambda\xi)Sh(\lambda(\xi-1)) - Sh(\lambda)S(\lambda\xi)S(\lambda(\xi-1))]$  is bounded with respect to  $\lambda$ , if  $\Re(\lambda)$  and  $\Im(\lambda)$  tend to  $+\infty$  and the quotient  $\frac{\mu}{2a\lambda^3}$  tends to zero if  $|\lambda|$  tends to  $+\infty$ . The other cases ( $\Re(\lambda)$  and  $\Im(\lambda)$  tend to  $-\infty$ ,  $\Re(\lambda)$  tends to  $\pm\infty$  while  $\Im(\lambda)$  tends to  $\mp\infty$ ) are treated similarly. Hence (4.25).

The multiplicity of the large eigenvalues follows from the fact that the roots of  $\sin(\lambda) = 0$  are simple as well as those of  $\sinh(\lambda) = 0$ . The result is proved using Rouché's Theorem. We follow the proof of Lemma 2.4 of the paper by Abdallah *et al.* [1].

We denote by  $h_\alpha(\lambda)$  the left-hand side of (4.29) and define, for  $N$  large enough, the curves:

$$\begin{cases} \Gamma_{\pm,n}^S := \left\{ z / |z \pm n\pi| = \frac{C_0}{n} \right\} \\ \Gamma_{\pm,n}^{Sh} := \left\{ z / |z \pm in\pi| = \frac{C'_0}{n} \right\}. \end{cases} \quad (4.30)$$

Our aim is to choose  $C_0$  such that, by Rouché's Theorem,  $h_\alpha$  has the same roots as  $\sinh$  (resp. as  $\sin$ ) inside the curve  $\Gamma_{\pm,n}^{Sh}$  (resp.  $\Gamma_{\pm,n}^S$ ) for every  $n > N$  where  $N$  is large enough.

The proof is written for  $z \in \Gamma_{+,n}^{Sh}$ . The rest is analogous.

The first step is to show that, if  $z \in \Gamma_{+,n}^{Sh}$ , then  $|Sh(z)| \geq \frac{C_0}{|z|}$ .

Writing  $z = in\pi + \rho_n e^{i\theta}$  with  $\rho_n = \frac{C_0}{n}$  and using trigonometric formulae lead to

$$\begin{aligned} |\sinh(z)|^2 &= |\sinh(in\pi + \rho_n e^{i\theta})|^2 = \sinh^2(\rho_n \cos(\theta)) \cos^2(\rho_n \cos(\theta)) + \cosh^2(\rho_n \cos(\theta)) \sin^2(\rho_n \cos(\theta)) \\ &= \rho_n^2 + o(1). \end{aligned} \quad (4.31)$$

Now  $|z|^2 = \rho_n^2 \cos^2(\theta) + (n\pi + \rho_n \sin(\theta))^2 \leq (n\pi + \rho_n \sin(\theta))^2$  and

$$\frac{C_0^2}{|z|^2} \leq \frac{C_0^2}{(n\pi + \rho_n \sin(\theta))^2} = \frac{C_0^2}{n^2 \pi^2} + o\left(\frac{1}{n^2}\right) = \frac{C_0^2}{n^2} + o(1) = |\sinh(z)|^2. \quad (4.32)$$

Since  $|Sh(z)| = 2e^{-\Re(z)}|\sinh(z)|$  and  $\Re(z) = \rho_n \cos(\theta)$  tends to zero when  $n$  tends to  $+\infty$

$$\exists N_1, n \geq N_1 \Rightarrow 2e^{-\Re(z)} \geq 1 \text{ and } |Sh(z)| \geq |\sinh(z)| \geq \frac{C_0}{|z|}. \quad (4.33)$$

The second step is to show that, if  $z \in \Gamma_{+,n}^{Sh}$ , then  $|h_\alpha(z) - S(z)Sh(z)| \leq |S(z)Sh(z)|$ .

By definition of  $h_\alpha(z)$ , it holds:

$$|h_\alpha(z) - S(z)Sh(z)| \leq \frac{1}{|z|} \cdot \left| \frac{\alpha\mu}{2a\lambda^2} \right| \cdot |S(z)| \cdot \left| Sh(\xi z)Sh((\xi - 1)z) - Sh(z) \frac{S(\xi z)S((\xi - 1)z)}{S(z)} \right|. \quad (4.34)$$

If  $z \in \Gamma_{+,n}^{Sh}$ ,  $\Im(z) = n\pi + \frac{C_0}{n} \sin(\theta) \geq n\pi - \frac{C_0}{n}$ .

The expression  $\frac{\alpha\mu}{2a\lambda^2} \left( Sh(\xi z)Sh((\xi - 1)z) - Sh(z) \frac{S(\xi z)S((\xi - 1)z)}{S(z)} \right)$  is bounded with respect to  $\Im(z)$ . Let us denote by  $C_0$  the real number such that

$$\exists C_1, \forall z \in \Gamma_{+,n}^{Sh}, \Im(z) \geq C_1 \Rightarrow \left| \frac{\alpha\mu}{2a\lambda^2} \right| \cdot \left| Sh(\xi z)Sh((\xi - 1)z) - Sh(z) \frac{S(\xi z)S((\xi - 1)z)}{S(z)} \right| \leq C_0. \quad (4.35)$$

At last, there exists  $N_2$  such that  $N_2\pi - \frac{C_0}{N_2} > C_1$ .

If  $N \geq \max\{N_1, N_2\}$ ,  $z \in \Gamma_{+,n}^{Sh} \Rightarrow |h_\alpha(z) - S(z)Sh(z)| \leq |S(z)Sh(z)|$ .  $\square$

**Remark 4.7.** The large roots of the characteristic equation (4.7) are close to the roots of (4.25), that is to say, either  $(n\pi)$  or  $(in\pi)$  with  $n \in \mathbb{Z}^*$ . The large eigenvalues  $\mu$  of the operator  $\mathcal{A}$  with  $\alpha > 0$  are thus near the eigenvalues of the operator  $\mathcal{A}_0$  given in Theorem 4.2.

**Remark 4.8.** (Case  $\beta > 0$ ) The latter Proposition still holds for Problem (1.4)–(1.6) with the additional term  $\beta\delta_\xi$ .

Another natural question which arises in this context is that of the continuity of the eigenvalues with respect to the parameter  $\alpha$ .

**Proposition 4.9.** (Continuity of the eigenvalues and eigenvectors of the operator  $\mathcal{A}$  with respect to  $\alpha$ )

Let  $\mathcal{A}$  be the operator defined by (2.3) and (2.4) in Section 2. Denote by  $\mu$  an eigenvalue of  $\mathcal{A}$  and by  $\lambda$  the complex numbers satisfying  $\lambda^4 = -\frac{b\mu + \mu^2}{a}$ . One of these values of  $\lambda$  is such that  $\exists k \in \mathbb{Z}, \arg(\lambda) \in [-\pi/4 + 2k\pi; \pi/4 + 2k\pi)$ . The others are  $-\lambda$  and  $\pm i\lambda$  and they are still denoted by  $\lambda$ .

Assume that  $\alpha > 0$ ,  $ab < 6$  and  $\xi \notin \mathbb{Q}$ .

Then  $\mu$  depends continuously on the parameter  $\alpha$  and the eigenvector  $\Phi(\mu, \cdot)$  converges pointwise to the eigenvector  $\Phi_n^\pm$  (up to a multiplicative constant) when  $\alpha$  tends to zero.

*Proof.* To state the continuity of  $\mu$  with respect to  $\alpha$ , we follow the proof of Remark 1 of the paper by Abdallah et al. [1]. We denote by  $g_\alpha(\lambda)$  the left-hand side of (4.7).

For a fixed  $\alpha$ , denote by  $\lambda_0$  a root of  $\mathcal{A}$ . Since  $\lambda_0$  is isolated, there exists  $\rho > 0$  such that

$$g_\alpha(z) \neq 0, \forall z \in \mathbb{C}, \text{ such that } 0 < |z - \lambda_0| < \rho. \quad (4.36)$$



Now,  $g_\alpha$  is a continuous function. Indeed the functions  $\lambda \mapsto \lambda^3$ ,  $\lambda \mapsto \sin(\lambda)$  and  $\lambda \mapsto \sinh(\lambda)$  are continuous functions on  $\mathbb{C}$ . As for  $\mu$ , it is the root of the second degree equation:

$$\mu^2 + b\mu + a\lambda^4 = 0. \quad (4.37)$$

Thus it is a continuous function of the coefficients of this equation *i.e.* a continuous function of the variable  $\lambda$ . Setting  $D = \{z \in \mathbb{C}, \text{ such that } |z - \lambda_0| = \rho\}$ , the continuity of  $g_\alpha$  implies that there exists a positive real number  $\kappa$  such that  $|g_\alpha(z)| \geq \kappa, \forall z \in D$ .

For a fixed positive real number  $\epsilon_0$ , we consider the mapping of two variables

$$H : [0, \epsilon_0] \times D \rightarrow \mathbb{C} : (\epsilon, z) \mapsto g_{\alpha+\epsilon}(z) - g_\alpha(z). \quad (4.38)$$

Since it is a uniformly continuous function and since  $H(0; z) = 0$  for all  $z \in D$ , we deduce the existence of a positive real number  $\delta$  such that

$$|H(\epsilon, z)| < \kappa, \forall (\epsilon, z) \in [0; \delta] \times D. \quad (4.39)$$

The last two estimates imply that

$$|g_{\alpha+\epsilon}(z) - g_\alpha(z)| < |g_\alpha(z)|, \forall (\epsilon, z) \in [0; \delta] \times D. \quad (4.40)$$

Hence Rouché's theorem allows to conclude that  $g_{\alpha+\epsilon}$  has the same number of roots as  $g_\alpha$  for all  $\epsilon \in [0; \delta]$ . Thus the root  $\lambda$  of  $g_\alpha$  is a continuous function of  $\alpha$  and, by composition,  $\mu$  is also a continuous function of  $\alpha$ .

To finish with the proof, note that, when  $\alpha$  tends to zero,  $\alpha b < 6$  then  $\mu + b \neq 0$  and the characteristic equation (4.7) becomes:  $\sin(\lambda) \sinh(\lambda) = 0$ . Thus  $\lambda$  tends either to  $(n\pi)$  or to  $(in\pi)$  with  $n \in \mathbb{Z}^*$ .

If it tends to  $(n\pi)$  for example,  $\sin(\lambda)$  tends to zero and the eigenfunction tends to:

$$\lim_{\lambda \rightarrow n\pi} \phi(\lambda, x) := \frac{1}{n^2\pi^2} \cdot e^{n\pi(\xi-2)} \sinh(n\pi) \sin(n\pi(1-\xi)) \sin(n\pi x). \quad (4.41)$$

Indeed the other terms tend to zero since  $\sin(\lambda)$  tends to zero and the following two expressions are bounded with  $\lambda$ :

$$\left\{ \begin{array}{l} \frac{1}{|\lambda|^2} \cdot e^{(\xi-2)|\lambda|} \sinh(\lambda) [\sin(\lambda(x-\xi)) - \sinh(\lambda(x-\xi))] \cdot H(x, \xi) \\ \frac{1}{|\lambda|^2} \cdot e^{(\xi-2)|\lambda|} \sinh(\lambda(1-\xi)) \sinh(\lambda x). \end{array} \right. \quad (4.42)$$

And if  $\lambda$  tends to  $(in\pi)$ ,  $\sinh(\lambda)$  tends to zero and the eigenfunction tends to:

$$\begin{aligned} \lim_{\lambda \rightarrow in\pi} \phi(\lambda, x) &:= \frac{1}{n^2\pi^2} \cdot e^{n\pi(\xi-2)} \sin(in\pi) \sinh(in\pi(1-\xi)) \sinh(in\pi x) \\ &= -i \frac{1}{n^2\pi^2} \cdot e^{n\pi(\xi-2)} \sinh(n\pi) \sin(n\pi(1-\xi)) \sin(n\pi x). \end{aligned} \quad (4.43)$$

Indeed the other terms tend to zero since  $\sinh(\lambda)$  tends to zero and the following two expressions are bounded with  $\lambda$ :

$$\begin{cases} \frac{1}{|\lambda|^2} \cdot e^{(\xi-2)|\lambda|} \sin(\lambda) [\sin(\lambda(x-\xi)) - \sinh(\lambda(x-\xi))] \cdot H(x, \xi) \\ \frac{1}{|\lambda|^2} \cdot e^{(\xi-2)|\lambda|} \sin(\lambda(1-\xi)) \sin(\lambda x). \end{cases} \quad (4.44)$$

Hence the announced result concerning the eigenfunctions.  $\square$

## 5. LOCALIZATION OF THE EIGENVALUES OF THE OPERATOR $\mathcal{A}$ FOR SMALL VALUES OF $\alpha$

The aim of this Section is to determine the localization of the eigenvalues of the operator  $\mathcal{A}$  for small values of  $\alpha$  compared with the case  $\alpha = 0$  already studied in the paper by Ammari *et al.* [4] and recalled in Theorem 4.2.

**Theorem 5.1.** (*Localization of the eigenvalues of the operator  $\mathcal{A}$  for small values of  $\alpha$* )

Let  $\mathcal{A}$  be the operator defined by (2.3) and (2.4) in Section 2 with  $\xi \notin \mathbb{Q}$ . Denote by  $\mu$  an eigenvalue of  $\mathcal{A}$  and by  $\Re(\mu)$  its real part.

The case  $\alpha = 0$  is already known (cf. Thm. 4.2). If  $b \leq 2\sqrt{a}\pi^2$ , then

$$\sup_{\mu \in \sigma(\mathcal{A})} \Re(\mu) = -\frac{b}{2}. \quad (5.1)$$

If  $b > 2\sqrt{a}\pi^2$ , then

$$\sup_{\mu \in \sigma(\mathcal{A})} \Re(\mu) = \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}). \quad (5.2)$$

Now, for any  $\epsilon > 0$ , there exists  $\alpha_0$  such that, if  $0 < \alpha < \alpha_0$ , it holds: if  $b \leq 2\sqrt{a}\pi^2$ , then

$$\Re(\mu) \leq -\frac{b}{2} + \epsilon. \quad (5.3)$$

If  $b > 2\sqrt{a}\pi^2$ , then

$$\Re(\mu) \leq \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}) + \epsilon. \quad (5.4)$$

*Proof.* Since the study is restricted here to the case of the small non-vanishing values of  $\alpha$ , the eigenvalues  $\mu$  of the operator  $\mathcal{A}$  are known to be such that the corresponding values of  $\lambda$  satisfy  $\lambda \notin \{k\pi, k \in \mathbb{Z}\}$ ,  $\lambda \notin \{ik\pi, k \in \mathbb{Z}\}$  and the characteristic equation:

$$2a\lambda^3 \sinh(\lambda) \sin(\lambda) - \alpha\mu [\sin(\lambda) \sinh(\lambda\xi) \sinh(\lambda(1-\xi)) - \sinh(\lambda) \sin(\lambda\xi) \sin(\lambda(1-\xi))] = 0 \quad (5.5)$$

(cf. Proof of Thm. 4.3).

The first remark is that this equation is invariant under the transform  $\lambda \mapsto i\lambda$  so it is enough to consider  $\lambda \in \{z \in \mathbb{C} - \{k\pi, k \in \mathbb{Z}\}, \exists l \in \mathbb{Z}, \arg(z) \in [-\pi/4 + 2l\pi; \pi/4 + 2l\pi]\}$ .

In order to compute the first degree Mac Laurin polynomial for the function  $\lambda(\alpha)$  *i.e.* to find  $\lambda_1$  such that  $\lambda = n\pi + \lambda_1\alpha + o(\alpha)$  around the value  $\alpha = 0$ , we write all the first degree Mac Laurin polynomials for the functions involved in equation (5.5) as functions of  $\lambda$  around  $\lambda = n\pi$  ( $n \in \mathbb{N}^*$ ):

$$\begin{cases} \lambda^3 = n^3\pi^3 + (3n^2\pi^2\lambda_1)\alpha + o(\alpha) \\ \sinh(\lambda) = \sinh(n\pi) + \cosh(n\pi)\lambda_1\alpha + o(\alpha) \\ \sin(\lambda) = (-1)^n\lambda_1\alpha + o(\alpha). \end{cases} \quad (5.6)$$

Hence  $2a\lambda^3 \sinh(\lambda) \sin(\lambda) = [2a(-1)^n n^3 \pi^3 \sinh(n\pi)\lambda_1]\alpha + o(\alpha)$ . Now,  $\exists n_0$  such that  $2an_0^2\pi^2 < b < 2a(n_0 + 1)^2\pi^2$  and for  $n > n_0$  :

$$\mu_n^{\pm, -} = \frac{1}{2} \left[ -b \pm i\sqrt{4an^4\pi^4 - b^2} \right] + o(1). \quad (5.7)$$

Inserting all these results into the characteristic equation (5.5) leads, for  $n$  such that  $n > n_0$ , to:

$$\lambda_1 = \frac{[-b \pm i\sqrt{4an^4\pi^4 - b^2}] \sin(n\pi\xi) \sin(n\pi(\xi - 1))}{4(-1)^n an^3 \pi^3} = \frac{[-b \pm i\sqrt{4an^4\pi^4 - b^2}] (-1)^n \sin^2(2n\pi\xi)}{8(-1)^n an^3 \pi^3}. \quad (5.8)$$

At last,  $\lambda = n\pi + \lambda_1\alpha + o(\alpha)$  is put into the expression of  $\mu_n^{\pm}$ , which leads after some calculations to:

$$\begin{cases} \mu_n^{\pm} = \frac{1}{2}[-b \pm i\sqrt{4an^4\pi^4 - b^2}] \pm i \frac{\sin^2(2n\pi\xi)}{2\sqrt{4an^4\pi^4 - b^2}} [-b \pm i\sqrt{4an^4\pi^4 - b^2}] \cdot \alpha + o(\alpha), \forall n > n_0 \\ \mu_n^{\pm} = \frac{1}{2}[-b \pm \sqrt{b^2 - 4an^4\pi^4}] \mp \frac{\sin^2(2n\pi\xi)}{2\sqrt{b^2 - 4an^4\pi^4}} [-b \pm \sqrt{b^2 - 4an^4\pi^4}] \cdot \alpha + o(\alpha), \forall n, 1 \leq n \leq n_0. \end{cases} \quad (5.9)$$

The real part of  $\mu_n^{\pm}$  is:

$$\begin{cases} \Re(\mu_n^{\pm}) = -\frac{b}{2} - \frac{1}{2} \sin^2(2n\pi\xi) \cdot \alpha + o(\alpha), \forall n > n_0 \\ \Re(\mu_n^{\pm}) = \frac{1}{2}[-b \pm \sqrt{b^2 - 4an^4\pi^4}] \pm \frac{\sin^2(2n\pi\xi)}{2\sqrt{b^2 - 4an^4\pi^4}} [b \mp \sqrt{b^2 - 4an^4\pi^4}] \cdot \alpha + o(\alpha), \forall n, 1 \leq n \leq n_0. \end{cases} \quad (5.10)$$

The eigenvalues of  $\mathcal{A}$  for small positive values of  $\alpha$  and large values of  $n$  are on the left of those we have for  $\alpha = 0$  (whose real part is  $-b/2$ ).

As for the case of small values of  $n$  ( $1 \leq n < n_0$ ), since  $[b \mp \sqrt{b^2 - 4an^4\pi^4}] > 0$  and  $\sqrt{b^2 - 4an^4\pi^4} > 0$ , the eigenvalues  $\mu_n^+$  of  $\mathcal{A}$  for small positive values of  $\alpha$  are on the right of those we have for  $\alpha = 0$  and the eigenvalues  $\mu_n^-$  of  $\mathcal{A}$  for small positive values of  $\alpha$  are on the left of those we have for  $\alpha = 0$ .

Moreover,  $0 \leq \frac{1}{2} \sin^2(2n\pi\xi) \cdot \alpha \leq \frac{\alpha}{2}$  and  $\left| \frac{\sin^2(2n\pi\xi)}{2\sqrt{b^2 - 4an^4\pi^4}} [b \mp \sqrt{b^2 - 4an^4\pi^4}] \right|$  is bounded for  $n \leq n_0$ .

Now, it holds, for all  $1 \leq n \leq n_0$ :

$$-b < \frac{1}{2}(-b - \sqrt{b^2 - 4a\pi^4}) < \frac{1}{2}(-b - \sqrt{b^2 - 4an^4\pi^4}) < -\frac{b}{2} < \frac{1}{2}(-b + \sqrt{b^2 - 4an^4\pi^4}) < \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}) < 0. \quad (5.11)$$

Thus, for all  $n \geq 1$  and  $\alpha$  sufficiently small:

$$\Re(\mu_n^{\pm}) < \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}). \quad (5.12)$$

Note that, if  $\xi$  belonged to  $\mathbb{Q}$ , there would exist  $(m, n) \in (\mathbb{N}^*)^2$  such that  $\xi = \frac{m}{n}$  and  $\sin(2n\pi\xi)$  would vanish (as well as  $\sin(2kn\pi\xi)$  for  $k \in \mathbb{Z}$ ). More calculations would be required to get the third (at least) degree MacLaurin polynomials of the functions involved in the characteristic equation since the second degree term for  $\lambda$  and for  $\mu_n$  also vanishes.  $\square$

**Remark 5.2.** (Case  $\beta > 0$ ) The first-degree Taylor polynomial linear approximation of  $\Re(\mu_n^\pm)$  if we add the term  $\beta\delta_\xi$  to get Problem (1.4)–(1.6) contains the additional term:

$$\frac{\pm 2 \sin^2(2n\pi\xi)}{\sqrt{b^2 - 4an^4\pi^4}} \cdot \beta. \quad (5.13)$$

Note that the real part of the eigenvalues  $\mu_n^+$  is larger than that of the same eigenvalues with  $\beta = 0$ . This confirms what we said in the introduction: the stiffness term does not improve the decreasing of the energy. If the SMA cables restrain the vibrations of girder bridges, it only comes from the damper term  $\alpha u_t \delta_\xi$ .

All that is true if  $\sin^2(2n\pi\xi)$  does not vanish, which happens for some values of  $\xi$ . That is why this result does not contradict the experiments of [12].

## 6. ENERGY DECREASING

**Definition 6.1.** (Functions  $\Psi_n^\pm$ )

Consider  $\mu$  such that  $\lambda$  defined by  $\lambda^4 = -\frac{b\mu + \mu^2}{a}$  satisfies the characteristic equation (4.7). It is already known that  $\mu$  is an eigenvalue of the operator  $\mathcal{A}$  and that one of these values of  $\lambda$  is such that  $\exists k \in \mathbb{Z}, \arg(\lambda) \in [-\pi/4 + 2k\pi; \pi/4 + 2k\pi)$ . The others are  $-\lambda$  and  $\pm i\lambda$ . Consider this  $\lambda$ . If its modulus  $|\lambda|$  tends to  $+\infty$ , then there exists  $n \in \mathbb{Z}^*$  such that  $\lambda$  tends to  $n\pi$  (respectively to  $in\pi$ ).

This integer  $n$  depends on  $\mu$ .

The function  $\Psi_n^\pm$  is defined on  $(0; 1)$  by

$$\Psi_n^\pm := \frac{1}{n^2\pi^2} e^{n\pi(\xi-2)} \sinh(n\pi) \sin(n\pi(1-\xi)) \sin(n\pi x)(1, \mu^\pm) \quad (6.1)$$

(respectively by

$$\Psi_n^\pm := -i \frac{1}{n^2\pi^2} e^{(\xi-2)|\lambda|} \sin(n\pi) \sin(n\pi x)(1, \mu^\pm). \quad (6.2)$$

**Theorem 6.2.** (Riesz basis for the operator  $\mathcal{A}$ )

Let  $\Phi(\mu, \cdot)$  still be defined as in Theorem 4.3 and  $\Psi_n^\pm$  given by the above definition. There exists  $n_0 \geq 1$ , such that

$$\sum_{|n| \geq n_0} \|\Phi(\mu, \cdot) - \Psi_n^\pm\|_{\mathcal{H}}^2 < \infty. \quad (6.3)$$

Thus, the root eigenvectors of  $\mathcal{A}$  form a Riesz basis of  $\mathcal{H}$ .

*Proof.* First the  $\Psi_n^\pm$  form an orthogonal basis (see the paper by Ammari *et al.* [4]).

The inner product in  $\mathcal{H}$  has two terms. Since the eigenvector  $\Phi(\mu, \cdot)$  is defined by  $\Phi(\mu, \cdot) := (\phi(\lambda, \cdot))(1, \mu)$  and since  $\mu$  has the same behaviour as  $\lambda^2$  for large values of  $\lambda$ , it is enough to consider the first term in the inner product.

Now the eigenfunction defined by (4.8) is made of three terms. Let us start with the third one. The second one is analogous with easier calculations. That is why we do not give details for this second term.

Let us prove that, if  $\lambda \rightarrow n\pi$  ( $n \in \mathbb{N}^*$ ), then

$$\int_0^1 \left| e^{(\xi-2)|\lambda|} \sin(\lambda) \sinh(\lambda(1-\xi)) \sinh(\lambda x) + e^{(\xi-2)n\pi} \sin(n\pi) \sinh(n\pi(1-\xi)) \sinh(n\pi x) \right|^2 dx = O\left(\frac{1}{n^2}\right). \quad (6.4)$$

The integrand is of the form  $|A \cdot B - A_0 \cdot B_0|^2 \leq 2|A - A_0|^2 \cdot |B|^2 + 2|A_0|^2 \cdot |B - B_0|^2$  with

$$|A_0|^2 := |e^{(\xi-2)n\pi} \sin(n\pi) \sinh(n\pi(1-\xi))|^2 = O(e^{-2n\pi}). \quad (6.5)$$

Since, for  $x \in (0, 1)$ ,  $\sinh(\lambda x) = O(e^{2|\Re(\lambda)|}) = O(e^{2n\pi})$  (cf. proof of Proposition 4.6 about the large eigenvalues), it holds:

$$\int_0^1 |B|^2 := \int_0^1 |\sinh(\lambda x)|^2 dx = O(e^{2n\pi}). \quad (6.6)$$

Thus it is enough to show the following two estimates:

$$|A - A_0|^2 := \left| e^{(\xi-2)|\lambda|} \sin(\lambda) \sinh(\lambda(1-\xi)) - e^{(\xi-2)n\pi} \sin(n\pi) \sinh(n\pi(1-\xi)) \right|^2 = O\left(\frac{1}{n^2} e^{-2n\pi}\right) \quad (6.7)$$

$$\int_0^1 |B - B_0|^2 := \int_0^1 |\sinh(\lambda x) - \sinh(n\pi x)|^2 dx = O\left(\frac{1}{n^2} e^{2n\pi}\right). \quad (6.8)$$

Coming back to the definition of the trigonometric functions, it holds:

$$\sin(\lambda) - \sin(n\pi) = \frac{1}{2i} \left[ (e^{i\Re(\lambda)} e^{-\Im(\lambda)} - e^{in\pi}) - (e^{-i\Re(\lambda)} e^{\Im(\lambda)} - e^{-in\pi}) \right]. \quad (6.9)$$

Then, using the mean-value Theorem as well as the asymptotic behaviour of the difference  $(\lambda - n\pi)$  obtained in the proof of Proposition 4.6 about the large eigenvalues:

$$\begin{aligned} & |e^{i\Re(\lambda)} e^{-\Im(\lambda)} - e^{in\pi}|^2 = |(e^{i\Re(\lambda)} - e^{in\pi}) + (e^{-\Im(\lambda)} - 1) \cdot e^{i\Re(\lambda)}|^2 \\ & \leq 2(e^{-\Im(\lambda)} - 1)^2 + 2|\cos(\Re(\lambda)) - \cos(n\pi) + i(\sin(\Re(\lambda)) - \sin(n\pi))|^2 \\ & \leq 2(e^{-\Im(\lambda)} - 1)^2 + 2(\cos(\Re(\lambda)) - \cos(n\pi))^2 + 2(\sin(\Re(\lambda)) - \sin(n\pi))^2 \\ & \leq \frac{C_0^2}{n^2} \exp\left(2\left(\frac{C_0}{n} + o(1/n)\right)\right) + 4(\Re(\lambda) - n\pi)^2 = O\left(\frac{1}{n^2}\right). \end{aligned} \quad (6.10)$$

Now (6.7) is of the form  $|E \cdot F - E_0 \cdot F_0|^2 \leq 2|E - E_0|^2 \cdot |F|^2 + 2|E_0|^2 \cdot |F - F_0|^2$  with

$$\begin{cases} |F - F_0|^2 := |e^{i\Re(\lambda)} e^{-\Im(\lambda)} - e^{in\pi}|^2 = O\left(\frac{1}{n^2}\right) \\ |E_0|^2 := |e^{(\xi-2)n\pi} \sinh(n\pi(1-\xi))|^2 = O(e^{-2n\pi}) \\ |E - E_0|^2 := |e^{(\xi-2)|\lambda|} \sinh(\lambda(1-\xi)) - e^{(\xi-2)n\pi} \sinh(n\pi(1-\xi))|^2 = O(e^{-2n\pi}) \\ |\sin(\lambda)|^2 = \left| \sin\left(n\pi + O\left(\frac{1}{n}\right)\right) \right|^2 = \left| (-1)^n \sin\left(O\left(\frac{1}{n}\right)\right) \right|^2 = O\left(\frac{1}{n^2}\right). \end{cases} \quad (6.11)$$

Hence (6.7). Analogously, it holds:

$$\sinh(\lambda x) - \sinh(n\pi x) = \frac{1}{2} \left[ (e^{\Re(\lambda)x} e^{i\Im(\lambda)x} - e^{n\pi x}) - (e^{-\Re(\lambda)x} e^{-i\Im(\lambda)x} - e^{-n\pi x}) \right]. \quad (6.12)$$

Using once more the mean-value Theorem as well as the asymptotic behaviour of the difference  $(\lambda - n\pi)$  obtained in the proof of Proposition 4.6 about the large eigenvalues:

$$\begin{aligned} & \left| e^{\Re(\lambda)x} e^{i\Im(\lambda)x} - e^{n\pi x} \right|^2 = \left| (e^{\Re(\lambda)x} - e^{n\pi x}) + (e^{i\Im(\lambda)x} - 1) \cdot e^{\Re(\lambda)x} \right|^2 \\ & \leq 2 (e^{\Re(\lambda)x} - e^{n\pi x})^2 + 2e^{2\Re(\lambda)x} \left[ (\cos(\Im(\lambda)x) - 1)^2 + \sin^2(\Im(\lambda)x) \right] \\ & \leq 4 \frac{C_0^2}{n^2} \exp \left( 2 \left( n\pi + \frac{C_0}{n} + o(1/n) \right) \right) \\ & = O \left( \frac{1}{n^2} e^{2n\pi} \right). \end{aligned} \quad (6.13)$$

Hence (6.8).

The first term of the eigenfunction has to be considered.

Since  $|\sin(\lambda)|^2 = O\left(\frac{1}{n^2}\right)$  and  $|e^{(\xi-2)|\lambda|} \sinh(\lambda) \sinh(\lambda(x-\xi)) H(x, \xi)|^2 = O(1)$  for  $x > \xi$ , it holds, if  $\lambda \rightarrow n\pi$  ( $n \in \mathbb{N}^*$ ):

$$\int_0^1 \left| e^{(\xi-2)|\lambda|} \sin(\lambda) \sinh(\lambda) \sinh(\lambda(x-\xi)) H(x, \xi) \right|^2 dx = O\left(\frac{1}{n^2}\right). \quad (6.14)$$

This ends the proof.  $\square$

Finally, the energy is proved to decay exponentially and the localization of the eigenvalues for small values of  $\alpha$  leads to a lower bound for the decay rate in that case. This bound is the optimal decay rate obtained for  $\alpha = 0$  in the paper by Ammari *et al.* [4] which means that the decay rate for small values of  $\alpha > 0$  is equal to that with  $\alpha = 0$  (*i.e.* without SMA cables). SMA cables do not restrain the vibrations of girder bridges if  $\alpha$  is close to zero. This is not in contradiction with the experiments of [12] but it means that the simple fact of adding SMA cables may not be enough.

**Theorem 6.3.** (*Exponential stability and decay rate for small values of  $\alpha$* )

*System (1.1)–(1.3) (presented in the introduction) is exponentially stable and, for any  $\epsilon > 0$ , there exists  $\alpha_0$  such that for any  $0 < \alpha < \alpha_0$ , for any solution  $U$  of (2.5) with initial data  $U_0 \in D(\mathcal{A})$ , there exist constants  $C > 0$  and  $\omega_0(\epsilon) < 0$  depending on  $a, b$  and  $\alpha$  such that:*

$$E(t) \leq C e^{2\omega_0(\epsilon)t} \|U_0\|_{D(\mathcal{A})}^2, \forall t > 0,$$

where  $E(t)$  is defined by (1.7) and

$$\begin{cases} \omega_0(\epsilon) := -\frac{b}{2} + \epsilon, & \text{if } b \leq 2\sqrt{a}\pi^2 \\ \omega_0(\epsilon) := \frac{1}{2}(-b + \sqrt{b^2 - 4a\pi^4}) + \epsilon, & \text{if } b > 2\sqrt{a}\pi^2. \end{cases} \quad (6.15)$$

*Proof.* According to Theorem 6.2, the system of eigenvectors of  $\mathcal{A}$  constitutes a Riesz basis. Consequently, by a standard argument (see the proof of Theorem 2.5 of the paper by Ammari *et al.* [4], for example), the optimal

energy decay rate is identified with the supremum of the real part of the eigenvalues of  $\mathcal{A}$ .

Thus the result follows from Theorem 5.1.  $\square$

**Remark 6.4.** (Case  $\alpha = 0$ )

In the paper by Ammari *et al.* [4], the systems of eigenvectors of  $\mathcal{A}$  and  $\mathcal{A}_0$  (obtained with  $\alpha = 0$ ) are proved to be quadratically close in  $VxL_2(0, 1)$  using the explicit expression of the eigenfunctions of both operators. Thus, it follows from Theorem 3 in Appendix D of the book by Pöschel and Trubowitz [19], that the system of eigenvectors of  $\mathcal{A}_0$  constitutes a Riesz basis. Consequently, by a standard argument (see the proof of Theorem 2.5 in the paper by Ammari *et al.* [4]), the optimal energy decay rate is identified with the supremum of the real part of the eigenvalues of  $\mathcal{A}_0$ . With the notation of Theorem 6.3, it means that  $\omega = \omega_0(0)$  when  $\alpha = 0$ .

## 7. CONCLUSION

One could have thought that adding the damping term  $\alpha\delta_\xi u_t$ ,  $\alpha > 0$ , would logically increase the decay rate of the energy compared with the case  $\alpha = 0$  already studied by Ammari *et al.* in [4].

In fact this is not obvious. If  $\alpha$  is small, we have proved that the decay rate does not change. Is it clear that it will increase with  $\alpha$ ? This is not clear either.

As announced in the introduction, let us prove that, if  $\alpha = 0$ , increasing the value of  $b$  may not increase the decay rate of the energy. Indeed, in that situation, the energy decreases exponentially and the decay rate is given by the maximum of the real part of  $\mu_n^+$  which is

$$\Re\mu_n^+ = \frac{1}{2} \left( -b + \sqrt{b^2 - 4an^4\pi^4} \right) = \frac{-2an^4\pi^4}{b + \sqrt{b^2 - 4an^4\pi^4}}, n = 1, 2, \dots, n_0 \quad (7.1)$$

where  $b \in \mathbb{R}^{+*} - \{2\sqrt{an^2\pi^2}, n \in \mathbb{Z}^*\}$  and  $n_0$  is such that  $2\sqrt{an_0^2\pi^2} < b < 2\sqrt{a}(n_0 + 1)^2\pi^2$ . Thus

$$\Re\mu_n^+ \leq \frac{-2a\pi^4}{b + \sqrt{b^2 - 4an_0^4\pi^4}}, n = 1, 2, \dots, n_0 \quad (7.2)$$

and the last bound tends to zero, if  $b$  tends to  $+\infty$ .

In conclusion, the decay rate of the energy may be better with  $\alpha > 0$  (with  $\alpha$  large enough) than with  $\alpha = 0$  *i.e.* Shape Memory Alloy cables may restrain the vibrations of girder bridges effectively.

The experiments of Liu *et al.* [12] would be confirmed mathematically if we were able to prove that there exist  $\alpha > 0$ ,  $\beta > 0$ ,  $\xi \in (0; 1)$  and  $\omega_1(\alpha, \beta, \xi) < \omega_0(0)$ , such that all the roots of the characteristic equation (4.24) are such that the corresponding values of  $\mu$  satisfy  $\Re(\mu) \leq \omega_1(\alpha, \beta, \xi)$ . We conjecture that this requires a value for  $\xi$  which is excluded in our theorems.

*Acknowledgements.* Thanks to S. Nicaise who suggested this problem to me as well as the reflection around the fact that increasing the value of  $b$  when  $\alpha = 0$  may not increase the decay rate. I also thank the referees for their attentive reading of the manuscript and their valuable remarks. At last I would like to dedicate this paper to D. Mercier with whom I have enjoyed doing mathematics for so many years. I wish him a happy retirement!

## REFERENCES

- [1] F. Abdallah, D. Mercier and S. Nicaise, Spectral analysis and exponential or polynomial stability of some indefinite sign damped problems. *Evol. Equ. Control Theory* **2** (2013) 1–33.
- [2] M. Akil, Y. Chitour, M. Ghader and A. Wehbe, Stability and exact controllability of a timoshenko system with only one fractional damping on the boundary. *Asymp. Anal.* **119** (2020) 221–280.
- [3] M. Akil, I. Issa and A. Wehbe, Energy decay of some boundary coupled systems involving wave Euler-Bernoulli beam with one locally singular fractional Kelvin-Voigt damping. *Math. Control Related Fields* (2021) doi:10.3934/mcrf.2021059.

- [4] K. Ammari, M. Dimassi and M. Zerzeri, The rate at which energy decays in a viscously damped hinged Euler-Bernoulli beam. *J. Diff. Equ.* **257** (2014) 3501–3520.
- [5] K. Ammari and M. Tucsnak, Stabilization of Bernoulli-Euler beams by means of a pointwise feedback force. *SIAM J. Control Optim.* **39** (2000) 1160–1181.
- [6] K. Ammari and M. Tucsnak, Stabilization of second order evolution equations by a class of unbounded feedbacks. *ESAIM: COCV* **6** (2001) 361–386.
- [7] W. Arendt and C.J.K. Batty, Tauberian theorems and stability of one-parameter of semi-groups. *Trans. Amer. Math. Soc.* **305** (2) (1988) 837–852.
- [8] A. Borichev and Y. Tomilov, Optimal polynomial decay of functions and operator semigroups. *Math. Ann.* **347** (2010) 455–478.
- [9] C. Castro and E. Zuazua, Exact boundary controllability of two Euler-Bernoulli beams connected by a point mass. *Math. Comput. Model.* **32** (2000) 955–969.
- [10] S. Cox and E. Zuazua, The rate at which energy decays in a damped string. *Partial Differ. Equ.* **19** (1994) 213–243.
- [11] B. Dekoninck and S. Nicaise, Control of networks of Euler-Bernoulli beams. *ESAIM : COCV* **4** (1999) 57–81.
- [12] A-R. Liu, C-H. Liu, J-Y. Fu, Y-L. Pi, Y-H. Huang and J-P. Zhang, A method of reinforcement and vibration reduction of girder bridges using shape memory alloy cables. *Int. J. Struct. Stab. Dyn.* **17** (2017) 1750076.
- [13] L. Majkut, Eigenvalue based inverse model of beam for structural modification and diagnostics. Part I: Theoretical formulation. *Latin Am. J. Solids Struct.* **7** (2009) 423–436.
- [14] D. Mercier and V. Régnier, Spectrum of a network of Euler-Bernoulli beams. *J. Math. Anal. Appl.* **337** (2007) 174–196.
- [15] D. Mercier and V. Régnier, Control of a network of Euler-Bernoulli beams. *J. Math. Anal. Appl.* **342** (2008) 874–894.
- [16] D. Mercier and V. Régnier, Boundary controllability of a chain of serially connected Euler-Bernoulli beams with interior masses. *Collect. Math.* **60** (2009) 307–334.
- [17] D. Mercier and V. Régnier, Decay rate of the Timoshenko system with one boundary damping. *Evol. Equ. Control Theory* **8** (2019) 423–445.
- [18] A. Pazy, Semigroups of linear operators and applications to partial differential equations. *Appl. Math. Sci.* **44** (1983).
- [19] J. Pöschel, E. Trubowitz, Vol. 130 of *Inverse Spectral Theory, Pure and Applied Mathematics*. Academic Press, Boston, MA (1987).



**Please help to maintain this journal in open access!**

This journal is currently published in open access under the Subscribe to Open model (S2O). We are thankful to our subscribers and supporters for making it possible to publish this journal in open access in the current year, free of charge for authors and readers.

Check with your library that it subscribes to the journal, or consider making a personal donation to the S2O programme by contacting [subscribers@edpsciences.org](mailto:subscribers@edpsciences.org).

More information, including a list of supporters and financial transparency reports, is available at <https://edpsciences.org/en/subscribe-to-open-s2o>.